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Characterization of Unsaturated Soils Using Elastic and Electromagnetic Waves

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**CHARACTERIZATION OF UNSATURATED SOILS USING ELASTIC AND
ELECTROMAGNETIC WAVES**

A Dissertation

Submitted to the Graduate Faculty of the
Louisiana State University and
Agricultural and Mechanical College
in partial fulfillment of the
requirements for the degree of
Doctor of Philosophy

in

The Department of Civil and Environmental Engineering

By
Bashar Alramahi
B.S., Birzeit University, 2002
M.S., Louisiana State University, 2004
December, 2007

To the memory of my father

&

To my mother and brother

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ABSTRACT

Recent advances in laboratory instruments and techniques enabled researchers to explore new aspects of the behavior of geomaterials and perform measurements that would be otherwise impossible to acquire using traditional geotechnical laboratory techniques. This dissertation is focused on utilizing elastic and electromagnetic wave measurements and SMCT imaging to non-destructively characterize different aspects of the behavior of unsaturated soils. A model that relates P-wave velocity in soils to the volumetric water content was used to develop a new methodology to determine the in situ density and moisture content of soils. It was numerically and experimentally verified to assess its validity and range of applicability. On the other hand, a triaxial apparatus that enables the measurement of P- and S- wave velocities in unsaturated soil specimens under controlled net stress and matric suction was also developed. Several verification experiments were performed using the apparatus and the results were compared to theoretical models as well as previous experimental results. Moreover, a drying cell was used to examine the effect of the presence of fine clay and silt particles on the elastic waves' velocity and the small strain stiffness of unsaturated soils. The results were confirmed by analyzing SMCT images of similar samples at different drying stages.

The proposed methodology yielded good predictions of the density and the moisture content of soils. However, different experimental and numerical error sources caused the predicted density and moisture content values to slightly differ from the measured values. For the majority of the tested specimens, the density was estimated within $\pm 10\%$ of the measured values while the water content was estimated within $\pm 20\%$. On the other hand, the experimental results from the new triaxial apparatus showed a

significant effect of matric suction on the recorded wave velocities. It was also documented that wave velocity values increase with increasing percentage of fine silt particles in the specimens.

The results of the drying cell experiments as well as the SMCT image analysis showed the profound effect of the presence of fine silt and clay particles on the small strain stiffness of unsaturated soils. The density of the pore fluid increased during drying due to the higher fine concentration. The concentration of fine particles was found to be significantly higher at areas close to the interparticle contact than in pore bodies away from the contacts causing an increase in the interparticle contact stiffness.

CHAPTER 1

INTRODUCTION

1.1 Background

In common engineering applications, soils are assumed to be either dry or fully saturated. However, in most practical cases, soils exist in a state between those two extreme states known as unsaturated condition. The presence of both air and water in the pore space introduces a new force balance that adds a very important component to the interparticle forces known as capillary forces. Since the strength and deformation characteristics of unsaturated soils are controlled by many factors including particle shapes, particle interaction, applied stresses and capillary forces, it is very important to understand the contribution of each of these factors to the overall behavior of geomaterials.

Recent advances in instrumentation and measurement techniques permitted better characterization of soil properties such as stiffness at different strain levels starting from large strains acquired by conventional laboratory testing to very small strains acquired by wave propagation techniques. Recent research advances gave the opportunity to bridge the gap between “low” stiffnesses typically measured using traditional laboratory tests and “high” stiffnesses (usually denoted G_{\max}) measured using geophysical techniques (Wood, 2004). Figure 1.1 illustrates the levels of shear stiffness corresponding to different strain levels, and the measurement methods used for different ranges.

Elastic wave propagation through soils is one of the most common geophysical techniques used to determine soils small strain stiffness. The velocity of elastic waves in granular media is determined by its stiffness; therefore, these waves provide a unique tool to monitor the effects of different factors such as degree of saturation, suction, or applied stresses on the small

strain stiffness in soils and other granular materials. Moreover, because of the very small strains associated with this type of waves, measurement could be performed without disturbing the soils or altering any ongoing processes.

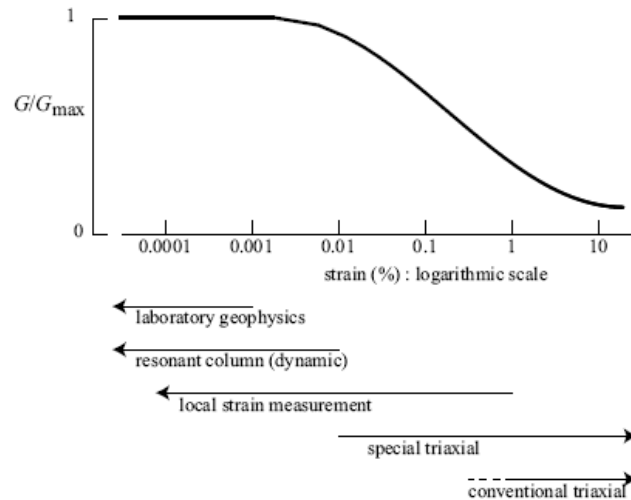


Figure 1.1 Shear stiffness (G) corresponding to different strain levels and the measurement methods (Atkinson and Sallfors, 1991)

In this dissertation, elastic waves along with electromagnetic waves and x-ray tomography were used in various studies to determine the effects of different factors such as the degree of saturation, matric suction and fine particle migration on the small strain stiffness of unsaturated soils. They were also utilized to develop a new non-destructive methodology to determine soils in situ density and moisture content.

1.2 Dissertation Objectives

This research is aimed at utilizing elastic and electromagnetic waves to characterize different aspects of the behavior of unsaturated soils. The main objectives of the dissertation are:

1. Assess the accuracy and range of applicability of a new methodology to determine the field density and moisture content of soils using elastic and electromagnetic waves.

2. Develop a new apparatus for the measurement of P- and S- wave velocities in unsaturated soil samples under controlled net stress and matric suction.
3. Perform verification experiments using that apparatus to determine the effects of matric suction on the P- and S- wave velocities through soils and compare the measurements to theoretical models.
4. Determine the effects of the presence of fine silt and clay particles in the pore fluid on the small strain stiffness of unsaturated soils during drying.
5. Monitor the changes in the properties of pore fluid containing fine silt and clay particles at different drying stages using synchrotron x-ray computed tomography.

1.3 Dissertation Layout

This dissertation consists of six chapters. The second Chapter presents a literature review covering the major subjects presented in the dissertation. Different aspects of unsaturated soils are discussed in this chapter including effective stress expressions, soil suction, and soil water characteristic curves. Chapter 2 describes the different types of elastic waves in soils and their measurement techniques, and presents concepts related to electromagnetic waves propagation in soils including the Time Domain Reflectometry (TDR) technique. Finally, Chapter 2 presents some of the current practices for determining the in situ density and moisture content in soils.

Chapter 3 describes a new methodology for determining in situ density and moisture content of soils using a combination of elastic and electromagnetic waves. The chapter includes a derivation of a semi-empirical model for the P-wave velocity in soils as a function of its volumetric water content, along with a description of the suggested methodology. It also presents the results of numerical and experimental studies performed to assess the validity and range of applicability of the suggested methodology.

The fourth Chapter is dedicated to describing a new apparatus for the measurement of P- and S- wave velocities in unsaturated soils under controlled net stress and matric suction. It includes a detailed description of the apparatus along with the results of experiments performed using this apparatus. It also presents a model for the determination of S- wave velocity considering the effects of matric suction along with a comparison between the experimentally measured and wave velocity values based on the proposed model.

The fifth Chapter presents a study of the effects of fine particle migration on the small strain stiffness of unsaturated soils during drying. It includes the description of a drying cell that was developed to perform several drying experiments to determine the effects of the presence of fine silt and clay particles in the pore fluid on the overall small strain stiffness of soils. The results of the experiments performed using this cell along with the measured wave velocities and the calculated stiffnesses are interpreted using synchrotron x-ray computed tomography. These micro-tomography results show how the changes in pore fluid properties during drying increase the fine concentration at the interparticle contacts changing the nature of the forces and increasing the soils stiffness.

The sixth chapter presents the conclusion of the performed studies along with the recommended future work.

CHAPTER 2

LITERATURE REVIEW

2.1 Stress, Suction, and Elastic Wave Propagation in Soils

2.1.1 Effective Stress in Soils

As a multi-phase particulate medium, the behavior of soils depends on the interaction of all physical and chemical micro scale forces between the soil particles, air, and water. Interparticle forces can be either repulsive or attractive. Repulsive forces take effect at very small separations between particles. For example, solvation forces act at separations less than 20 Å. They are caused by the “hard shell” effect of the molecules as the particles are getting closer. On the other hand, electrostatic forces known as Born repulsion act at even smaller separations. It is caused by the overlap of electron clouds between the adjacent atoms at contact points between particles. Although these forces act at the atomic level, not the particle level, they are often referred to as ‘interparticle reaction forces’ (Santamarina et al. 2001). Moreover, the hydration energy of particle surfaces and inter-layer cations causes large repulsive forces at small distances between unit layers. Hydration repulsions decay rapidly with separation distance varying inversely as the square of the distance (Mitchell 1993).

On the other hand, when particle edges and corners are oppositely charged, electrostatic attraction develops between them. This attraction is believed to be one of the causes of the adherence of dry fine grained particles. At any one time there may be more electrons on one side of the atomic nucleus than the other, creating instantaneous dipoles whose oppositely charged ends attract each other. These are the van der Waals bonds. Electromagnetic attractions are significant long range attractions in clays resulting from frequency dependent dipole interactions. Although the magnitude of these forces are difficult to compute in natural soils, it is known that

they vary inversely as the fourth power of the separation distance between particles (Mitchell 1993).

Another source of strong long range attraction, known as Coulombian force, develops between oppositely charged ions. The strength of this force is a function of the amplitude of the charge of both ions and inversely related to the separation distance. However, if both ions were similarly charged a repulsive force of similar strength will develop between them.

Cementation is a chemical bonding mechanism that can be treated as a short range attraction where covalent and ionic bonds occur at spacing less than 3 Å. Very high contact stresses between particles could expel adsorbed water and cations and cause mineral surfaces to come close together providing an opportunity for cold welding. However, the absence of such cohesion in over-consolidated silts and sands argues against such pressure induced bonding. Since particle size is much larger than clays, interparticle contact forces should be orders of magnitude larger as a result of fewer contacts per unit volume in coarser soils. Cementation may also occur naturally from precipitation of calcite, silica, alumina or iron oxides or other organic or inorganic oxides.

Another source of particle attraction is the capillary stresses in unsaturated soils. Since water is attracted to soil particles and because water can develop surface tension, capillary menisci form between particles in partially saturated soils. Capillary forces will be discussed later in this section. A quantitative measure of the interaction of the mentioned forces is not presently possible. However, a simplified expression for the inter-granular pressure can be derived by considering the force balance through a horizontal surface through a saturated soil at some depth (Figure 2.1). The vertical equilibrium of forces can be expressed as:

$$\sigma_a + Aa + A'a_c = ua + Ca_c \quad (2.1)$$

where a is the average total cross sectional area, a_c is the effective area of interparticle contacts, σ is the applied vertical stress, u is the hydrostatic pressure, A is the long term attractive stress (i.e., van der Waals and electrostatic attractions), A' is the short range attractive stress (i.e., primary valence bonding and cementation), and C is the repulsive stress resulting from hydration and Born repulsion. Dividing all the terms in Equation 2.1 by “ a ” converts all the forces to stresses per unit area of the cross section (Equation 2.2).

$$\sigma = (C - A') \frac{a_c}{a} + u - A \quad (2.2)$$

The term $(C - A') \frac{a_c}{a}$ represents the intergrain force divided by the gross area, referred to as interangular pressure (σ_i'). Equation 2.2 can be re-written as:

$$\sigma_i' = \sigma + A - u \quad (2.3)$$

It should be noted that u is the pore water pressure at the true interparticle zones, which is different than the pore water pressure measured by a piezometer or other pressure measurement devices (u_0). To enable a more accurate estimate of effective stresses Bernoulli's equation is used to derive a relationship between u and u_0 as follows:

$$u = u_0 - Z\gamma_w - h_s\gamma_w \quad (2.4)$$

where Z is the elevation difference between the piezometer and the point in question, γ_w is the unit weight of water, and h_s is the osmotic head resulting from the difference in ionic concentrations between points near soil particles and points away from them. Assuming no elevation difference, and using the expression for u from Equation 2.4, Equation 2.4 can be expressed as:

$$\sigma_i' = \sigma + A - u_0 + h_s\gamma_w \quad (2.5)$$

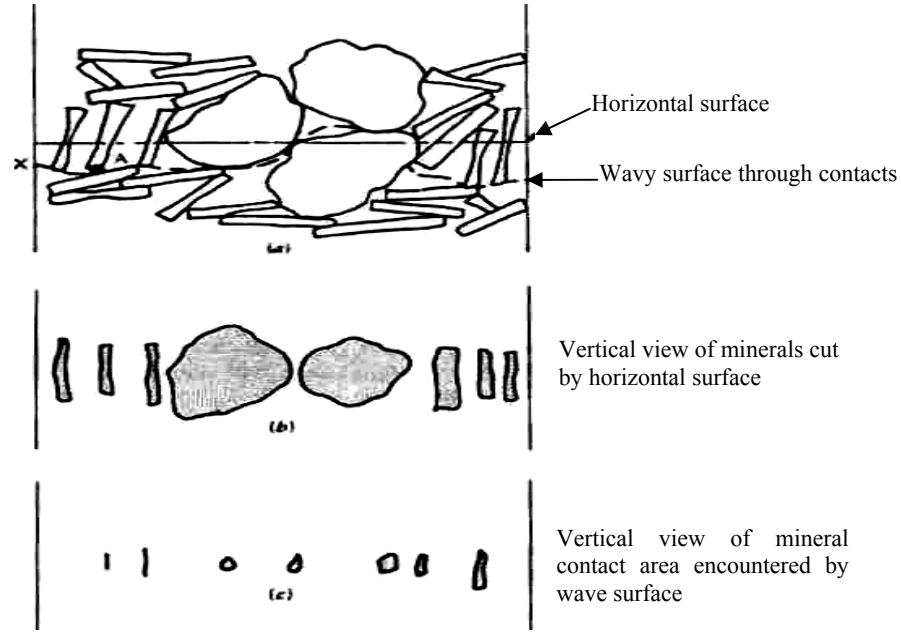


Figure 2.1. Surface through a soil mass (Mitchell, 1993)

Since the salt concentration between particles will be greater than at points away from them, the osmotic term $h_s \gamma_w$ will be negative. It has been labeled as R in previous studies (Lambe 1960, Mitchell 1962, Balasubramonian, 1972). Consequently, Equation 2.5 can be written as:

$$\sigma'_i = \sigma + A - u_0 - R \quad (2.6)$$

Comparing the result in Equation 2.6 with the commonly used Terzaghi effective stress equation ($\sigma' = \sigma - u_0$), the difference in the two expressions can be written as:

$$\sigma'_i - \sigma' = A - R \quad (2.7)$$

Both equations give similar results if A and R were very small, as in the case of sands, silts, low plasticity clays, or when $A \approx R$. However, in cases where either A or R is large or when they are significantly different, the results of the two equations would not be similar. Skempton (1960) stated that although the simple Terzaghi's effective stress equation does not give the true effective stress, it gives an excellent approximation for the case of saturated soils.

2.1.2 Effective Stress in Unsaturated Soils

In fully saturated or fully dry soils, the Terzaghi's effective stress principle defines the effective stress as the difference between the total stress and the pore fluid pressure. In the case of three phase systems (unsaturated soils) the Terzaghi's effective stress principle is no longer valid. The presence of air in the soil pores induces a new force balance that incorporates capillary forces acting at the soil-air-water interface. Such forces have a direct effect on the forces acting on particle contacts and highly influence the macroscopic behavior of the soil.

The effective stress in unsaturated soils is dependent on the relative amounts of total stress (σ), air pressure (u_a) and water pressure (u_w). Many researchers proposed different forms of effective stress relationships. Table 2.1 presents a summary of some of the equations suggested for the effective stress in unsaturated soils.

The equation suggested by Bishop (1959) is the most general form since it includes a term for the pressure of the gas phase. It retains the understanding that the effective stress is made up of two components, one resulting from total normal pressure ($\sigma - u_a$), and the other from pressures exerted by the fluid in the soil pores ($u_a - u_w$) (Jennings and Burland, 1962). The magnitude of the χ parameter is equal to unity for a saturated soil and null for a dry soil. Its values at different degrees of saturation can be determined experimentally (e.g., Figure 2.2). It also depends on the wetting history, loading path, soil type and internal structure of the soil (Fredlund and Rahardjo, 1993).

The first four Equations in Table 2.1 are equivalent when the pore air pressure is the same ($\beta' = \chi = \psi = \beta$). Only Bishop's form references the total and pore water pressure to the pore air pressure. The other equations simply use gauge pressures which are referenced to the external atmospheric pressure (Fredlund and Rahardjo, 1993). Jennings and Burland (1962) suggested

that under a critical degree of saturation, Bishop's equation did not provide an adequate description for the relationship between volume change and effective stress for most soils. The critical degree of saturation was estimated to be approximately 20% for silts and sands and 85% to 90% for clays. Moreover, Morgenstern (1979) stated that the Bishop's equation proved to "have little impact on practice. The parameter χ when determined for volume change behavior was found to differ when determined for shear strength".

Table 2.1. Suggested effective stress equations for unsaturated soils.

Researcher	Equation	Notes
Croney (1958)	$\sigma' = \sigma - \beta' u_w$	B' = a measure of the number of bonds under tension.
Bishop (1959)	$\sigma' = (\sigma - u_a) + \chi(u_a - u_w)$	X = a parameter related to the degree of saturation of the soil
Aitchinson (1961)	$\sigma' = \sigma - \Psi \cdot p''$	P'' = pore water pressure deficiency. Ψ = a parameter with values ranging from zero to one
Jennings (1961)	$\sigma' = \sigma - \beta \cdot p''$	P'' = negative pore water pressure taken as a positive value. β = an experimentally measured statistical factor of the same type as the contact area.
Richards (1966)	$\sigma' = \sigma - u_a + \chi_m(h_m - u_a) + \chi_s(h_s - u_a)$	χ_m, χ_s = effective stress parameters for matric and solute suctions, respectively. h_m, h_s = matric and solute suctions, respectively.
Atchinson (1965)	$\sigma' = \sigma + \chi_m p_m'' + \chi_s p_s''$	χ_m, χ_s = soil parameters ranging from 0 to 1 depending on the stress path. P_m'', p_s'' = matric and solute suctions, respectively.

Another important limitation in Bishop's equation is the use of soil properties in the description of a stress state. Morgenstern (1979) stated that since the effective stress is a stress variable, it should be related to equilibrium considerations only, however, the parameter χ bears on constitutive behavior. Constitutive behavior is normally used in geomaterials to link equilibrium considerations to deformations, but in this case it is introduced directly into the stress variable. Moreover, this equation mixes global and local conditions. Both pressure and total stress are boundary actions in Terzaghi's effective stress equation for saturated media; however, the pore water pressure in unsaturated soils causes a local action at the particle level (Cho and Santamarina, 2001).

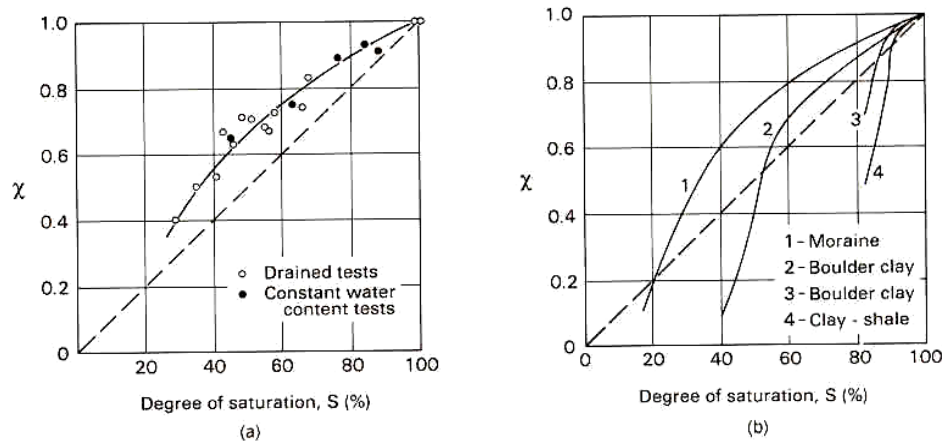


Figure 2.2. The relationship between the χ parameter and the degree of saturation S . (a) χ values for a silt (after Donald, 1961); (b) χ values for compacted soils (after Blight, 1961).

Many researchers suggested using independent stress state variables such as $(\sigma - u_a)$ and $(u_a - u_w)$. Numerous pairs of variables have been proposed in the literature (e.g., Coleman 1962; Bishop and Blight 1963; and Matyas and Radhakrishna 1968). After conducting a theoretical stress analysis, Fredlund and Morgenstern (1977) concluded that any two of three possible normal stress variables can be used to describe the stress state of unsaturated soils. The following three different combinations can be used as stress state variables: $(\sigma - u_a)$ and $(u_a - u_w)$,

$(\sigma - u_w)$ and $(u_a - u_w)$ or $(\sigma - u_a)$ and $(\sigma - u_w)$. In a three dimensional stress analysis, the stress state variables of unsaturated soil form two independent stress tensors. Using the first couple of stress variables, the net normal stress tensor in the Cartesian coordinate system is defined as:

$$\begin{pmatrix} \sigma_x - u_a & \tau_{yx} & \tau_{zx} \\ \tau_{xy} & \sigma_y - u_a & \tau_{zy} \\ \tau_{xz} & \tau_{yz} & \sigma_z - u_a \end{pmatrix} \quad (2.8)$$

and the matric suction tensor is

$$\begin{pmatrix} u_a - u_w & 0 & 0 \\ 0 & u_a - u_w & 0 \\ 0 & 0 & u_a - u_w \end{pmatrix} \quad (2.9)$$

Under stable field conditions, the total normal stress exceeds the pore air pressure which in turn exceeds the pore water pressure, $(\sigma > u_a > u_w)$. This condition is generally true for field conditions. However, under unique loading conditions such as blasting during mining operations, the pore air pressure is suddenly increased to a value exceeding the total stress, resulting in a limiting stress state condition (Fredlund and Rahardjo, 1993). Consequently, the soil skeleton disintegrates when its tensile strength is reached.

Burland (1964) indicated that the state of stress in a soil remains constant if the values of the net stress $(\sigma - u_a)$ and suction $(u_a - u_w)$ are not altered. Fredlund and Morgenstern (1977) acknowledged that suitable sets of stress state variables are those that produce no distortion or volume change of an element when the individual components of the stress state variables are modified but the stress state variables themselves are kept constant. They conducted a series of “null” experiments to test different stress state variables. Similar “null” tests related to the shear

strength of unsaturated silt were also performed by Bishop and Donald (1961). It led to the development of the axis translation technique that will be discussed later in Section 2.1.6.

2.1.3 Micro-Scale Force Balance in Unsaturated Soils

Surface tension (T_s) is a property of the air water interface that results from the intermolecular forces acting on molecules in the contractile skin (Fredlund and Rahardjo 1993). It is often defined as the maximum energy level a fluid can store without breaking apart (Lu and Likos, 2004). It has units of energy per surface area (J/m^2) or force per length (N/m). The surface tension of pure water depends on temperature as illustrated in Figure 2.3.

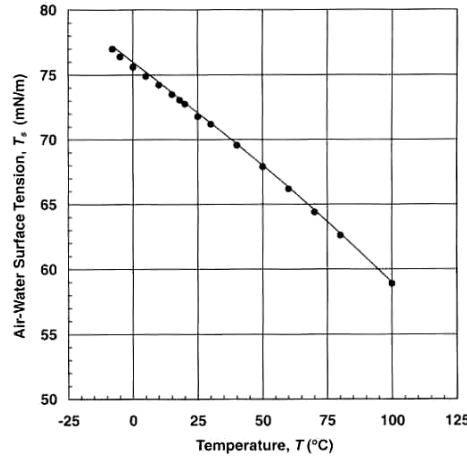


Figure 2.3. Surface tension of air-water interface as a function of temperature (Weast et al., 1981)

Figure 2.4 depicts an idealized geometry of the air-water interface between two spherical soil grains with radii r_1 and r_2 . Considering the horizontal force balance in the horizontal direction, three forces are considered: surface tension along the interface described by r_1 , surface tension along the interface described by r_2 , and air and water pressure applied on either side of the interface. The first and second forces can be expressed as (Lu and Likos, 2004):

$$F_1 = 4r_3T_s \sin \alpha \quad (2.10)$$

$$F_2 = -4r_1T_s \sin \alpha \quad (2.11)$$

The projection of the air and water pressure u_a and u_w in the horizontal direction (assuming $r_2 = r_3$) is:

$$F_3 = 4r_1r_2(u_a - u_w)\sin\alpha \quad (2.12)$$

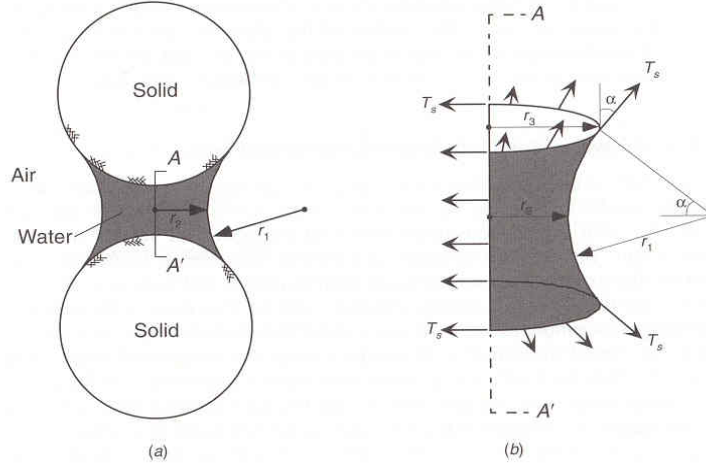


Figure 2.4. Idealized air-water interface geometry in unsaturated soil. (a) water meniscus between two spherical soil particles and (b) free-body diagram for water meniscus (Lu and Likos, 2004).

Balancing the three forces leads to:

$$T_s(r_2 - r_1) = (u_a - u_w)r_1r_2 \quad (2.13)$$

or

$$u_a - u_w = T_s \left(\frac{1}{r_1} + \frac{1}{r_2} \right) \quad (2.14)$$

Equation 2.14, known as Laplace's equation, describes the relationship between the matric suction ($u_a - u_w$) and the water meniscus geometry (r_1 and r_2). It shows that the pressure across the air-water-solid interface depends on the relative magnitudes of r_1 and r_2 , and could be positive, negative, or zero. In most cases, the matric suction is positive as r_1 is mostly less than r_2 under unsaturated conditions (Lu and Likos, 2004).

Cho and Santamarina (2001) conducted a particle level study on unsaturated soils. They considered different modes of packing and derived their effective stress equations. Two cases were considered: simple cubic packing (coordination number 6, void ratio 0.91), and tetrahedral packing (coordination number 12, void ratio 0.34). The effective stresses for these two cases are shown as follows:

$$\sigma'_{eq} = \frac{\pi \cdot T_s}{4R} \left[2 - \left(\frac{8}{9} G_s w \right)^{1/4} \right] \quad \text{Simple cubic packing} \quad (2.15)$$

$$\sigma'_{eq} = 2\sqrt{2} \sigma'_{eq, cubic} = \frac{\pi \cdot T_s}{\sqrt{2}R} \left[2 - \left(\frac{8}{9} G_s w \right)^{1/4} \right] \quad \text{Tetrahedral packing} \quad (2.16)$$

where R is the radius of the spherical particles, G_s is the specific gravity, and w is the water content. The authors also noted that the soil structure for uniformly graded sands at the microscale can be constrained between the simple cubic packing and the tetrahedral packing. Equations 2.15 and 2.16 are plotted in Figure 2.5 assuming particle radius (R) of 1 μm , specific gravity of 2.65 and surface tension of 72.75 mN/m.

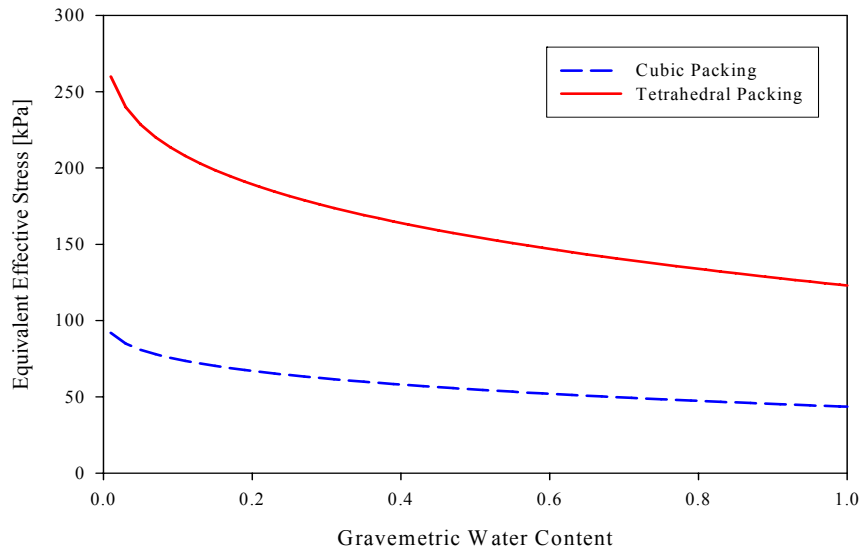


Figure 2.5. Equivalent effective stress for simple cubic and tetrahedral packing.

2.1.4 Soil Suction

Total soil suction quantifies the thermodynamic potential of soil pore water relative to a reference potential of free water. Free water is defined as water containing no dissolved solutes, having no interactions with other phases that impart curvature to the air water interface, and having no external forces other than gravity (Lu and Likos, 2004). The total suction ψ of a soil consists of two components: osmotic suction π and the matric suction ($u_a - u_w$). Osmotic suction is the component that results from the presence of dissolved salts in the pore water. It is present in saturated as well as unsaturated soils. If the salt content in a soil changes, there will be a change in its overall volume and shear strength (Fredlund and Rahardjo, 1993).

The matric suction arises from the combined effects of capillarity ($u_a - u_w$) and short term adsorption. It is of greater importance when studying the engineering behavior of unsaturated soils. Figure 2.6 illustrates the relative importance of changes in osmotic suction as compared to matric suction (Fredlund and Rahardjo, 1993). One can notice that unless fine silt or clay particles were present in the pore liquid, osmotic suction stays almost unchanged throughout the range of water content, and the change in the total suction can be attributed solely to the change in matric suction.

2.1.5 Soil Water Characteristic Curves

The soil water characteristic curve (SWCC) is an important relationship in unsaturated soils. It describes the relationship between the soil suction and water content. In more specific terms, the SWCC describes the thermodynamic potential of the soil pore water relative to that of free water as a function of the amount of water adsorbed by the soil system (Lu and Likos, 2004). The amount of water can be expressed as gravimetric water content (w), volumetric water content (θ) or degree of saturation (S). Typical SWCC's for sand, silt, and clay are shown in Figure 2.7.

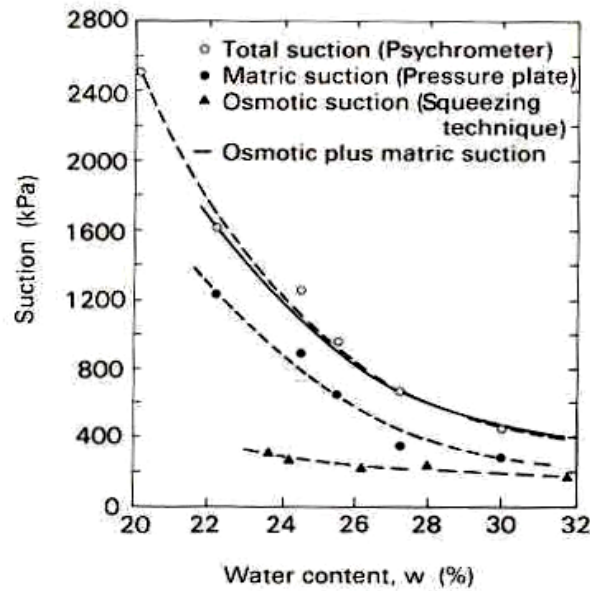


Figure 2.6. Total matric and osmotic suction measurements on compacted Regina Clay (Krahn and Fredlund, 1972).

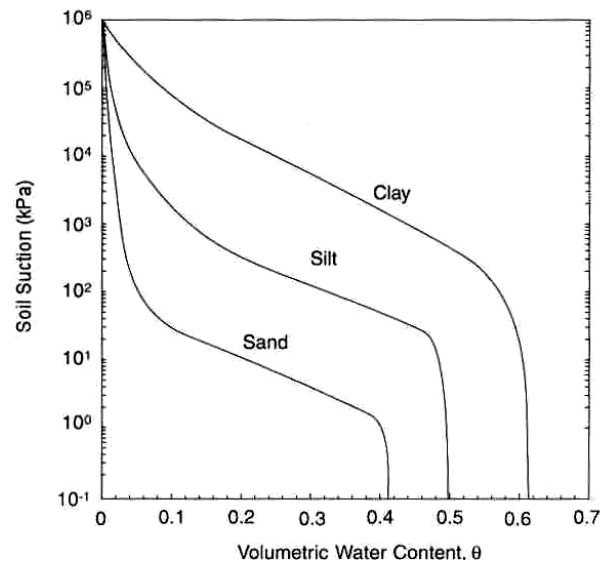


Figure 2.7. Representative SWCC for sand, silt and clay (Lu and Likos, 2004).

Two processes are described in the SWCC, adsorption (wetting), and desorption (drying). The curves for the wetting and drying do not usually coincide due to hysteretic behavior that occurs between the two processes. This behavior is attributed to the irregularity of pore shapes (Hanks, 1992) resulting in what is known as the “ink-bottle” effect. This concept is illustrated in Figure 2.8 that represents a capillary tube with an irregular radius (a small radius r , and a larger

radius R). During a wetting process (Figure 2.8a) the tube is initially empty and the height of the capillary rise is controlled by the smaller radius (r), and stops where the larger radius is encountered. The value of the matric suction here is: $\frac{2T_s}{r}$ (where T_s is the surface tension). On the other hand, during a drying process where the tube is initially filled with water, the height of the capillary rise might extend beyond the large radius (R), however, the matric suction at equilibrium in this case is still $\frac{2T_s}{r}$. Therefore, for the same value of matric suction, the wetting and drying processes have different water content values, causing the hysteretic behavior.

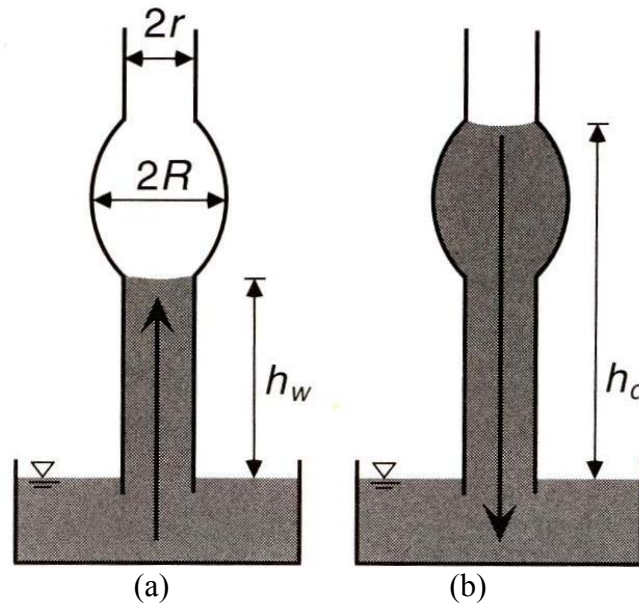


Figure 2.8. The ink bottle effect illustrated in a capillary tube model (Lu and Likos, 2004)

Due to this irregularity, losing the pore water in the drying process and filling the pores with water in the wetting process have completely different paths causing the hysteretic behavior. For this reason, a differentiation must be made between the wetting and drying characteristic curves. More water is generally retained by soil during a drying process than what is absorbed during the wetting process under the same value of suction. Knowing the SWCC of

an unsaturated soil helps in the assessment of many of its properties like the coefficient of permeability, shear strength, volumetric strain and the pore size distribution (Jian and Jian-lin, 2005).

2.1.6 Soil Suction Measurements Techniques

In typical unsaturated field conditions, the air pressure (u_a) is equal to the atmospheric pressure, and the pore water pressure is generally negative. In the laboratory, it is difficult to perform negative pore water pressure measurements without including the effect of cavitation. For this reason some indirect techniques are utilized to measure soil suction. The most popular laboratory technique is known as the axis translation technique developed by Hilf (1956). It is based on the assumption that the state of stress of a soil remains constant if the values of the net stress ($\sigma - u_a$) and suction ($u_a - u_w$) are not altered (Burland 1964). In this technique, the net stress and suction are controlled by applying a positive water pressure, and increasing the air pressure and the total stress by the same amount to achieve the desired net stress and matric suction. The reference or “axis” for matric suction is “translated” from the condition of atmospheric air pressure and negative water pressure to the condition of atmospheric water pressure and positive air pressure without changing the state of stress of the soil (Lu and Likos, 2004).

Axis translation is accomplished by separating the air and water phases of the soil specimen. This is achieved using high air entry (HAE) porous stones in the drainage system. When saturated, HAE materials prevent the seepage of air into the drainage system as long as the difference between the air pressure and water pressure is less than the air entry value. Macari and Hoyos (2001) and Hoyos and Macari (2001) developed a cubical true triaxial device where unsaturated soils can be tested under controlled suction using the axis translation technique for a wide range of stress paths. They conducted a series of drained true triaxial experiments on silty sand and concluded that matric suction has a significant influence on the stress-strain behavior of

unsaturated soils. It was found to exert a paramount influence on the size, position, and shape of the potential failure envelopes in the octahedral stress plane. This study also examined the influence of the loading rate on stress-strain behavior of the specimen (Figure 2.9). They stated that loading should be at a rate slow enough to ensure that the water menisci would not be disturbed resulting in the loss of matric suction.

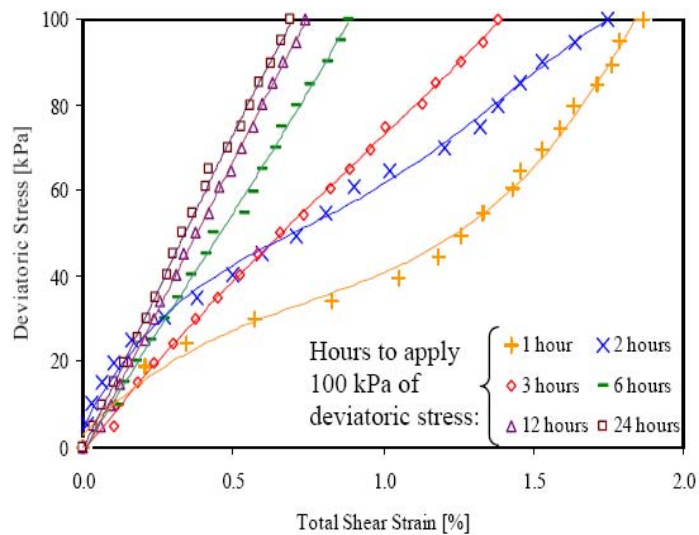


Figure 2.9. Effect of the loading rate on the stress-strain behavior of unsaturated soils. (After Hoyos and Macari, 2001)

2.2 Elastic Wave Propagation in Soils

Elastic waves can be defined as small mechanical perturbations that traverse particulate media without causing permanent effects or altering on-going processes (Santamarina et al. 2001). The physical interpretation of the measurements of these waves permits inferring important information about the particulate material and its processes (Choi et al. 2004). In an infinite elastic medium, there are two types of body waves: Compressional or primary waves (P-waves) where the particle motion and the wave propagation are in the same direction (Figure 2.10a), and shear or secondary waves (S-waves) where the particle motion is perpendicular to the direction of wave propagation (Figure 2.10b).

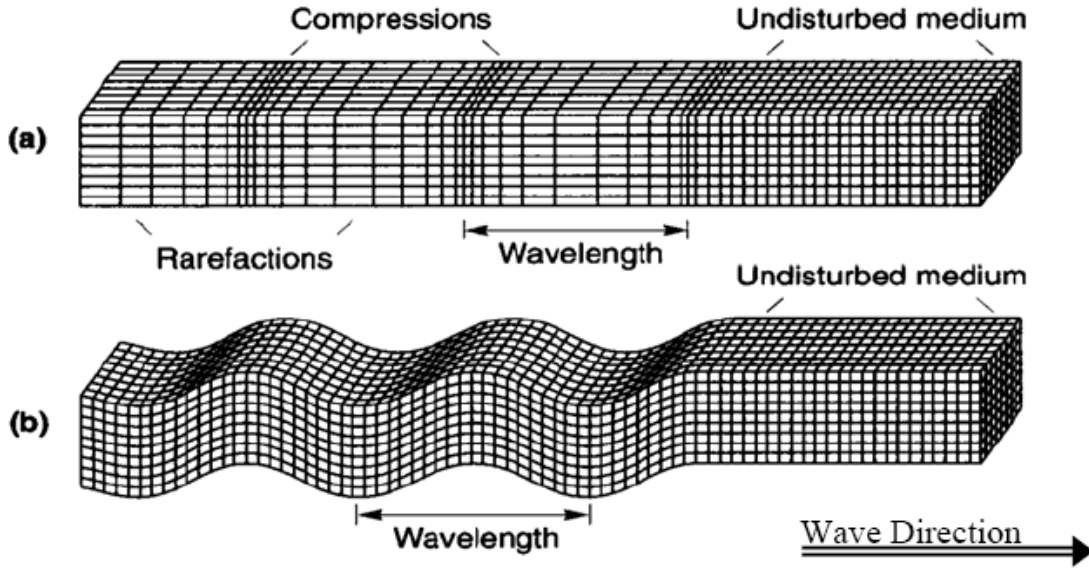


Figure 2.10. Types of body waves (a) P-wave (b) S-wave (Kramer 1995)

The propagation of elastic waves in geophysical studies involves strain levels that are lower than the threshold strain of the soil. The propagation velocity through a medium is proportional to its stiffness and inversely proportional to its density. The shear wave velocity (V_s in Equation 2.17) depends on the shear modulus of the soil (G_{soil}) which is only dependent on the skeleton shear stiffness and not influenced by the bulk stiffness of the pore fluid. For this reason, S-waves are preferred for the characterization of saturated soils (Santamarina et al. 2005).

$$V_s = \sqrt{\frac{G_{\text{soil}}}{\rho_{\text{soil}}}} \quad (2.17)$$

where ρ_{soil} is the mass density of the soil mass. G_{soil} is determined by the state of stress, degree of cementation, and by processes that alter the interparticle contacts such as capillary forces and electrical forces. On the other hand, the propagation velocity of P- waves (V_p in Equation 2.18) is proportional to the constraint modulus M_{soil} .

$$V_p = \sqrt{\frac{M_{\text{soil}}}{\rho_{\text{soil}}}} = \sqrt{\frac{B_{\text{soil}} + \frac{4}{3}G_{\text{soil}}}{\rho_{\text{soil}}}} \quad (2.18)$$

where B_{soil} is the bulk modulus of the soil. The bulk stiffness of the mineral grains B_g is much greater than the bulk stiffness of the granular skeleton B_{sk} . Furthermore the bulk stiffness of de-aired fluids B_{fl} (or B_w in the case of water) is also greater than the bulk stiffness of the skeleton. However, even minute quantities of air in the fluid phase drastically reduce the bulk modulus of fluid mixture (Richart et al. 1970). Detailed equations for the soils bulk modulus will be presented in Section 3.2.

When the soil is fully saturated, the bulk modulus is dominated by the bulk modulus of water which is very high, resulting in a very high propagation velocity. Hence, this kind of wave is not preferred when studying saturated soils. Figure 2.11 shows the increase in the wave velocity with increasing the effective stresses and cementation. It also shows the decrease in the wave velocity with increasing the degree of saturation which is caused by the decrease in capillary forces with the increase in the degree of saturation.

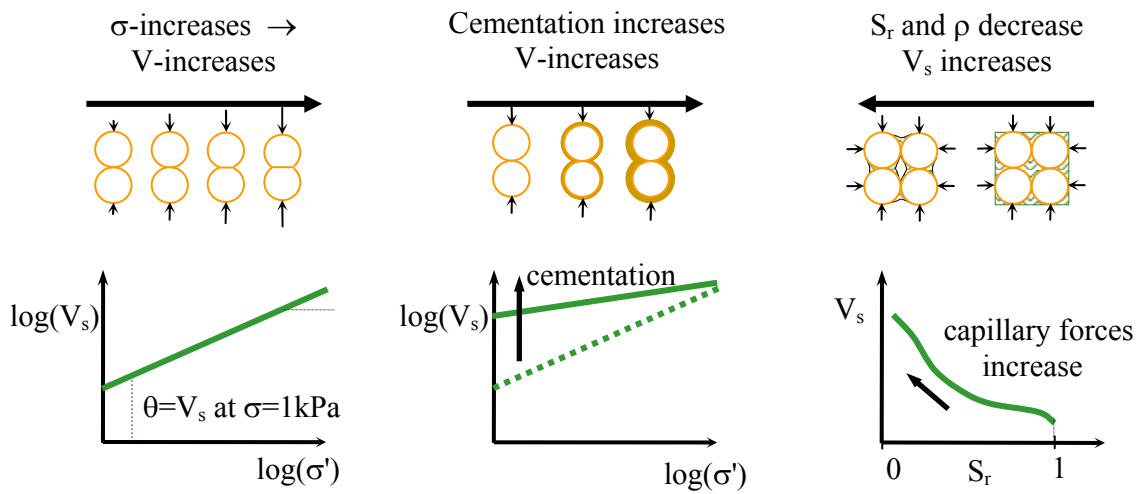


Figure 2.11. Effect of inter-particle forces on wave velocity (Fratta et al., 2001)

White (1983) presented a generalized relationship between P-wave velocity and effective isotropic stress σ' in a simple cubic packing as:

$$V_p = \left[\frac{3E_s^2}{8(1-\nu_s^2)^2} \right]^{\frac{1}{6}} \cdot \left[\frac{6}{\pi \cdot \rho_s} \right]^{\frac{1}{2}} \cdot \sigma'^{\frac{1}{6}} \quad (2.19)$$

where E_s , ν_s , ρ_s are the Young's modulus, Poisson's ratio, and mass density of the spherical particles. For random particle arrangements, Roessler (1979) and Stokoe (1991) presented a predictive empirical equation for the shear wave velocity in saturated or dry particulate materials as:

$$V_s = \theta \left(\frac{\sigma'_{\parallel} + \sigma'_{\perp}}{2p_r} \right)^{\beta} \quad (2.20)$$

where θ and β depend on particle type and agreement (Figure 2.12), σ'_{\parallel} and σ'_{\perp} are the stresses in the direction parallel and perpendicular to the direction of wave propagation, and $p_r = 1$ kPa is the reference pressure. Santamarina et al. (2005) extended this expression to unsaturated particulate materials as:

$$V_s \approx V_{s(\text{for } S_r=1.0)} \left[1 + \frac{(u_a - u_w) \cdot S_r}{0.75 \sigma'_v} \right] \quad (2.21)$$

where u_a and u_w are the pore air and pore water pressures respectively, S_r is the degree of saturation and σ'_v is the vertical effective stress. The ratio of the P-wave velocity to the S-wave velocity can be estimated using Poisson's ratio ν as in Equation 2.22. Figure 2.13 illustrates the P to S wave velocity ratios for ν range of 0 to 0.5.

$$\nu = \frac{0.5 \left(\frac{V_p}{V_s} \right)^2 - 1}{\left(\frac{V_p}{V_s} \right)^2 - 1} \quad (2.22)$$

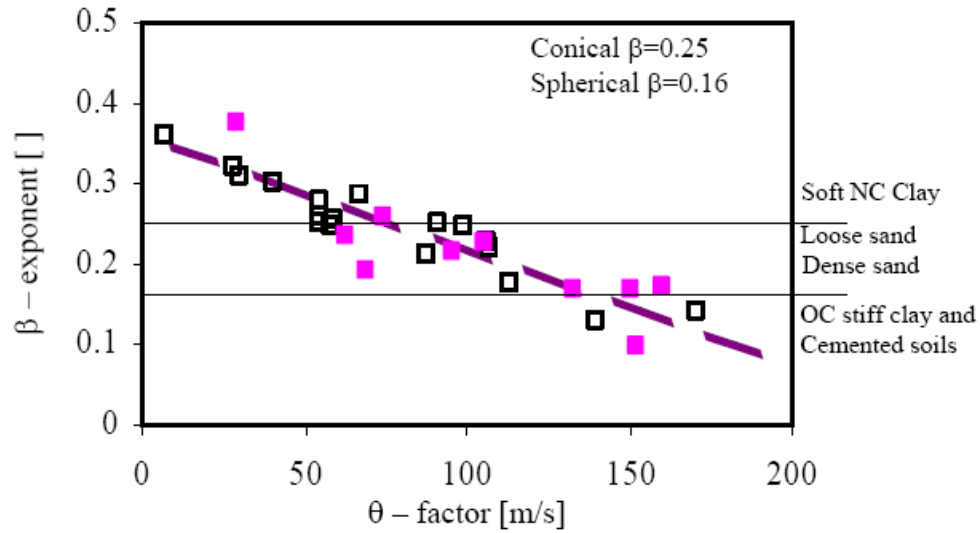


Figure 2.12. Relationship β -exponent and θ -factor for different types of soils (after Santamarina et al. 2001)

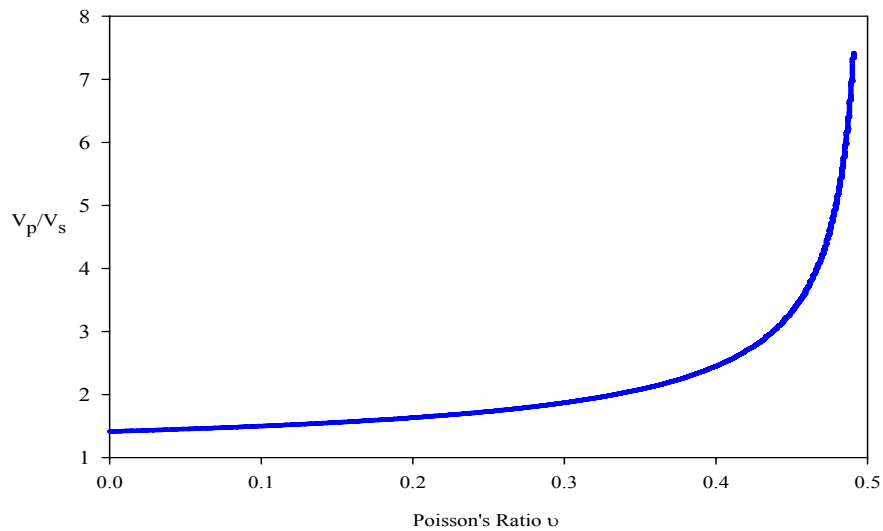


Figure 2.13. The ratio of P to S-wave velocities at different Poisson's ratio values

2.2.1 Measurement of Wave Velocity Using Bender Elements

Piezoelectricity is a phenomenon resulting from lack of crystal symmetry or from the electrically polar nature of crystals. When a mechanical load is applied to a piezo material, the lattice distorts the dipole moment of the crystal and a voltage is generated. The voltage output increases with crystal asymmetry. On the other hand, when a voltage is applied, the crystal deforms (Lee and Santamarina, 2005).

Bender elements are composed of two layers of thin piezoelectric sheets glued to opposite sides of a conductive metal shim (Figure 2.14). They can be used either as actuators or sensors. Piezo motors (actuators) convert voltage and charge to force and motion. When an AC voltage is applied to the bender element, one layer expands while the other contracts producing a bending typically in the order of hundreds to thousands of microns and a force from tens to hundreds of grams. Altering the current on the sides of element results in reversing the bending action, which generates a vibration and hence a mechanical wave. On the other hand, piezo generators (sensors) convert force and motion to voltage and charge. When a mechanical force causes a suitable polarized 2-layer element to bend, one layer is compressed and the other is stretched causing a charge to develop across each layer in an effort to counteract the imposed strains (Piezo systems, 2005).

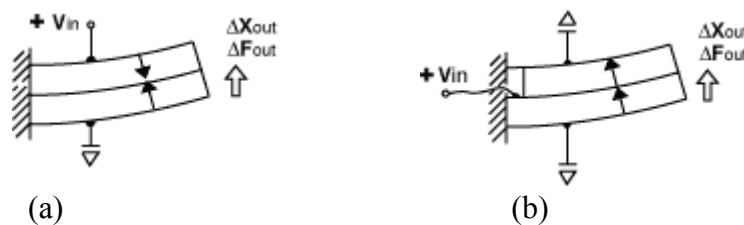


Figure 2.14. Schematic of a piezoelectric bender element. a) series connection, b) parallel connection (Piezo systems, 2005).

Bender elements can be configured for either series or parallel operation. In series configuration, the voltage is applied to all piezo layers at once, while in the parallel configuration the supply voltage is applied to each layer individually. Series configuration develops twice the voltage as the parallel, but provides only half the displacement for the same applied voltage. Therefore it is recommended to use parallel bender elements as sources and series bender elements as receivers.

Piezoelectric elements can be used in geotechnical testing to send and receive P and S-waves. The type of wave generated depends on mounting configuration of the bender element. The most common form is cantilever bender elements mounted on both ends of a triaxial or an Oedometer cell (Figure 2.15) to generate and receive S-waves (e.g. Blewwett et al. 1999, Cho and Santamarina 2001, Choi et al. 2004).

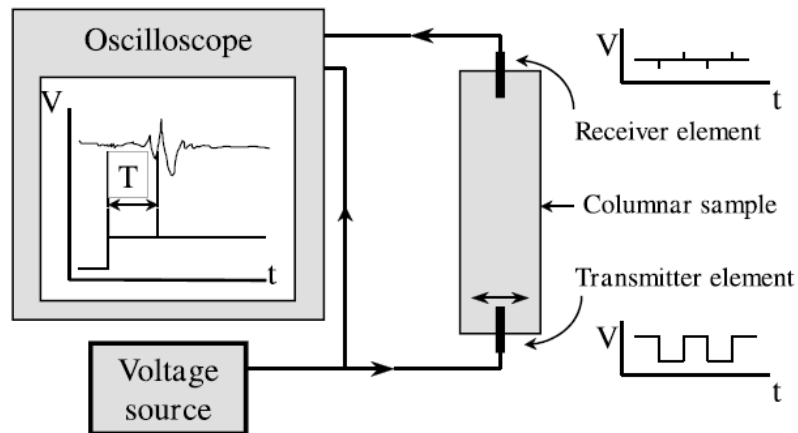


Figure 2.15. Schematic representation of shear velocity measurement technique (Blewwett et al. 1999)

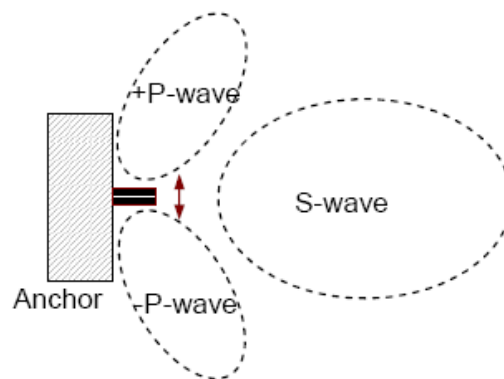


Figure 2.16. Schematic representation of the waves generated by a vibrating bender element (Lee and Santamarina, 2005).

Bender elements also generate two P-wave side lobes normal to their plane (Figure 2.16). Therefore, P-wave velocity can be measured, by placing sensors in the transverse direction of the specimen. P-waves can also be generated using other kinds of piezoelectric materials. Two-layer circular bending disks with constrained edges can be used (Figure 2.17). When a voltage is applied to the disk, a drum-like movement is generated, creating a P wave. Another disk can be used on the other side of the specimen to receive the waves.

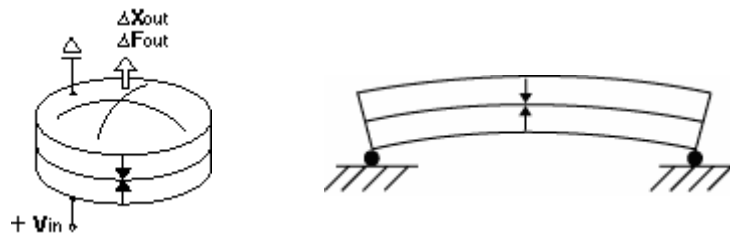


Figure 2.17. Schematic of a two layer piezoelectric disk

The voltage supplied to the source element and the one generated by the receiver are recorded using an oscilloscope. The time difference between the emitted and received signals can be then evaluated. Knowing the distance between the source and the receiver, the wave velocity can be calculated. Example source and receiver traces are depicted in Figure 2.18.

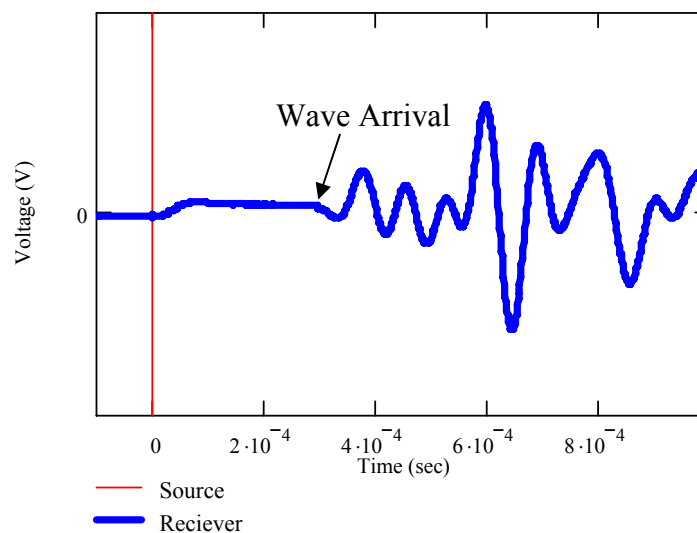


Figure 2.18. Example bender element traces

2.3 Electromagnetic Wave Parameters and Time Domain Reflectometry

2.3.1 Soils Electrical Properties

Dielectric permittivity is a measure of the polarization of a dielectric material when exposed to an external electric field (Mojid et al. 2003). It is a physical quantity that defines the materials ability to store electric charges and determines the phase velocity of electromagnetic radiation through the medium. The expression for the dielectric permittivity is defined as (Topp et al. 1980):

$$\kappa^* = \kappa' + j \left(\kappa'' + \left(\frac{\sigma_{dc}}{\omega \epsilon_0} \right) \right) \quad (2.23)$$

where κ^* is the complex dielectric permittivity, κ' and κ'' are the real and imaginary parts of the dielectric permittivity, respectively, σ_{dc} is the zero frequency conductivity, ω is the angular frequency, ϵ_0 is the free space permittivity (8.85×10^{-12} F/m), and the imaginary term j^2 is equal to -1. Davis and Annan (1977) indicated that in the 1 MHz to 1 GHz frequency range the real part of the dielectric permittivity does not appear to be strongly dependent on frequency. They also indicated that the dielectric loss κ'' was considerably less than κ' in this frequency range (e.g., Figure 2.19). The propagation velocity (V) of an electromagnetic wave in a transmission line of a known length is expressed as (Topp et al. 1980):

$$V = \frac{c}{\sqrt{\kappa' \frac{1 + \sqrt{1 + \tan^2 \delta}}{2}}} \quad (2.24)$$

where c is the velocity of an electromagnetic wave in free space ($3 * 10^8$ m/s) and $\tan(\delta)$ is the loss tangent defined as:

$$\tan \delta = \frac{\kappa'' + \left(\frac{\sigma_{dc}}{\omega \epsilon_0} \right)}{\kappa'} \quad (2.25)$$

In the case of most geotechnical applications, the loss tangent has been very small ($\ll 1$), then Equation 2.25 can be simplified to:

$$V = \frac{c}{\sqrt{k'}} \quad (2.26)$$

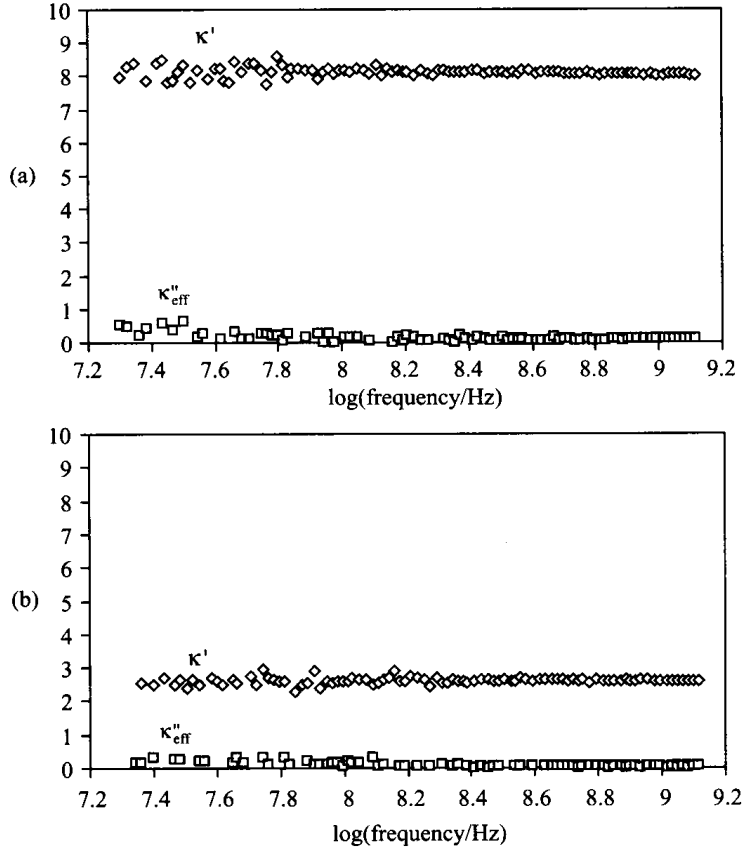


Figure 2.19. Dielectric permittivity spectra for solids: (a) a shale and (b) dry kaolinite (Santamarina et al., 2001)

When Equation 2.26 is valid (i.e. low loss condition), κ' is referred to as the apparent dielectric permittivity κ_a . From Equation 2.26, if the apparent dielectric permittivity of the medium is known, the velocity of the pulse can be determined, or if the velocity of the pulse is known, the dielectric permittivity can be calculated.

For a multiphase system such as soils, the bulk dielectric function depends on the volume fraction of soil solids, water, and air in the medium. The dielectric permittivity is typically 80 for liquid water (at 20° C), 3 to 8 for soil minerals, and unity for air. One can notice that the amount of water present in the soil controls the value of the bulk dielectric permittivity due to the high value of its dielectric permittivity. Furthermore the bulk dielectric permittivity is relatively insensitive to soil composition and texture due to the large disparity of dielectric permittivity values (Jones et al., 2001).

2.3.2 TDR - Principle of Operation

Time Domain Reflectometry (TDR) is a non-destructive method that uses electromagnetic pulses to measure the travel time along probes inserted in soils. Such travel time is then used to calculate the electromagnetic wave velocity and the dielectric permittivity and to estimate the volumetric water content which is defined as the ratio of the volume of water to the total volume. TDR has long been used in the telecommunications industry to detect faults and discontinuities in data transfer cables. It was first used by Topp et al. (1980) to measure soil water content.

The TDR system sends a high frequency pulse signal to a sensor through a coaxial cable (Figure 2.20). Reflections are recorded from any discontinuities in impedance along the path of the signal (Mojid, 1998). The trace of the reflected pulses is recorded and the travel time is measured. In geotechnical applications, the reflections are caused by the mismatch in the electromagnetic impedance between the cable and the soil along the waveguide. Knowing the time between the two reflections and the length of the probes, the travel velocity in the soil medium can be calculated. A typical TDR trace is shown in Figure 2.21 and the reflections as the pulse travels in the soil are illustrated.

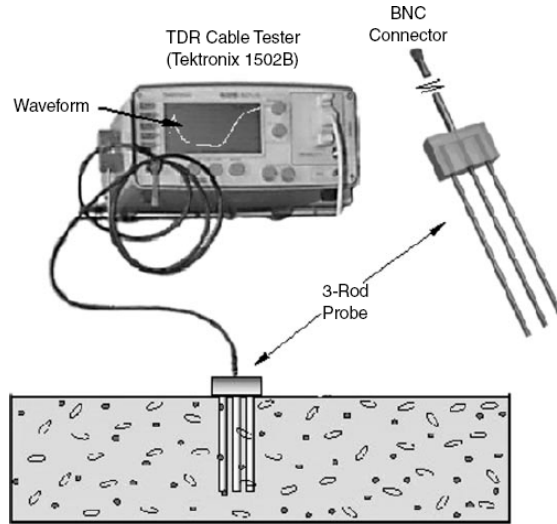


Figure 2.20. TDR System with probes vertically embedded in surface soil layer (Jones et al. 2001)

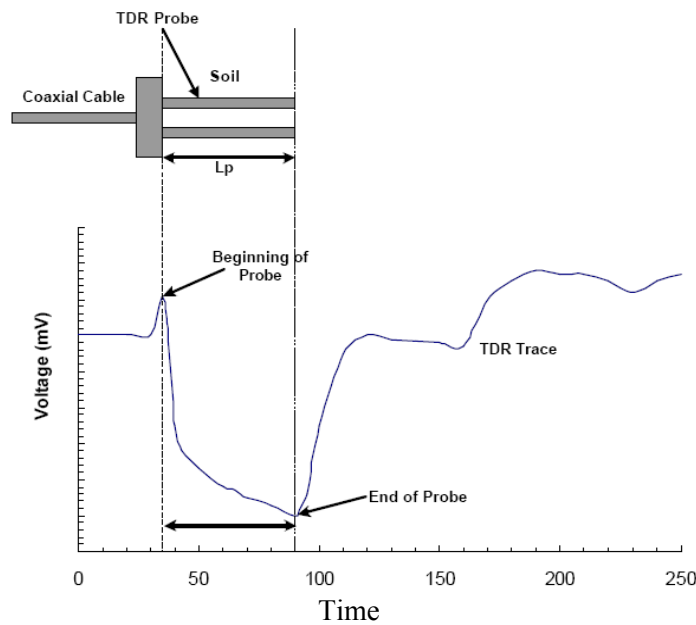


Figure 2.21. Typical TDR trace and reflection points. (after Benson and Bosscher 1999)

After analyzing the TDR trace, the time required for the electromagnetic pulse to travel to the end of the probe and the bulk dielectric permittivity is obtained using the following relationship:

$$\kappa_b = \left(\frac{c}{v} \right)^2 = \left(\frac{ct}{2L} \right)^2 \quad (2.27)$$

where t is the time required for the pulse to travel the length of the embedded probe and come back, and L is the length of the probe. As mentioned earlier, the value of the bulk dielectric permittivity is mainly controlled by the fraction of water in the medium. Several researchers have developed equations relating the soil dielectric permittivity to the volumetric water content (θ_v). Some common equations are listed in Table 2.2. It can be noticed that while some equations use the dielectric permittivity as the only parameter to calculate the volumetric water content others incorporate other parameters such as the porosity (n) and the soil bulk density ρ . Although good correlations were obtained using only the dielectric permittivity, adding these physical properties could yield better correlations since the volumetric water content is highly influenced by the porosity and bulk density. Figure 2.22 shows a comparison between these correlations. For the practical range of volumetric water content of 0 to 0.3, the values predicted by the different equations are very close to each other.

Table 2.2. Evaluation of volumetric water content using TDR measurements

Researcher	Equation
Topp et al. (1980)	$\theta_v = -5.3 \cdot 10^{-2} + 2.92 \cdot 10^{-2} \cdot \kappa - 5.5 \cdot 10^{-4} \cdot \kappa^2 + 4.3 \cdot 10^{-5} \cdot \kappa^3$
Wensink (1993)	$\theta_v = -1.294 + \sqrt{1.486 + 0.063 \cdot \kappa}$
Mixture equation ($\beta \approx 0.5$)	$\theta_v = \frac{\kappa^\beta - (1-n) \cdot \kappa_s^\beta - n \cdot \kappa_a^\beta}{\kappa_w^\beta - \kappa_a^\beta}$
Malicki et al. (1996)	$\theta_v = \frac{\sqrt{\kappa - 0.819} - 0.168 \cdot \rho - 0.159 \cdot \rho^2}{7.17 + 1.18 \cdot \rho}$

Sources: Topp et al. (1980); Benson and Bosscher (1999); Jones et al. (2001); Noborio (2001); Santamarina et al. (2005)

The electrical conductivity of soils can also be determined using waveforms of the TDR system. The Giese and Tiemann (1975) thin section approach gives the most accurate results for the soil electrical conductivity (Zegelin, 1989, Benson and Bosscher, 1999, Jones et al., 2001). The Giese-Tiemann's equation can be written as:

$$EC(S/m) = \frac{\varepsilon_0 c}{L} \frac{Z_0}{Z_c} \left(\frac{2V_0}{V_f} - 1 \right) \quad (2.28)$$

where Z_0 is the characteristic probe impedance and Z_c is the TDR cable tester output impedance (typically 50 Ω). V_0 is the incident pulse voltage and V_f is the return pulse voltage after multiple reflections have died out (Figure 2.23). The probe characteristic impedance is determined using a separate calibration procedure by immersing the probe in deionized water with known dielectric permittivity ε (Equation 2.29).

$$Z_0 = Z_c \sqrt{\varepsilon} \left(\frac{V_1}{2V_0 - V_1} \right) \quad (2.29)$$

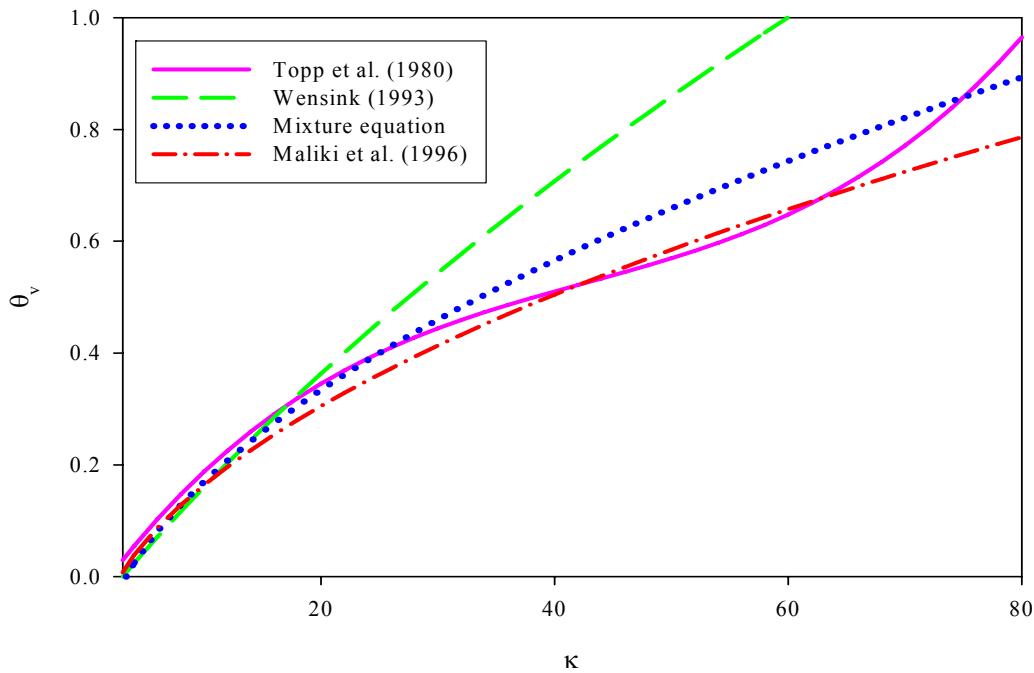


Figure 2.22. Comparison of different TDR equations.

Jones et al. (2001) listed the benefits of using TDR over other soil water content measurement methods as:

- Superior Accuracy ($\pm 2\%$ of θ_v).
- Minimal calibration requirements; usually no soil-specific calibration is needed.

- Lack of radiation hazard.
- Excellent spatial resolution.
- Measurements are simple to acquire.

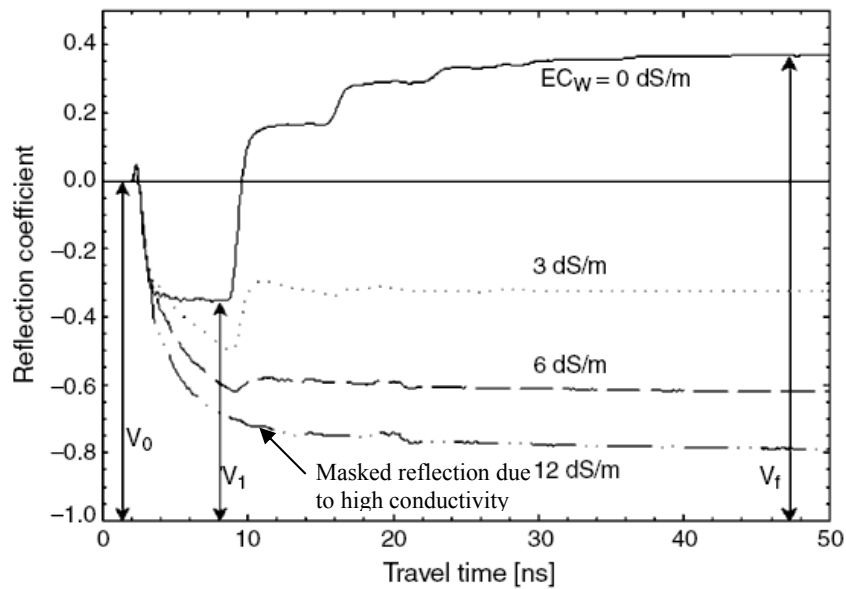


Figure 2.23. Determination of electrical conductivity from TDR traces (Jones et al. 2001).

2.3.3 The Effect of Soils Electrical Conductivity on TDR Measurements

Many conflicting results have been published in the literature to describe the effect of electrical conductivity on the estimated dielectric permittivity and volumetric water content and volumetric water content of the medium (Nadler et al., 1999). Some researchers stated that increasing the salinity of the soil solution can cause over-estimation of the water content (Dalton, 1992; Noborio et al., 1994; Wyseure et al., 1997). They argued that the increase in the electrical conductivity causes an increase in the imaginary part of the dielectric permittivity causing the apparent dielectric permittivity to be overestimated, therefore, the volumetric water content to be over-estimated. Wyseure et al. (1997) presented Equation 2.30 to estimate the increase in the

apparent dielectric permittivity (κ_a) caused by the increase in the soils electrical conductivity (σ). On the other hand, Robinson et al. (1999) presented Equation 2.31 to calculate κ_a .

$$\kappa_a = \varepsilon_r + 14.31\sigma \quad (2.30)$$

$$\sqrt{\kappa_a} = \sqrt{\varepsilon_r} + 1.72\sigma \quad (2.31)$$

where ε_r is the relative dielectric permittivity. Other researchers found no change in the estimated water content (Topp et al., 1980; Dalton and van Genuchten, 1986; Nadler et al., 1991) whereas some researchers found a decrease in the water content with the solution salinity (Herkelrath et al., 1991; Nielsen et al., 1995; Vogeler et al., 1996).

Increasing the conductivity of the medium causes the attenuation of the electromagnetic pulse as it passes through it. This situation results in the reduction of the voltage returned to the cable tester (V_2 in Figure 2.24). Dalton et al. (1984) conducted a study on ten different soil columns with different conductivity values and compared the waveforms to determine the effect of the soils electrical conductivity on TDR measurements. Figure 2.24 shows typical results of their study as reported by O'Connor et al. (1999).

The attenuation of the waveform at high conductivity (salinity) values results in masking the reflection at the end of the probes causing some difficulties in determining the travel time. Dalton et al. (1992) also showed that 8 dSm^{-1} is the value above which over-estimation of volumetric water content occurs. Wyseure et al. (1997) stated that increasing the electrical conductivity causes the loss of energy and hinders the propagation velocity of the electromagnetic pulse. Similar result was obtained by Hook et al. (2004). They concluded that the travel time of the pulse increases with increasing the pore water salinity and that the variability of travel time measurements increases, which leads to uncertainties in the measured water content. Hook et al. (2004) suggested that it would be impossible to form a general

correction for the effect of salinity on the measured volumetric water content under spatially variable field conditions even if the electrical conductivity is known. In their study, the effect of salinity on water content error were not apparent until the salinity of the sands under test reached 25 dSm^{-1} . The effect of electrical conductivity on the attenuation of the signal could be decreased using shorter probes with a lower practical limit of 10 to 15 cm (Nichol et al., 2002).

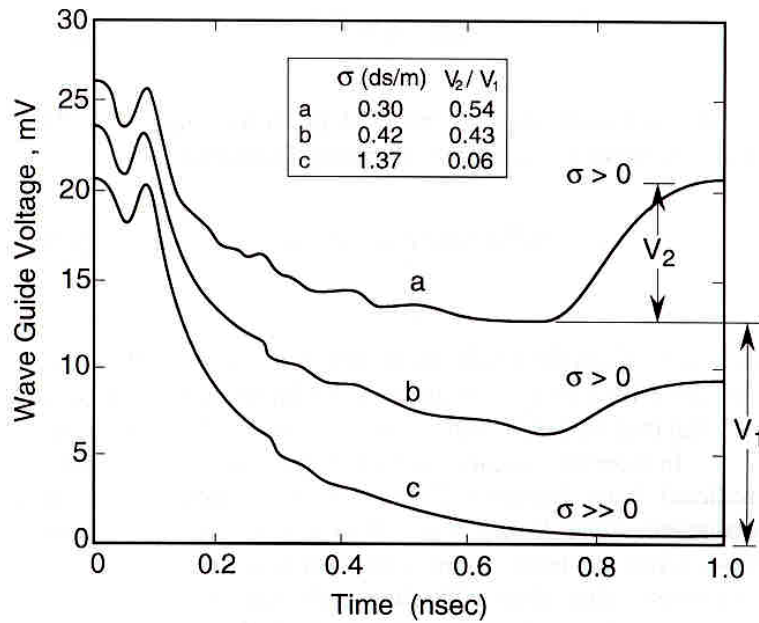


Figure 2.24. The attenuation of TDR waveforms with increasing the electrical conductivity

Mojid et al. (2003) investigated the effect of the electrical conductivity on the apparent dielectric permittivity (ϵ_a). They found that even at low values of EC the apparent dielectric permittivity is equal to the real part of the complex dielectric permittivity (ϵ_r), which no longer applies at high EC values. A considerable increase in the imaginary part (ϵ_i) and decrease in the real part of the dielectric permittivity was observed for electrolyte solutions. The resulting ϵ_a was much higher than ϵ_r . The value of ϵ_i increased sharply with decreasing the frequency of the applied electric field.

2.3.4 TDR Probe Configurations

As the electromagnetic pulse is emitted, it travels through two types of transmission lines: a coaxial cable where the electromagnetic field is constrained within the annulus between the inner and outer conductors, and parallel rods transmission line where such constraint no longer exists (O'Connor et al., 1999). The geometry of the rods in the TDR probe is of great importance as it controls the degree of disturbance in the soil, the sampling volume, and the resolution of the system. The aim of an optimal TDR probe design is to obtain a representative sampling volume with sufficient resolution and minimal soil disturbance. Unfortunately, these factors are not complimentary and compromises must be made to achieve the optimal design. For instance, increasing the spacing between the rods increases the sampling volume. However, it yields in a higher scatter of the electromagnetic field energy and a lower resolution. Moreover, increasing the number of the rods yields a better resolution but causes more disturbance.

The length of TDR rods ranges from few centimeters to 50 cm. The practical lower limit to maintain accuracy of the measurements is generally 10 to 15 cm (Heimovaara, 1993). Shorter probes can be also used successfully if the volumetric water content is consistently high or with higher frequency TDR systems (Persson and Wraith, 2002).

Many TDR probe designs have been suggested. Topp and Davis (1985) found that parallel rod transmission lines with center to center spacing of 50 mm are a practical compromise. They showed that the volume of the soil measured was a cylinder whose axis lies midway between the rods with a diameter of 1.4 times the spacing between the rods. Some designs have been proposed using two to seven rods. Another less common probe configuration is the parallel plate probe (Robinson and Friedman, 2000) which is typically used to generate a highly uniform electrical field between plates. Figure 2.25 shows some common TDR probe configurations.

Knowing the sampling volume of the TDR probe is of great importance. Ferre et al. (1998) defined the sampling volume as the region of porous material contributing to the TDR measurements. Changes in the medium outside this volume have no significant impact on the response of the instrument. Many researchers attempted to visualize the sampling volume to determine what portion of a soil is being measured (Robinson et al., 2003; Baker and Lascano, 1989; Zegelin et al., 1989; Knight, 1992; Peterson et al., 1995).

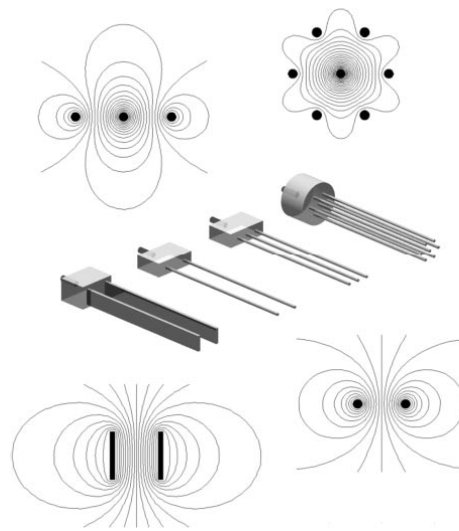


Figure 2.25. TDR Probe Configurations (Jones et al., 2001)

Ferre et al. (1998) compared the sampling volumes of balanced two and three rod probes. They found that at the same spacing, an increase in the rod diameter causes a marginal improvement in the uniformity of the distribution of the sensitivity within the sample area. The three-rod probe had a reduced sample area and more energy around the central probe. For this reason, Robinson et al. (2003) suggested that the two rod probes are preferable for field work.

Plates are used as an alternative to the rods in some TDR probes to produce a more uniform energy distribution. This is done to reduce the bias of the measurements for areas close to the rods (Robinson et al. 2003). However, using the plates causes more soil disturbance.

Robinson and Friedman (2000) conducted a comparison of four different probe geometries (two and three rod probes, and two and three plate probes) and stated that at least one wetting and drying cycle is needed for the probes to settle in the ground. Their results showed that improved measurement precision is obtained using plates instead of rods. However, the plates might be difficult to install and they are only suitable for non saline soils with low bulk EC due to a high geometric factor causing rapid waveform attenuation.

2.4. Methods of Field Evaluation of Mass Density and Water Content

2.4.1. Traditional Field Methods

Many methods have been traditionally used to determine the field density and moisture content. These methods include the sand cone and the rubber balloon. Alternative non destructive methods have also been developed including the nuclear density gauge.

The sand cone method (ASTM D1556-00) consists of a glass or a plastic jar with a metal cone attached to its top (Figure 2.26). The jar is filled with dry uniform Ottawa sand with known dry density (ρ_{sand} – by laboratory calibration). After taking the combined weight of the jar, the cone and the sand (W_1), a small hole is excavated in the area where the density is to be determined. The weight of the excavated soil is recorded (W_2). Then the jar is placed on top of the hole allowing the sand to flow out of the jar to fill the hole. The weight of the jar, the cone and the remaining sand is then determined (W_3). The weight of the sand that filled the cone and the hole (W_4) is calculated as the difference between W_1 and W_3 . Knowing the weight of the sand required to fill the cone (W_5), the weight of the sand that filled the hole (W_6) can be calculated as the difference between W_4 and W_5 . The volume of the hole can be calculated as:

$$V_{\text{hole}} = \frac{W_6}{\rho_{\text{sand}}} \quad (2.32)$$

Knowing the weight of the excavated soil, and the volume of the hole, the field density can be calculated as:

$$\rho_{\text{field}} = \frac{W_2}{V_{\text{hole}}} \quad (2.33)$$

The rubber balloon method (ASTM D2167-94 – Figure 2.27) is similar to the sand cone method, where a hole is excavated and the weight of the soil is recorded. The volume of the hole is determined by inserting a rubber balloon in the excavated hole and filling the balloon with water from a calibrated vessel. The balloon takes the shape of the hole as it is filled with water and the volume can be directly read from the vessel.

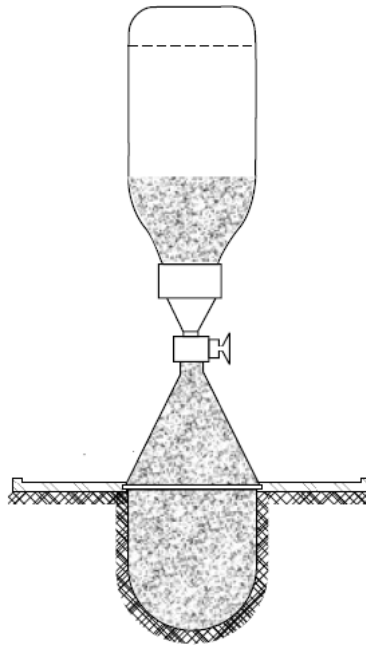


Figure 2.26. Schematic of the sand cone test (Multiquip, 2004)

Although these two methods are widely accepted for the field density, they are time consuming and destructive in nature. Furthermore, some inaccuracies may arise when using these methods such as poor laboratory calibration of the density of the Ottawa sand, the volume of the cone, and the loss of some of the soil excavated from the hole could result in the poor

estimation of the field density. Moreover, the cone is positioned on a perfectly flat surface during the laboratory calibration, but this is hardly ever the case in the field. The roughness of the soil surface in the field could result in some error in the estimated field density. Also, it should be noted that it is not possible to ensure that the density of the Ottawa sand poured in the hole is exactly the same as the calibration density. On the other hand, it is difficult to get the water balloon to take the exact shape of the excavated hole, resulting in some uncertainty in the volume reading.



Figure 2.27. The water balloon apparatus.

The nuclear density gauge (ASTM D6758-02- Figure 2.28) is a fast and fairly accurate non destructive method to measure the field density and moisture content. Gamma rays are emitted from a radioactive isotope source (Cesium 137) at the soil surface (backscatter) or from a probe inserted into the soil (direct transmission - Multiquip, 2004). After passing through the soil, a detector collects the Gamma rays and determines the soil density by examining the absorbance of the radiation through the medium. Due to its radioactive nature, the nuclear density gauge requires special safety procedures and certifications to transport and operate.

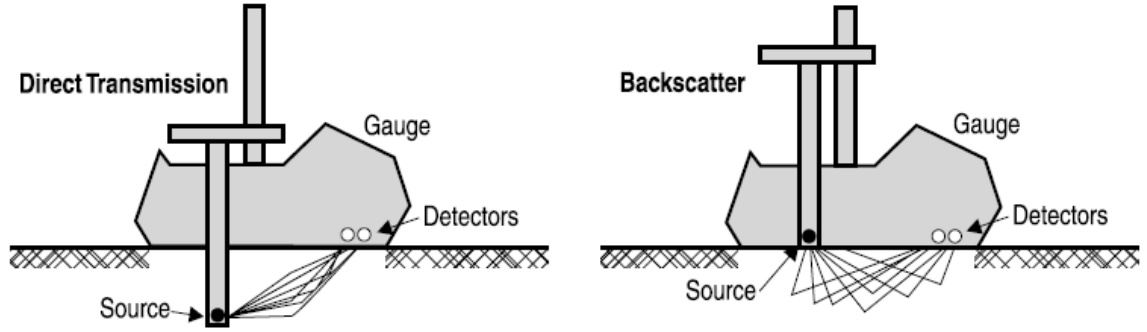


Figure 2.28. Schematic of the nuclear density test (Multiquip, 2004)

2.4.2 The Purdue TDR Method

Drnevich and his co-workers developed a methodology to determine the field density and moisture content using TDR measurements (The Purdue Two Step Method - Siddiqui and Drnevich 1995, Drnevich et al. 2001, Drnevich et al. 2003, ASTM D6780-02). Two measurements are required: a field measurement and a compaction mold measurement. Assuming a similar moisture content in the mold and in the field, the density and moisture content in the field can be calculated. Drnevich et al. (2001) suggested using Equation 2.34 to calculate the water content from the soils apparent dielectric constant.

$$w = \frac{1}{b} \left[\frac{\rho_w}{\rho_d} \sqrt{\kappa} - a \right] \quad (2.34)$$

where ρ_w is the density of water, ρ_d is the soils dry density, κ is the soils bulk dielectric permittivity measured by TDR, a and b are calibration constants. Knowing the density of the soil in the compaction mold, Equation 2.34 can be rewritten as:

$$w_{\text{mold}} = \frac{\sqrt{\kappa_c} - a \frac{\rho_c}{\rho_w}}{b \frac{\rho_c}{\rho_w} - \sqrt{\kappa_c}} \quad (2.35)$$

where κ_c is the apparent dielectric permittivity in the compaction mold and ρ_c is the density of the soil in the mold. Assuming the water content to be the same in the field and the compaction mold, the field dry density can be calculated as:

$$\rho_{d,f} = \frac{\sqrt{\kappa_f}}{\sqrt{\kappa_c}} \cdot \frac{\rho_c}{1+w} \quad (2.36)$$

This procedure uses two soil dependent calibration constants (a and b). They are experimentally determined for different soils through a series of compaction tests. Drnevich et al. (2003) states that based on tests on a variety of ordinary soils the value of a is close to one and b is approximately equal to 9.

Although this method gives quite accurate results, it still requires a field calibration which includes excavating the soil and preparing a compaction mold and repeating the procedure for different soils and different moisture contents in the field. Furthermore, a series of laboratory compaction tests are required to determine the values of “a” and “b” which makes this method time consuming and not very feasible when rapid field density and moisture content measurements are sought.

Yu and Drnevich (2004) developed an alternative method to eliminate the need for the compaction mold measurement, it is known as the Purdue One Step Method. It uses the electrical conductivity that can be calculated from the TDR traces along with the bulk dielectric permittivity to estimate the water content and dry density (Equations 2.37 and 2.38).

$$\rho_d = \frac{d \cdot \sqrt{\kappa_a} - b \cdot \sqrt{EC_b}}{a \cdot d - c \cdot b} \rho_w \quad (2.37)$$

$$w = \frac{c \cdot \sqrt{\kappa_a} - a \cdot \sqrt{EC_b}}{b \cdot \sqrt{EC_b} - d \cdot \sqrt{\kappa_a}} \quad (2.38)$$

where EC_b is the bulk electrical conductivity. a , b , c and d are calibration coefficients. This method still requires a laboratory calibration including a series of compaction tests to determine the six coefficients for different soils.

CHAPTER 3

COMBINED TDR AND ELASTIC WAVE VELOCITY MEASUREMENTS TO DETERMINE IN SITU DENSITY AND MOISTURE CONTENT

3.1 Introduction

In situ density and water content of soils are two essential properties in the design and evaluation of many earthwork projects. They are used as criteria for acceptance and quality control of most compaction projects including pavement systems, embankments, and clay liners. Therefore, it is imperative to have robust, accurate, and reliable methods to ensure the proper evaluation of these quality control parameters. Conventional and well established practices for the evaluation of the soil density and water content include the use of simple techniques such as the sand cone and water balloon methods which involve removing the soil for the evaluation of water content (ASTM D1556-00; ASTM D2167-94). These methods are widely accepted; however, they are time consuming and require the collection of specimens which makes the process destructive. A very well known alternative is the use of the nuclear density gauge (ASTM D3017-04) to determine both the in situ density and water content. This device allows for the rapid determination of these parameters by monitoring the attenuation of the energy generated by a radiation source. The radioactive nature of the device requires special transportation, operation safety procedures, and personnel certification.

Due to limitations in the time required to collect large number of specimens and the safety-related issues of these devices, both researchers and research & development companies have been trying to develop new non-nuclear technologies for rapid evaluation of in situ water content and mass density. There were many attempts toward the development of alternative techniques with different levels of success. For example, the GeoGauge (ASTM D6758-02)

measures the small-strain stiffness of soils at depths one to two times the depth of its base. It can be used as an indirect measure of the quality of compacted materials (Alshibli et al. 2005; Sawangsurinya et al. 2003 and 2004). Other devices include the Pavement Quality Indicator (PQI) and Soil Quality Indicator (SQI – Transtech, Inc.) which has been developed to evaluate the quality of asphalt pavements and compacted soils (including mass density) and the Electrical Density Gauge (Electrical Density Gauge, LLC). The SQI uses electrical impedance spectroscopy and a calibrated artificial neural network algorithms to back-calculate the soil density and moisture content. The Electrical Density Gauge uses a single frequency complex impedance of soil to correlate for density and water content in bases and sub-bases. These methods are still under performance evaluation.

A technique that is gaining acceptance and has been standardized is the Purdue TDR method which was developed by V. Drnevich and co-workers (Siddiqui and Drnevich 1995; Lin et al. 2000; ASTM D6780-02). It utilizes data collected using the TDR technique for soil compacted in a mold and in situ at the same water content to estimate the in situ moisture and dry density. The procedure involves driving four spikes into the soil surface using a template (Drnevich et al. 2003; Yu and Drnevich 2004). Then, the multiple-rod-head-probe TDR system is placed on the top of the spikes to measure the electromagnetic wave properties. The measurement procedure also includes extracting a soil specimen, placing it in a compaction mold, and measuring electromagnetic wave properties as a way to calibrate the measurements. Based on the two sets of measurements, the water content and dry density are calculated. Although promising results were obtained using this technique, the process of calibration and preparing compaction molds in the field makes it time consuming and not practical when rapid field measurements are sought. Moreover, the method uses four probes driven into the soil to

perform the TDR measurements that can cause disturbance to the sampled volume and alter the density of in situ soils.

In this chapter, an alternative non-destructive, non-nuclear method for determining the in situ density and moisture content of soils is presented. It combines the soils dielectric permittivity data acquired using TDR with P-wave velocity measurements. A model that correlates the P-wave velocity through soils to their volumetric water content and degree of saturation was used to back-calculate the density and moisture content. An experimental program was conducted to validate the proposed methodology and assess its range of applicability. The following sub-Sections provide a detailed description of the proposed methodology, model derivation, numerical validation and the results of the experimental program along with an analysis of the sources of error and limitation of the proposed methodology.

3.2 Model Derivation

In this Section, a model for the P-wave travel velocity in soils is derived. Then TDR and P-wave velocity measurements will be used to calibrate the model to find the parameters necessary to calculate the density and moisture content of the soil. First, some basic definitions of the soil phase relations are introduced (Figure 3.1). The soil density ρ is defined as:

$$\rho = (1 - n) \cdot \rho_s + n \cdot S_r \cdot \rho_w + n(1 - S_r) \cdot \rho_a \quad (3.1)$$

where ρ_s , ρ_w and ρ_a are the density of the solids, water, and air phases; and n , S_r and θ_v are the porosity, degree of saturation and volumetric water content, respectively. Knowing that the density of air is much smaller than the density of solid and water phases, Equation 3.1 may be reduced to:

$$\rho = (1 - n) \cdot \rho_s + n \cdot S_r \cdot \rho_w = (1 - n) \cdot \rho_s + \theta_v \cdot \rho_w \quad (3.2)$$

Porosity (n), degree of saturation S_r (or the volumetric water content $\theta_v = nS_r$) and ρ_s (or G_s) are needed to calculate the soil density. G_s ($=\rho_s/\rho_w$) is the specific gravity of solids and varies in a narrow range of values (from 2.6 to 2.8 for most soil minerals). The mass density of water is assumed to be known. Then, any methodology for the evaluation of mass density and water content should measure at least two parameters (e.g., n and S_r assuming that G_s is known or its value is estimated).

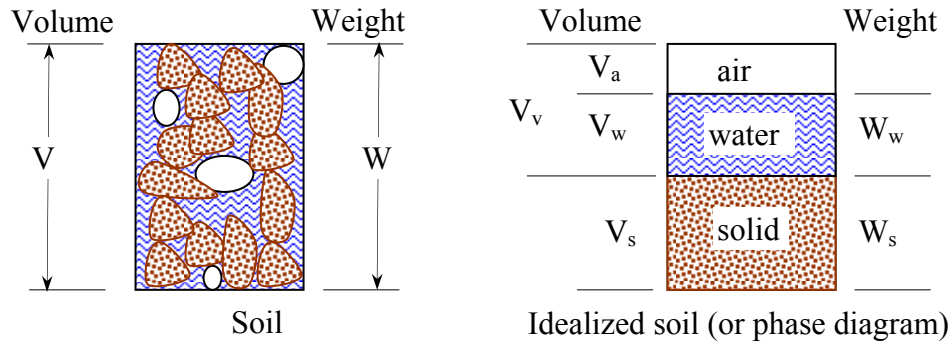


Figure 3.1. Phase diagram: Definitions

As presented in Section 2.2, the P-wave velocity in soils is expressed as:

$$V_p = \sqrt{\frac{M_{\text{Soil}}}{\rho_{\text{Soil}}}} = \sqrt{\frac{B_{\text{Soil}} + \frac{4}{3}G_{\text{Soil}}}{\rho_{\text{Soil}}}} \quad (3.3)$$

The soil bulk modulus B_{soil} is a function of the bulk modulus of the different soil components (i.e., soil minerals, pore fluid and the stiffness of the skeleton) and their relative quantities. Therefore, in order to quantify the overall soil stiffness, every component will be examined separately, and then the different components will be added together to form an expression for B_{soil} . The pore fluid consists of water (with a stiffness value $B_w = 2.19 \times 10^9$ Pa), and air (with a stiffness value $B_a = 1.42 \times 10^5$ Pa). The pore fluid stiffness can be calculated as a function of air and water stiffnesses in addition to the degree of saturation (S_r) as:

$$B_{fl} = \left(\frac{S_r}{B_w} + \frac{1-S_r}{B_a} \right)^{-1} \quad (3.4)$$

The soil fluid suspension stiffness is defined as the stiffness of the solid particles and fluid mixture (without the contribution of the solid particle contacts). It is a function of the stiffness of solid mineral grains ($B_g = 3.5 \times 10^9$ Pa), the fluid stiffness B_{fl} and the porosity (n). B_{sus} can be expressed as:

$$B_{sus} = \left(\frac{n}{B_{fl}} + \frac{1-n}{B_g} \right)^{-1} \quad (3.5)$$

The stiffness of the skeleton B_{sk} depends on the effective stress (Roessler 1979; Stokoe and Lee 1991; Vinale *et al.* 2001; Fratta and Santamarina 2002), capillary forces (Wu *et al.* 1984; Qian *et al.* 1993; Cho and Santamarina 2001), and cementation (Fernandez and Santamarina 2001). It is related to the soil shear modulus according to the following equation:

$$B_{sk} = G_{soil} \frac{2}{3} \frac{1+\nu}{1-2\nu} \quad (3.6)$$

where ν is Poisson's ratio (0.10 to 0.25 for elastic wave propagation in soils). The soil shear modulus G_{soil} depends not only on the state of stress, but also on the capillary forces (i.e., degree of saturation) and the degree of cementation. As the proposed methodology is to be applied in the top 15 cm of freshly compacted soils, the influence of effective vertical stresses and cementation are insignificant with respect to the capillary forces (Cho and Santamarina 2001). Therefore the soil shear modulus will be expressed as a function of the degree of saturation. The challenge here is to find an equation that represents this relationship considering that two competitive effects take that place as the degree of saturation changes. Matric suction decreases and the cross-section of the water phase between soil particles increases as the degree of

saturation increases (Lu and Likos 2004). Equation 3.7 will be used to express the shear stiffness:

$$G_{\text{soil}}(S_r) = G_o \cdot \left(S_r^{\frac{m}{1-m}} - 1 \right)^{\frac{1}{m}} \quad (3.7)$$

where m is an experimentally-determined parameter that is related to the pore size distribution of the soil and G_o is the soil shear modulus at degree of saturation $S_r = 2^{(1-m)/m}$. This equation follows a similar relationship as van Genuchten's relationship for soil-water retention curves (van Genuchten, 1980) as it depends on the type of soils (Wu et al., 1984; Qian et al., 1993; Cho and Santamarina, 2001). Typical normalized shear modulus $\left(\frac{G_{\text{soil}}}{G_o} \right)$ curves at different degrees of saturation and “ m ” values are illustrated in Figure 3.2. The semi-empirical relationship defined by Equation 3.7 combines both the effect of reducing suction when saturation increases while the capillary cross sectional area increases. At 100% saturation or zero suction, the soil shear modulus is null. Finally, the stiffness of the soil B_{soil} includes the stiffnesses of the suspension and the skeleton:

$$B_{\text{soil}} = B_{\text{sus}} + B_{\text{sk}} \quad (3.8)$$

Substituting Equations 3.4 to 3.8 into Equation 3.3, the P-wave velocity of a shallow freshly compacted soil may be expressed as (Fratta et al., 2005):

$$V_p = \sqrt{\frac{\frac{1}{n \left(\frac{S_r}{B_w} + \frac{1-S_r}{B_a} \right) + \frac{1-n}{B_g}} + \frac{7}{3} G_o \left(S_r^{\frac{m}{1-m}} - 1 \right)^{\frac{1}{m}}}{\rho_w \cdot G_s \cdot (1-n) + \rho_w \cdot n \cdot S_r}} \quad (3.9)$$

$$V_p = \sqrt{\frac{\frac{1}{\frac{\theta_v}{B_w} + \frac{n - \theta_v}{B_a} + \frac{1 - n}{B_g}} + \frac{7}{3} G_o \left(\left(\frac{\theta_v}{n} \right)^{\frac{m}{1-m}} - 1 \right)^{\frac{1}{m}}}{\rho_w \cdot G_s \cdot (1 - n) + \rho_w \cdot \theta_v}} \quad (3.9a)$$

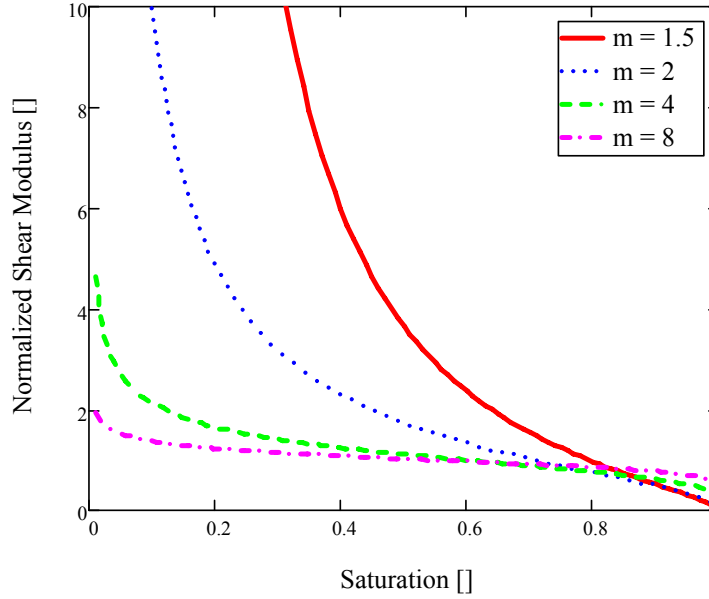


Figure 3.2. Relationship between normalized skeleton shear stiffness and degree of saturation for different m values

Equation 3.9 predicts the variation of the P-wave velocity as the skeleton shear stiffness changes with the degree of saturation. Experimental studies have shown that such behavior depends on the grain size distribution and texture of the grains (see for example Wu et al., 1984; Qian et al., 1993; Cho and Santamarina, 2001). Figure 3.3 illustrates how different model parameters, namely m and G_o , can be used to fit materials that might exhibit different rate of changes and magnitudes in the response. Note that as the volumetric water content reaches the numerical value of the porosity, the P-wave velocity rapidly increases. This is a common observation in both soils and weathered rocks.

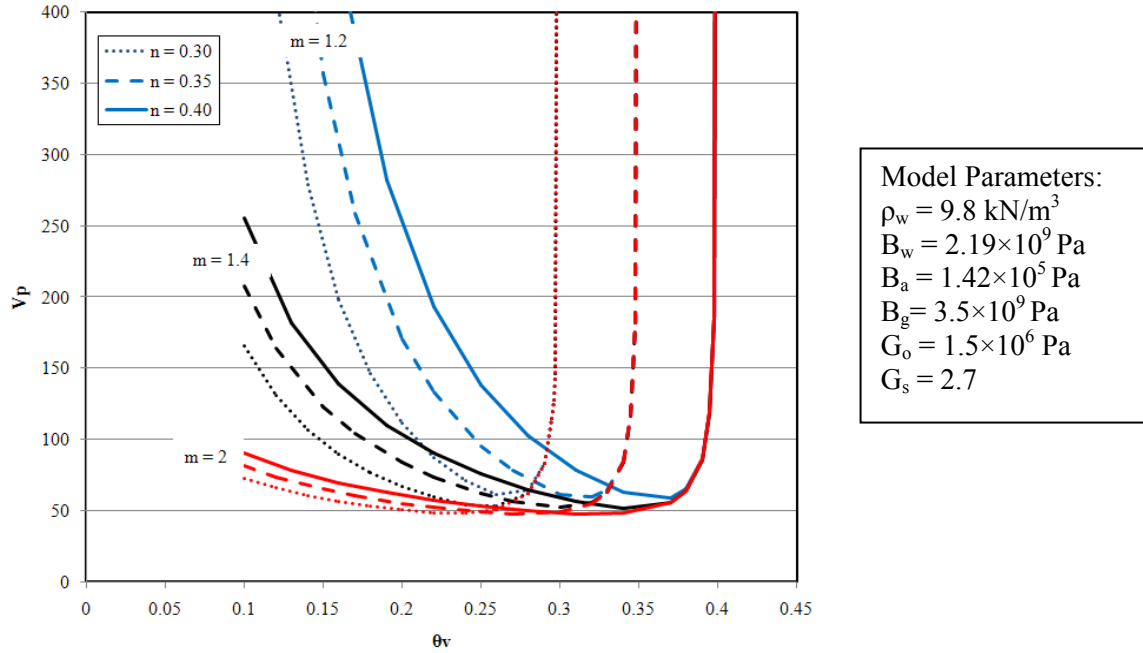


Figure 3.3. Modeled P-wave velocity profiles for different values of the parameter m and porosity.

Parameters n , S_r (or θ_v), G_o and m in Equation 2 should be determined by taking consecutive measurements of the P-wave velocity in the soil while increasing its moisture content. The following verification procedure is proposed to validate the proposed technique:

- Compact a soil specimen in the compaction mold
- Determine the weight and volume of the compacted specimen and calculate its bulk density
- Determine the compacted soil gravimetric water content by oven drying method
- Carefully drill a water injection hole at the center of the mold, care should be taken to minimize the disturbance of the surrounding soil
- Insert the TDR probes into the soil and measure the initial volumetric water content.
- Determine the initial P-wave velocity using surface accelerometers

- Increase the water content in the soil by injecting water into the hole
- Measure the volumetric water content and the P-wave velocity
- Repeat the measurements with increasing water content. About four points are required to acquire a solution for the inversion procedure

After taking multiple measurements, the parameters can be determined using a minimum error inversion procedure that includes fitting the wave velocity model to the measured points. The main assumption of this method is that the porosity of the soil is constant during the process. After determining n and S_r , Equation 3.2 can be used to calculate the density of the soil. The gravimetric water content can be calculated using Equation 3.10 (e is the void ratio):

$$w = \frac{S_r e}{G_s} \quad (3.10)$$

3.3 Inversion Procedure and Numerical Validation

3.3.1 Description of the Inversion Procedure

The model described by Equation 3.9 relates the P-wave velocity to the soil volumetric water content. The next step is to use the experimental data obtained by the proposed methodology to calculate the parameters necessary to calculate the soils density and moisture content. This was achieved by a minimum error inversion procedure. Such procedure considers a range of values for n , G_0 and m and uses the proposed semi-empirical model to calculate the P-wave velocity for every combination of the three variables within the defined ranges. The calculated values are then compared to the P-wave velocity values obtained experimentally. The error value is defined as the difference between measured and calculated P-wave velocity values:

$$\text{Error} = \frac{1}{N} \sqrt{\sum_{i=1}^N \left(V_{P_i}^{\langle \text{measured} \rangle} - V_{P_i}^{\langle \text{calculated} \rangle} \right)^2} \quad (3.11)$$

where $V_{p_i}^{<\text{measured}>}$ is the experimentally measured P-wave velocity, $V_{p_i}^{<\text{calculated}>}$ is the P-wave velocity calculated using the proposed model and N is the number of data points. The combination of model variables that yields the minimum error is then selected as the solution for the inversion problem. Some constraints are also used to assist in the convergence of the inversion procedure; that is, the values of porosity “ n ” should always be equal or larger than the highest measured volumetric water content because the volumetric water content becomes equal to the porosity when the soil is fully saturated. Moreover, the value of parameter m should be more than unity to ensure the inverse relationship between V_p and θ_v .

An algorithm was written using Matlab to perform the inversion, back-calculate the model parameters and use them to calculate the soils density and water content. For a set of ranges of m , G_o , and n values, an array of V_p values was calculated using Equation 3.9. The error value was then determined as the difference between the calculated and the experimentally measured P-wave velocity values. For each combination, the determined error value was stored in a three dimensional data array, where the x , y and z axes represent m , G_o , and n values, respectively. The algorithm then compares the error values and selects the combination of model parameters that yield a minimum error. The inverted porosity value can then be used to calculate density and water content.

3.3.2 Numerical Validation of the Inversion Procedure

A numerical validation was performed to assess the proper operation of the inversion procedure. The proposed model was used to generate a synthetic data set of P-wave velocities corresponding to assumed values of θ_v , n , G_o and m . White random noise was added to velocity values ($\pm 5\%$) and to the volumetric water content ($\pm 2\%$). These errors are commonly accepted in the measurement of wave velocity and volumetric water content (Hagedoorn 1964; Noborio 2001). Table 3.1 presents the synthetic values with and without the random noise. Synthetic data

were then used in the inversion procedure and the calculated n , G_o and m values were compared with the original values used to generate the data. The values of the assumed and inverted porosity and water content values are also listed in Table 3.1. Figure 3.4 illustrates the synthetic P-wave velocity versus volumetric water content as the soil is inundated with and without the random noise. A comparison of the assumed and inverted values of porosity, water content, density and dry density values is illustrated in Figure 3.5.

Table 3.1. Synthetic values, noisy data and inverted parameters

#	Specific gravity	Assumed Porosity	Assumed water content	Initial θ_v	Initial noisy θ_v	Initial noisy V_p	Inverted porosity	Inverted water content
1	2.6	0.35	0.10	0.169	0.166	107.0 m/s	0.346	0.094
2	2.6	0.40	0.14	0.218	0.215	128.9 m/s	0.388	0.130
3	2.7	0.37	0.12	0.204	0.204	191.4 m/s	0.374	0.121
4	2.7	0.42	0.16	0.251	0.252	144.4 m/s	0.414	0.159
5	2.8	0.40	0.14	0.235	0.231	81.3 m/s	0.398	0.142
6	2.8	0.45	0.18	0.277	0.280	145.8 m/s	0.451	0.189

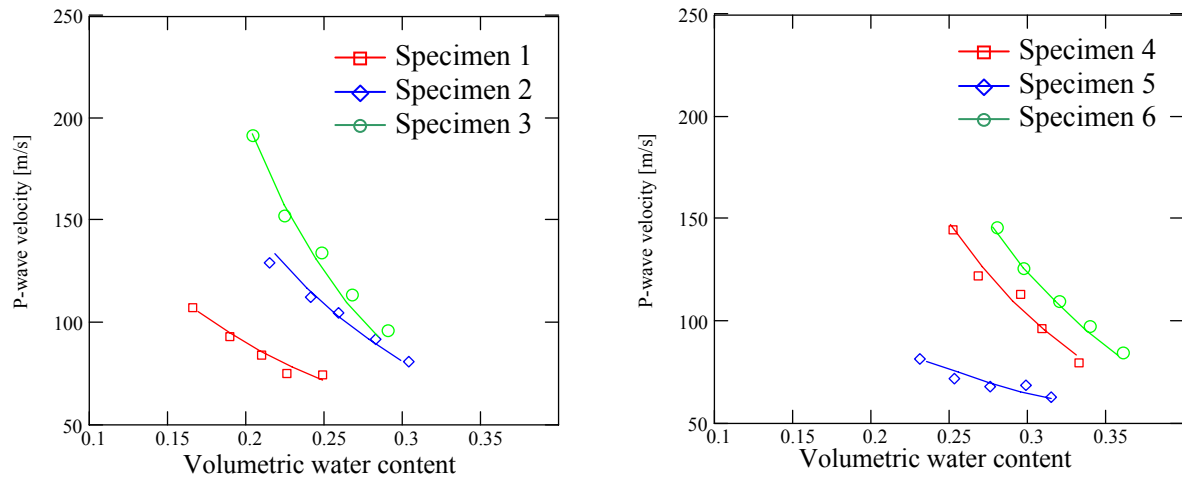


Figure 3.4. Modeled $V_p - \theta_v$ data with (symbols) and without (lines) uniform random noise.

It can be seen from Figure 3.5 that the inversion procedure yielded close estimates of the porosity, density, dry density and moisture content. However, due to an area of similar error

values in the inversion error matrix, some difference between the modeled and the inverted values are noticed. This issue is discussed in detail in the following Section.

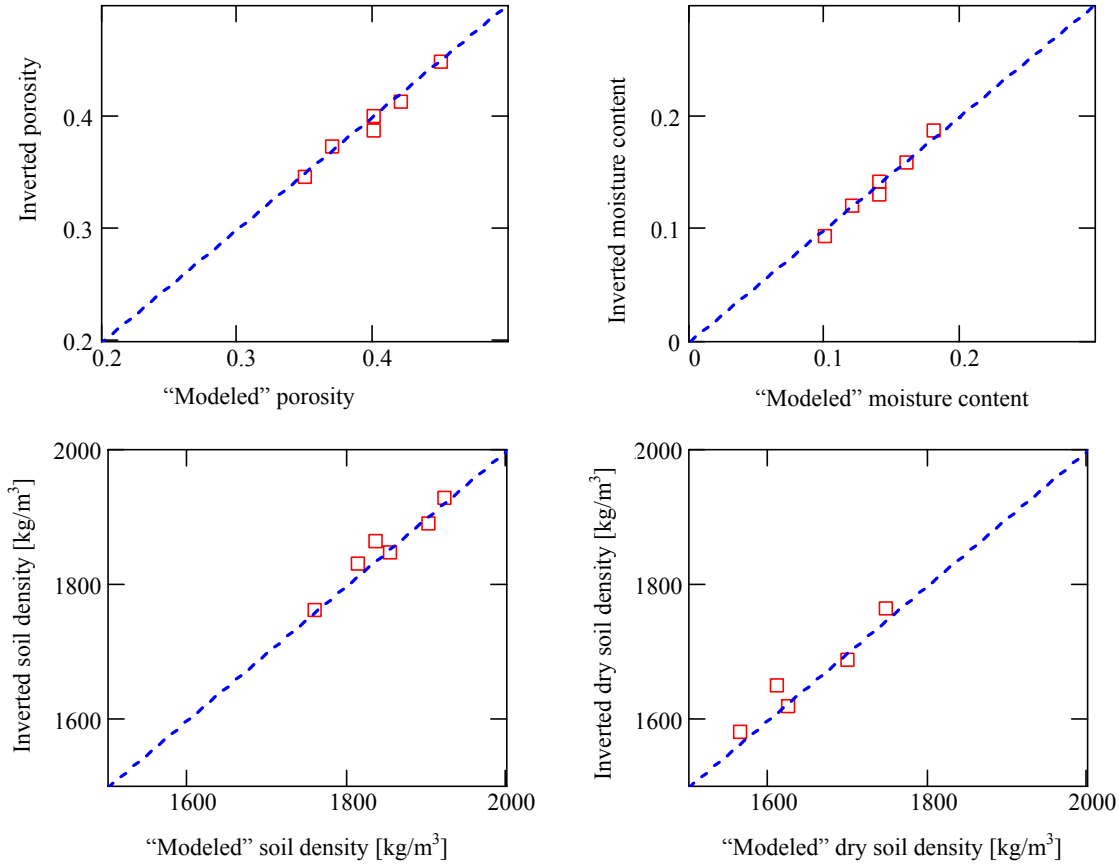


Figure 3.5. Synthetic data versus inverted parameters: (a) porosity, (b) water content (c) density and (d) dry density.

3.3.3 Inversion Procedure Sensitivity Analysis

As mentioned earlier, soils' density and water content are calculated from the collected volumetric water content and P-wave velocity data by means of a minimum error inversion procedure. The inversion algorithm determines the combination of n , m , and G_0 values that yield the minimum difference between the measured and calculated P-wave velocity values. In order to investigate the effects of model parameters on the solution convergence, a sensitivity analysis of the suggested model and inversion procedure is performed.

In this analysis, a set of assumed parameters was used to generate a θ_v - V_p curve using the proposed model (Figure 3.6). Then, the inversion procedure was conducted on data points selected along the generated curve, and the inverted model parameters were obtained using the inversion procedure. Figure 3.7 illustrates the selected data points and the fitted curve obtained from the inversion procedure. The results were used to perform a detailed analysis of the inversion procedure and to examine the behavior of the inversion error values.

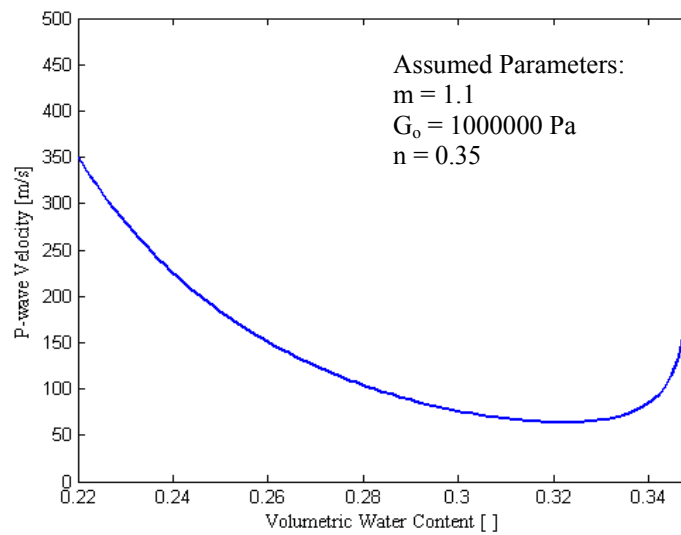


Figure 3.6. θ_v - V_p curve used for the sensitivity analysis

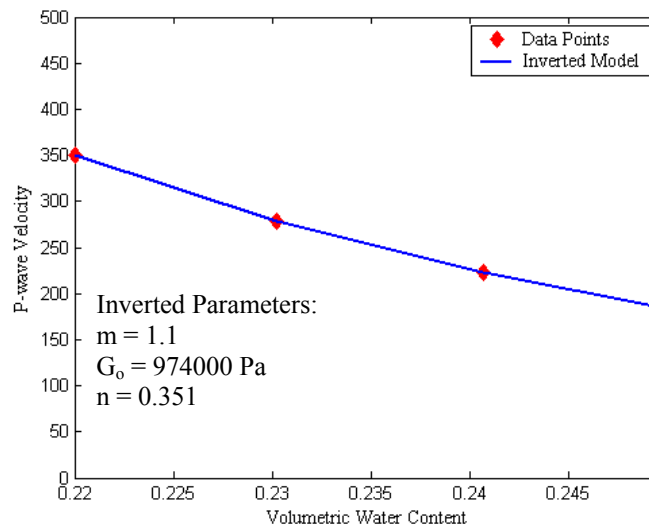


Figure 3.7. Selected data points and inverted model response

During the inversion process, a 3D data matrix was generated to store the error values generated for every combination of the three model parameters. To facilitate the visualization of the error surfaces, three 2D planes representing the error values for the combination of two model parameters while the third was held constant are displayed. Figure 3.8 illustrates the x,y and z sections corresponding to constant values of n , m , and G_o values, respectively.

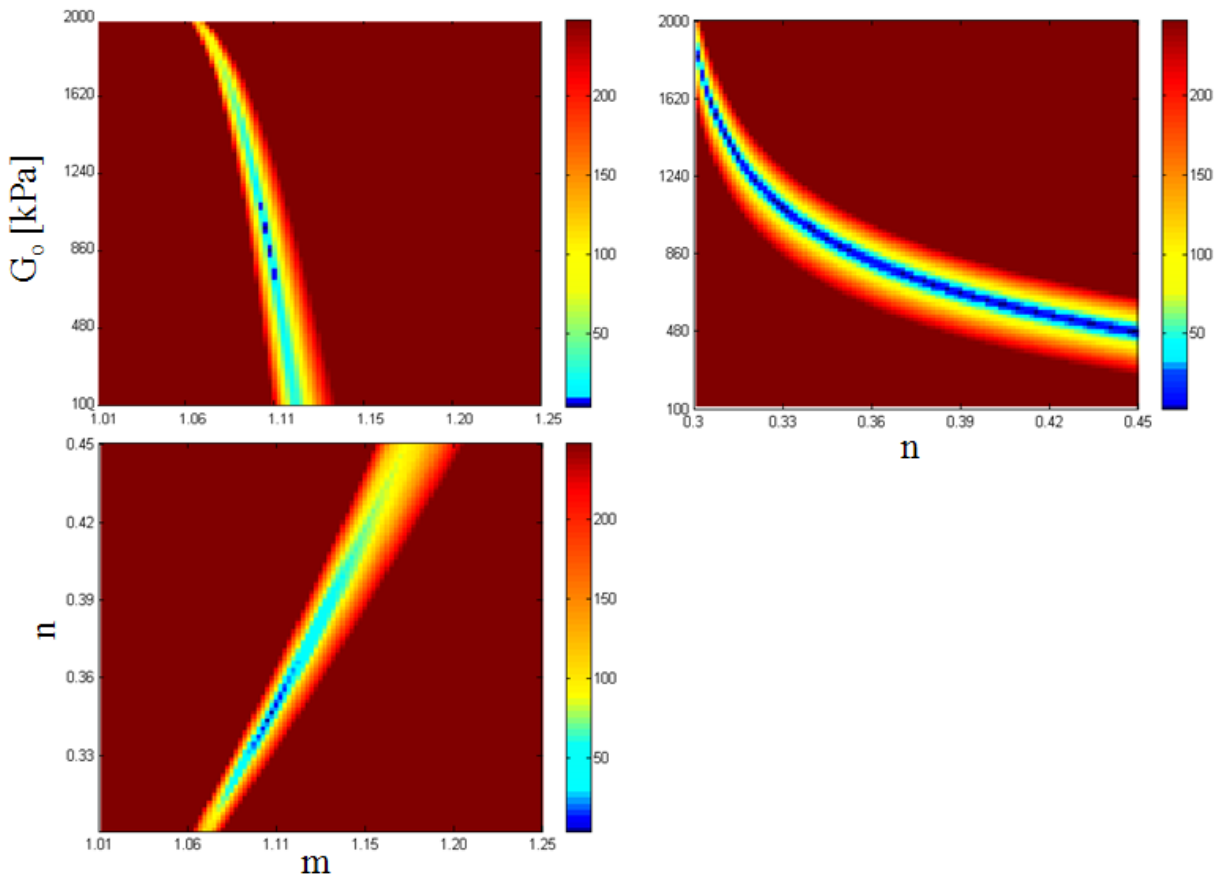


Figure 3.8. 2-D projections of error matrix at: (a) constant n , (b) constant m , (c) constant G_o

It can be seen from the error values in Figure 3.8 that while a minimum error value can be clearly identified in the cases of constant n and constant G_o , it is difficult to identify such value for the case of constant m . In this case, a band of similar error values can be identified along a wide range of n and G_o values. In other words, different combinations of n and G_o could yield similar error values making it difficult to identify a unique minimum error solution (i.e.,

non-unique solutions are obtained). While in most cases a single minimum error solution can be still obtained using this method, the presence of similar error values can introduce uncertainty in the obtained solution.

To further investigate how the error values are affected by the different model parameters, further analyses were performed where after the model parameters were obtained, each parameter was studied separately. Therefore, two parameters were held constant while the third was changed to determine how prominent the effect of that parameter is on the calculated error value. Figure 3.9 illustrates the effect of changing different model parameters on calculated error values. The higher change in the error values makes it easier to identify the minimum error point.

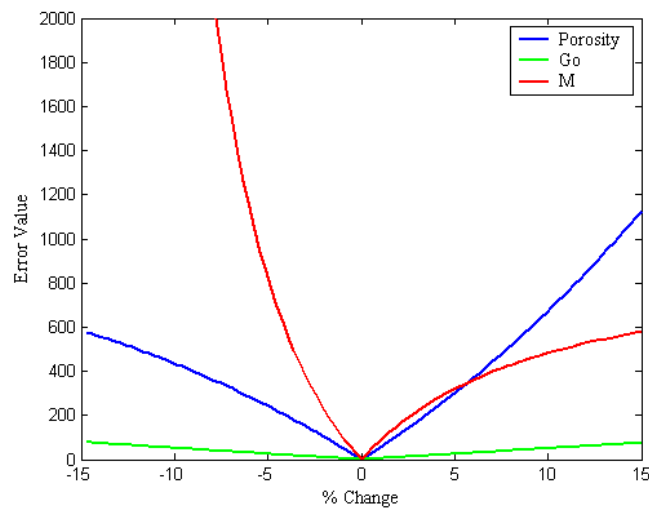


Figure 3.9. Change in error values with changing individual model parameters

Figure 3.9 demonstrates that the value of parameter m has the most profound effect on the error values, where a slight change in m will cause a sharp change in the calculated error values. This helps in the convergence of this parameter in the inversion procedure where it is relatively easy to recognize the value of parameter m that yields the minimum error. The value of parameter n has a smaller effect on the obtained error where smaller error values were obtained for the range of change of $\pm 5\%$. On the other hand, it can be seen that the value of G_0 has the

smallest effect on the obtained error values. Such small change in the calculated error value could cause some difficulty in obtaining the correct G_o value because of its close similarity to the error values obtained at neighboring G_o values.

3.4 Experimental Setup and Soils Description

In order to verify the validity and range of applicability of the proposed procedure, an experimental program was implemented where the procedure was applied to a variety of soil types consisting of different mixtures of sand, silt and clay. The following sub-Sections include a description of the soils used in the experiments, calibration of the TDR system and the results of the different verification experiments.

3.4.1 Description of Soils

Four basic types of soils were used in the experimental program: Kaolinite clay (C), fine construction sand (S), silt (M1) and silty sand (M2). Figure 3.10 depicts the grain size distribution curves for the four soil types. The properties and USCS classification of the four soil types are summarized in Table 3.2. These different types of soils were selected due to the range in grain size, plasticity, and electrical conductivity. Grain size distribution controls the capillary forces exerted in the particles (Cho and Santamarina, 2001), plasticity indicates water absorption and swelling potential (Mitchell, 1993), and electrical conductivity controls the range of application of the TDR methodology to evaluate volumetric water content in soils (Robinson et al., 2003).

Table 3.2. Properties of Basic Soil Types

Soil Label	Specific Gravity (G_s)	Liquid Limit	Plasticity Index	USCS Classification	D_{50} (mm)	C_c	C_u
C	2.59	43.6	18.2	CL	0.0012	N/A	N/A
M1	2.64	35.5	9.0	ML	0.015	6.64	21.25
M2	2.68	17.5	2.7	SM	0.21	0.66	440
S	2.65	N/A	N/A	SP	0.31	1.22	1.83

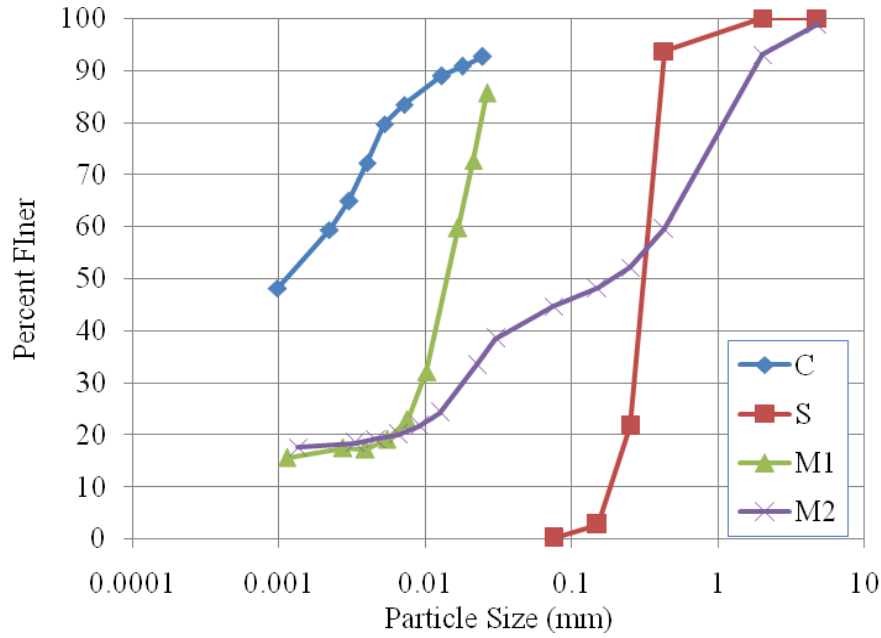


Figure 3.10. Particle size distribution for the basic soil types.

3.4.2 TDR System Calibration

The TDR system used in this study was model 6050X3K1-MiniTrase Kit manufactured by Soilmoisture Equipment Corp. It used a two probe configuration. The probes were 150 mm long and 50 mm apart. The system used a manufacturer-developed curve to estimate the volumetric water content values (θ_v) from the measured dielectric permittivity value (κ – Figure 3.11). In order to verify the proper performance of the TDR system, a series of experiments were performed to compare the volumetric water content values obtained by the TDR system with the values obtained by weight and volume measurements. This was achieved by compacting different soils in a cylindrical mold, recording their weight and volume to calculate the density and obtaining the gravimetric water content according to oven drying procedure. These measurements were used to calculate the soils volumetric water content. TDR readings were taken in the mold and the obtained θ_v values were compared with calculated values. Table 3.3 presents the different soil mixtures used in the calibration procedure.

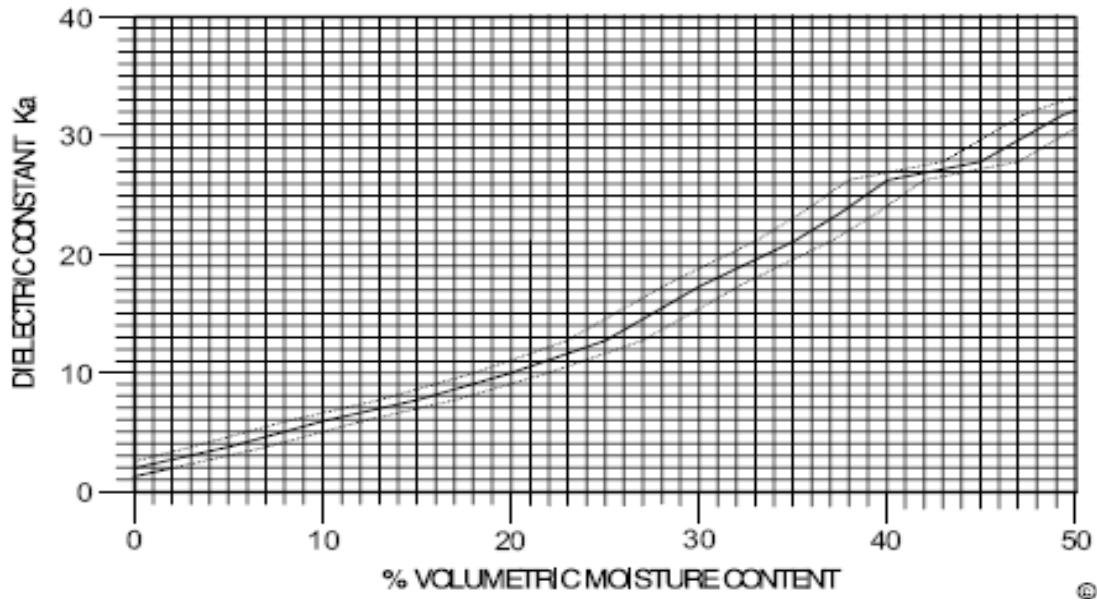


Figure 3.11. θ_v - κ curve developed by Soil Moisture Inc.

Table 3.3. Soil Combinations Used for TDR Calibration

Specimen Name	Percent of Soil			
	C	S	M1	M2
C	100	0	0	0
S	0	100	0	0
M1	0	0	100	0
M2	0	0	0	100
C-M1	67	0	33	0
C-S	67	33	0	0
C-S-M1	33.3	33.3	33.3	0
M1-C	33	0	67	0
M1-S	0	33	67	0
S-C	33	67	0	0
S-M1	0	67	33	0
C-S-M2	33.3	33.3	0	33.3
M2-S-C	25	25	0	50
M2-S	0	50	0	50

The results of the TDR system calibration are displayed in Figure 3.12. It can be seen that the volumetric water content values obtained by TDR measurements over-estimate θ_v by about 14%. Therefore, a correction factor of 0.88 was applied to the readings of the TDR system

in order to obtain more accurate θ_v values. The difference between the calculated and measured θ_v values could arise from the difference of soil types used in the development of the manufacturer curve and the ones used in the laboratory calibration.

3.4.3 Experimental Setup Description

An experimental setup was developed to implement the proposed procedure (Figures 3.13 and 3.14). It consisted of a TDR system to measure the volumetric water content of the soil specimens, two piezoelectric accelerometers (model: PCB 353B16) to measure the P- wave velocity, and a digital oscilloscope (model: Fluke 192 Scopemeter) to record the wave arrival data from the accelerometers. A special compaction mold made of plastic was used in this procedure since using a metal mold could result in the interference of the propagation of the electromagnetic waves through the compacted soils. The mold has a diameter of 150 mm and a height of 200 mm.

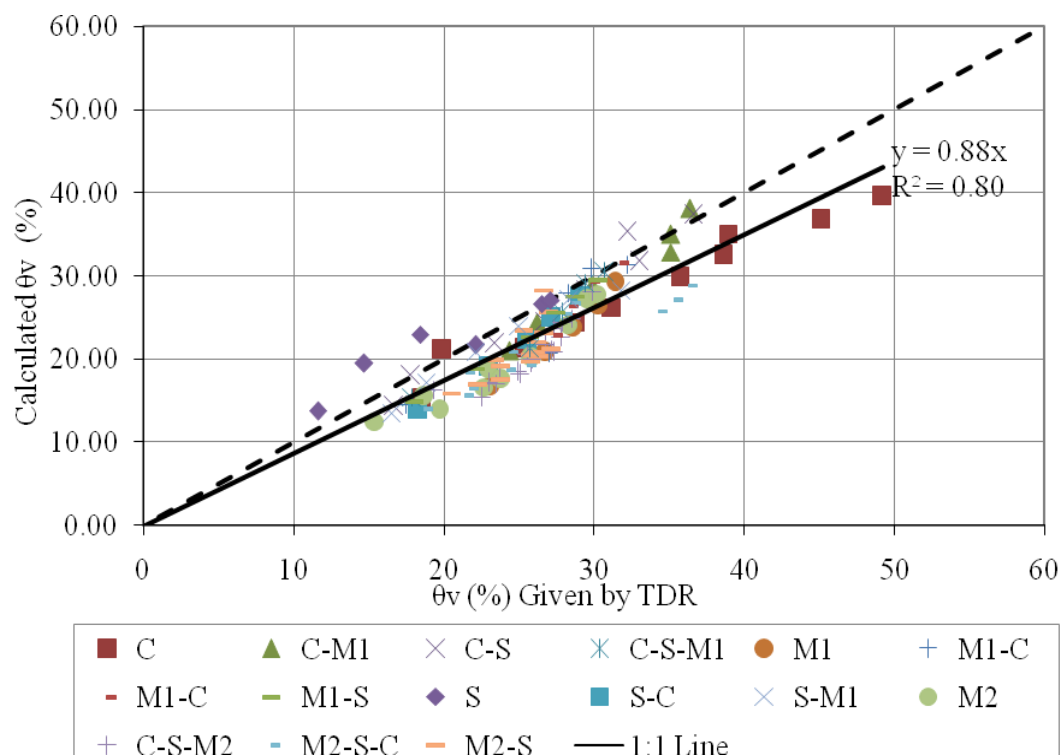


Figure 3.12. TDR System Calibration

A water injection hole was drilled at the center of the mold after the soil was compacted. Accelerometers and TDR rods were embedded into the soil. In some stiff soils, a special guide template was used to insure that the TDR rods remain straight while being driven into the soil. The volumetric water content was measured using the TDR system. Then, a P-wave was generated by dropping a metal ball along the axis of the accelerometers. Knowing the distance between the accelerometers, the wave velocity was determined by measuring the time difference between the arrivals. Figure 3.14 shows typical accelerometer traces. After recording the volumetric water content and the wave arrival data, the water content of the soil was increased by adding water to the injection hole. After the water was allowed to dissipate into the soil, another wave velocity and volumetric water content measurements were made. This procedure was repeated until enough data points were generated to apply the inversion procedure. After the inverted parameters were obtained, the density and moisture content were calculated and compared to the values measured by weight-volume relationships.

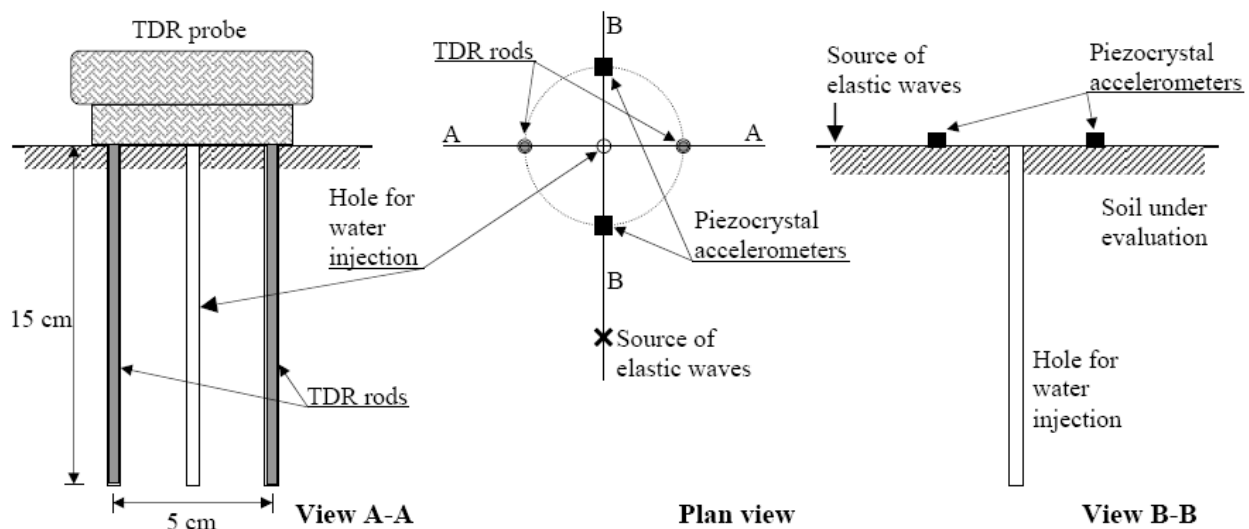


Figure 3.13. Schematic of the proposed test setup for the evaluation of in situ density and moisture content by means of combined electromagnetic and elastic wave propagation (Fratta et al. 2005).

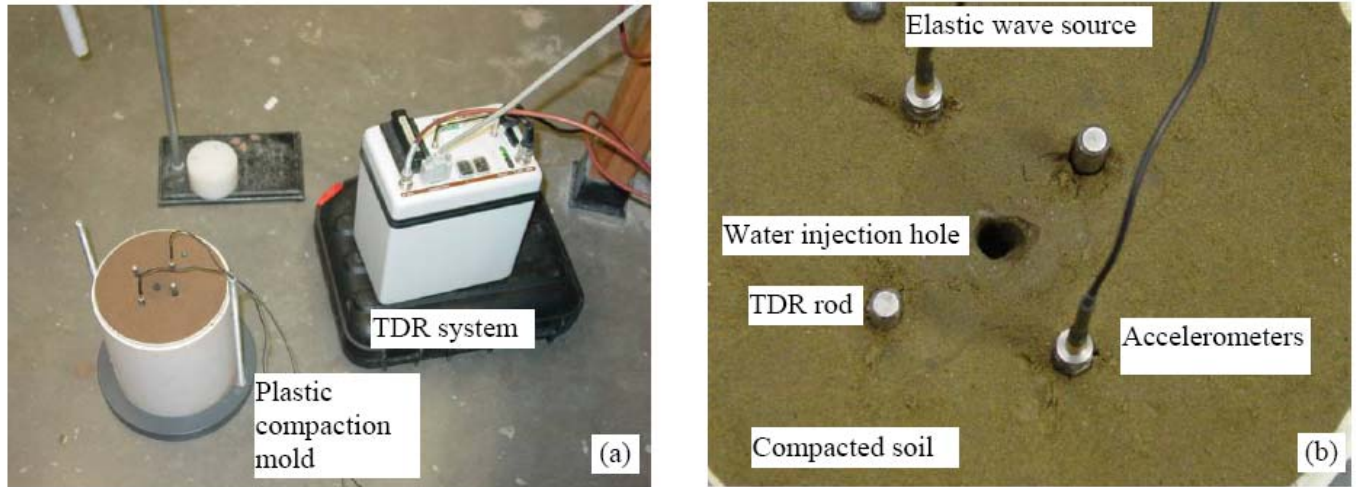


Figure 3.14. Application of the testing methodology to laboratory testing: (a) Plastic compaction mold and TDR systems and (b) Detail of sensor and water injection hole in a compacted soil.

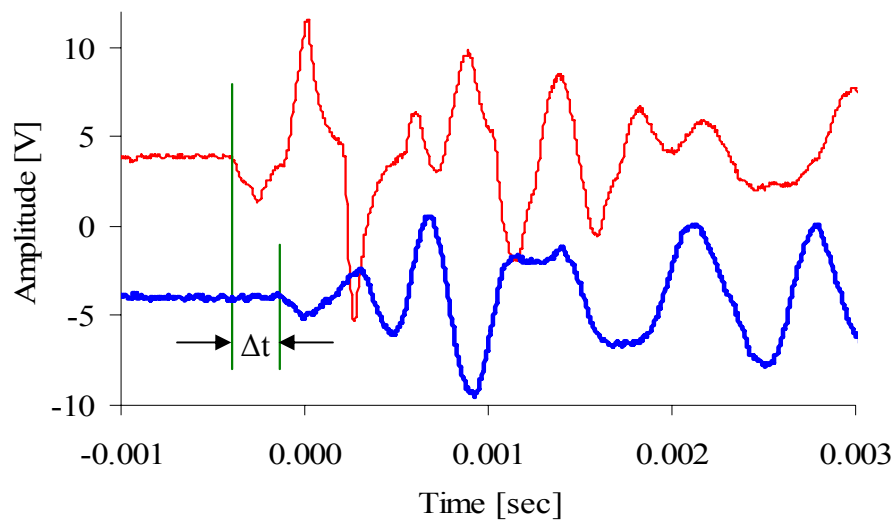


Figure 3.15. Typical accelerometer traces

3.4.4 Specimens Description

To verify the accuracy and range of applicability of the proposed methodology and inversion procedure, a testing program was developed that includes testing different mixtures of the four basic soil types described in Section 3.4.1. To prepare the soil specimens, the soil mix components were thoroughly mixed to obtain a homogeneous soil mixture. Then, water was

added to achieve the desired initial water content. Table 3.4 summarizes properties of the tested specimens. Soils were compacted in three layers. They were compacted in the mold at three compaction energy levels and different initial water content values. The different compaction efforts were achieved by varying the number of blows used to compact the soil layers. The energies used were: 239, 398 and 592 kJ/m³, which correspond to 15, 25 and 37 blows/layer, respectively, using the standard Proctor compaction hammer.

Table 3.4. Summary of tested soil specimens.

Specimen Name	Percent of Soil				Blows per layer	Number of experiments
	C	S	M1	M2		
C-S-M1-15	33.3	33.3	33.3	0	15	6
C-S-M1-25	33.3	33.3	33.3	0	25	4
C-S-M1-37	33.3	33.3	33.3	0	37	8
M1-C-S-15	25	25	50	0	15	4
M1-C-S-25	25	25	50	0	25	5
M1-C-S-37	25	25	50	0	37	5
M1-S-15	0	50	50	0	15	5
M1-S-25	0	50	50	0	25	5
M1-S-37	0	50	50	0	37	5
M1-15	0	0	100	0	15	4
M1-25	0	0	100	0	25	5
M1-37	0	0	100	0	37	5
C-M2-15	33	0	0	67	15	5
C-M2-25	33	0	0	67	25	8
C-M2-37	33	0	0	67	37	2
M2-15	0	0	0	100	15	3
M2-15	0	0	0	100	15	4
M2-15	0	0	0	100	15	5
						88

3.5 Experimental Results and Discussion

The proposed methodology to determine the density and moisture content was performed on all the soil specimens listed in Table 3.4. The inversion algorithm was then used to estimate the three model parameters: n , m and G_o . Bulk density, dry density and moisture content were

then calculated and compared to values obtained by weight and volume measurements. This Section presents typical experimental data along with their inversion results. A summary of the results of all experiments and a comparison of all the measured and calculated parameters are also presented.

After each experiment was performed, the traces from the accelerometers were analyzed to obtain the P-wave velocity at the different θ_v values for all experiments. These values were plotted to ensure that the proper trend was obtained. The experimental data was then input into the MATLAB inversion program in order to calculate the model parameters. A wide range of the three model parameters was initially set in the inversion program in order to find the approximate region where the minimum error solution was located. The ranges of the three model parameters were narrowed down in the second run in order to “fine tune” the solution. A narrower range of values results in a smaller step size yielding an improved data resolution. Figure 3.16 illustrates typical experimental V_p - θ_v data along with the curve obtained by the inversion procedure. It also lists the measured and calculated model parameters.

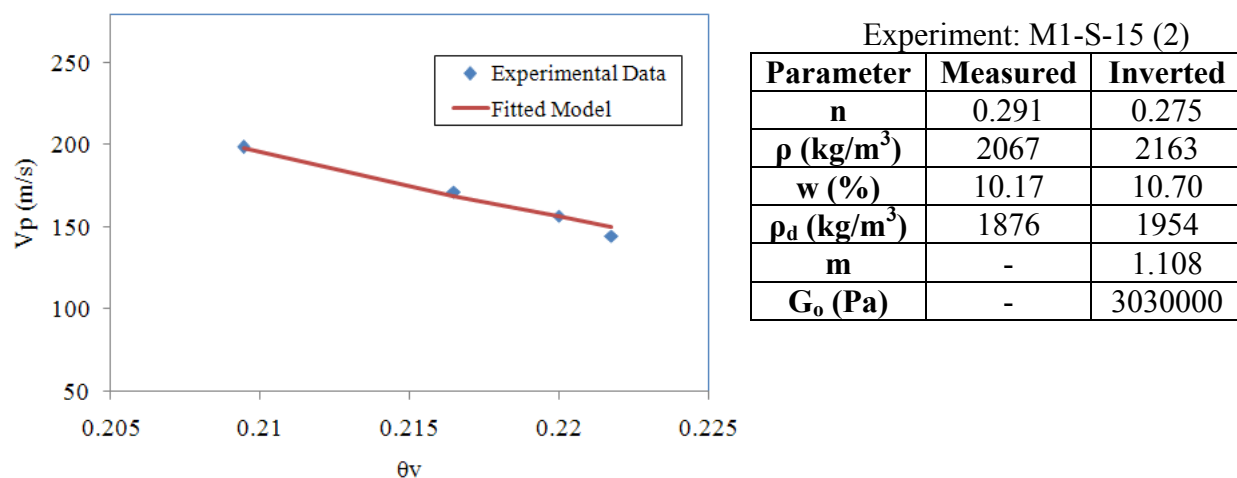


Figure 3.16 Typical Experimental Results, fitted curve and comparison between the measured and calculated parameters.

The inversion procedure was performed on the data obtained from all soil specimens (i.e., Table 3.4). Calculated density, dry density, and water content values were compared with the values obtained from weight and volume measurements. The results are depicted in Figure 3.17 along with the histograms of the obtained error values. Moreover, the error percentage between the measured and calculated values was calculated for each set of experiments (same soil type with same number of blows) and plotted in Figure 3.18, the error percentage was calculated as (X represents the value of the density or water content):

$$E(\%) = \text{Average} \left[\frac{|X_{\text{measured}} - X_{\text{calculated}}|}{X_{\text{measured}}} \cdot 100\% \right] \quad (3.12)$$

It can be seen from Figure 3.17 and 3.18 that the proposed methodology was able to predict the soils density and dry density within $\pm 10\%$ of the measured values in most cases. It can also be seen that there was no noticeable tendency to over- or under-estimate that density value. The data points were nearly evenly distributed about the 1:1 line in Figure 3.17a. On the other hand, most of the predicted water content values were within $\pm 20\%$ of the measured values. It was also noticed that the majority of the calculated water content values lied above the 1:1 line indicating that the proposed methodology tends to over-estimate the water content. These errors could be caused by many factors including measurement errors and numerical inversion errors.

3.5.1 Sources of Error

The error in the calculated density and moisture content values is a result of a combination of different error sources taking place either during data collection and interpretation or during the inversion procedure, the main sources of error can be summarized as:

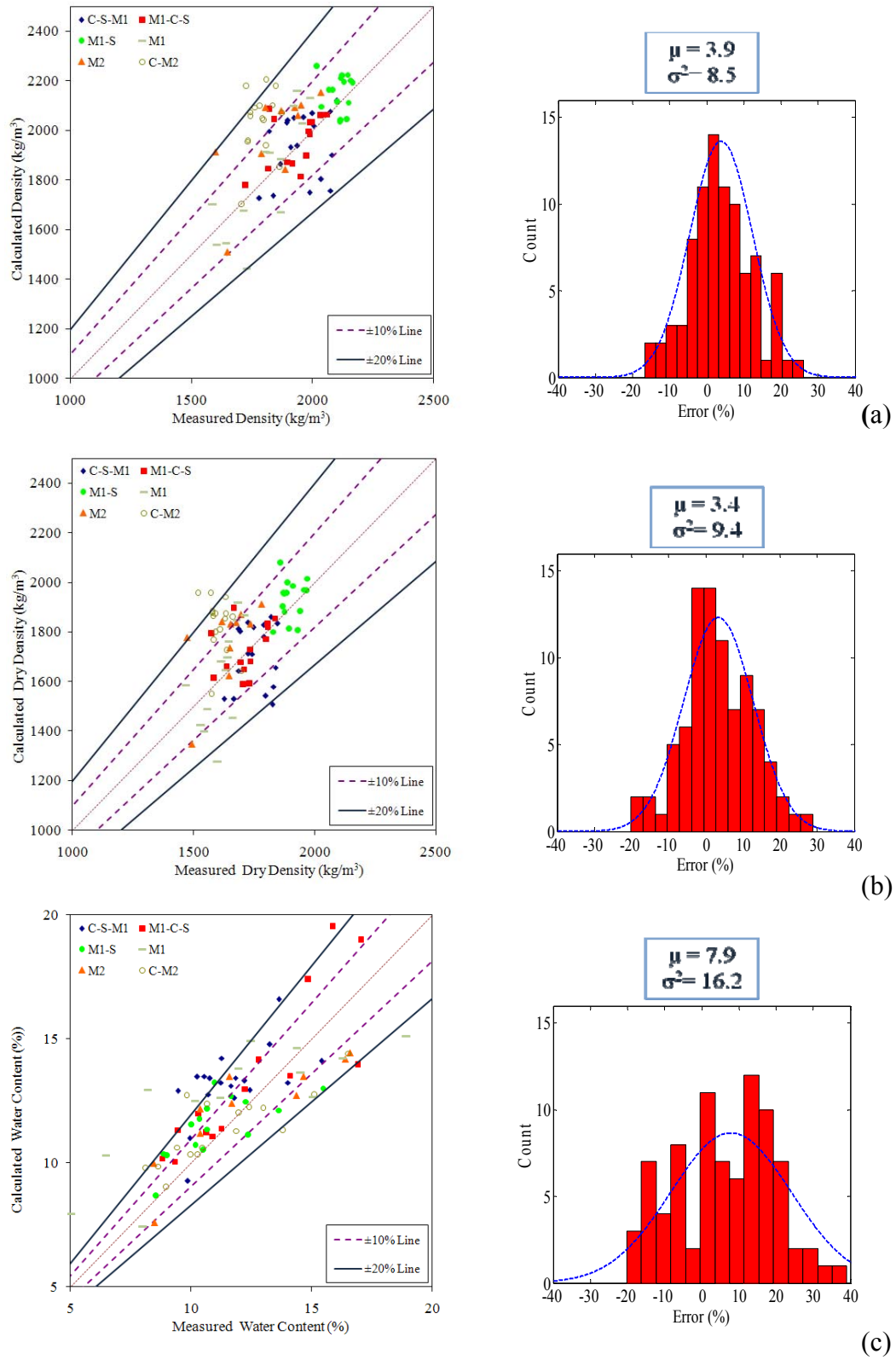


Figure 3.17. Comparison between measured and calculated values along with error distributions
(a) density (b) dry density (c) water content

- *P-wave velocity and volumetric water content estimation:* such errors can arise during the evaluation of the travel times of P waves and electromagnetic waves resulting in errors up to about 10%.
- *Non-homogeneous distribution of injected water:* the value of the volumetric water content obtained by the TDR system represent the average water content value along the TDR sampling volume, however, during the water injection process, the distribution of water along this volume might not be homogenous resulting in a measurement error. Moreover, non-homogeneity causes the sampling area of the TDR to be different from the sampling area of the P-wave velocity. The output from the TDR system represents the average values along the depth of 150 mm whereas P-wave velocity is measured near the surface.
- *Local saturation of the region around the injection hole:* increasing the soils water content yields a reduction in P-wave velocity, however when the soil is fully saturated, the P-wave velocity travels in water at a much higher velocity (~ 1500 m/s). Therefore, having a locally saturated zone (even a few millimeters in size) around water injection hole could results in errors in the measured P-wave velocity.
- *Inversion procedure errors:* as shown earlier, it is relatively easy to identify the minimum error point in the G_o -m and n-m planes. However, in the n- G_o planes, a band of similar error values was noticed. The presence of this “similar error zone” makes it difficult to identify the inversion solution because of the close similarity to the calculated error values for the points around it along that band.

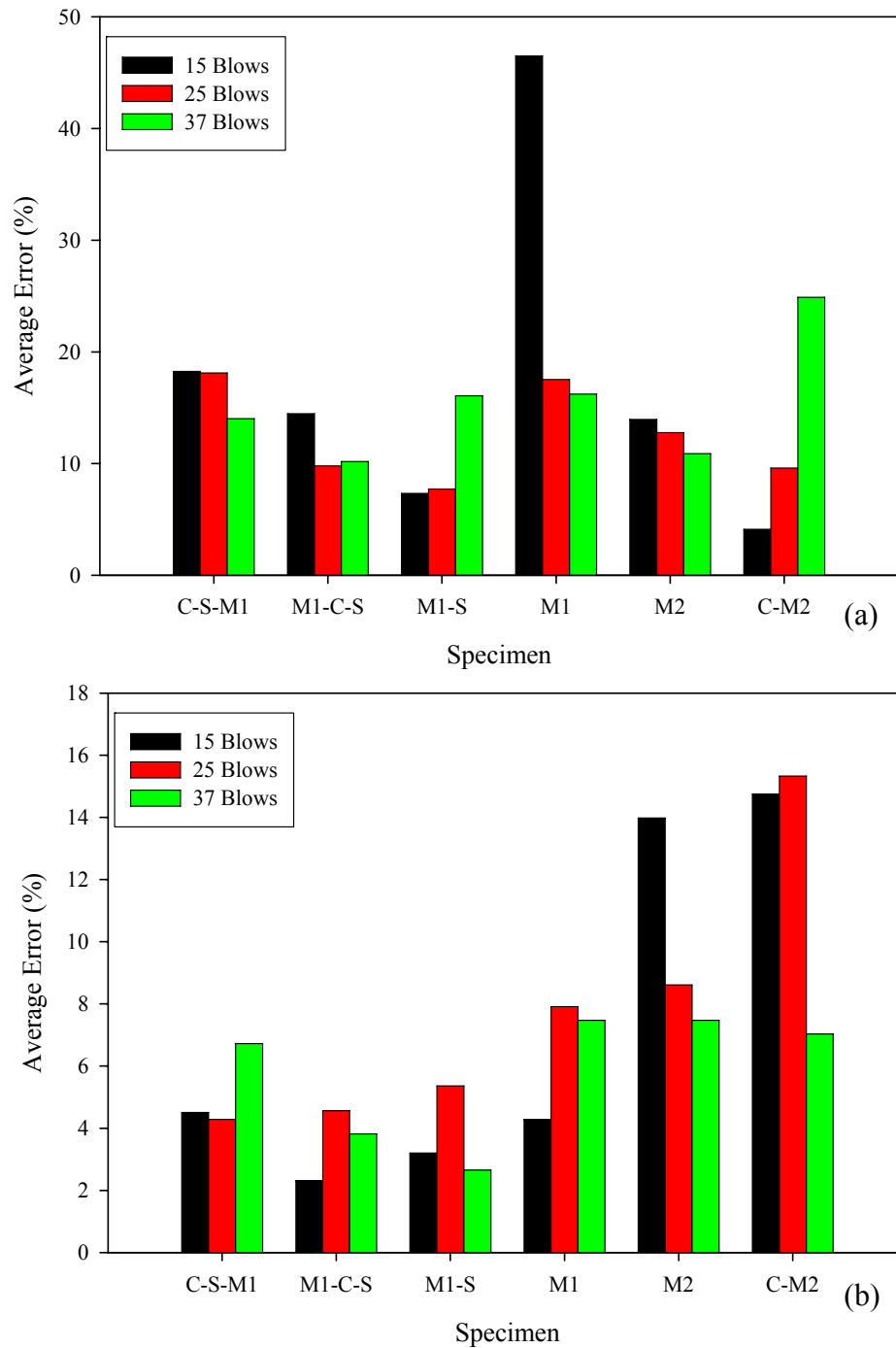


Figure 3.18. Error percentage calculated for all specimens (a) water content, (b) density

3.5.2 Effect of Injection Hole(s) Configuration

During implementation of the suggested experimental procedure, the volumetric water content of the soil should be increased. This was achieved through injecting water in single or

multiple injection holes carefully drilled in the sampled volume. Initially, a single injection hole was used to achieve this purpose (Figure 3.19a). This configuration was chosen to be able to raise the volumetric water content without causing much disturbance to the surrounding soils. However, due to the long time required for the water to dissipate into the soil, and in an attempt to achieve a more homogeneous water distribution, two sets of experiments were carried out using a multiple injection hole configuration: M2 and C-M2. This configuration consisted of five smaller injection holes distributed around the sampled volume (Figure 3.19b).

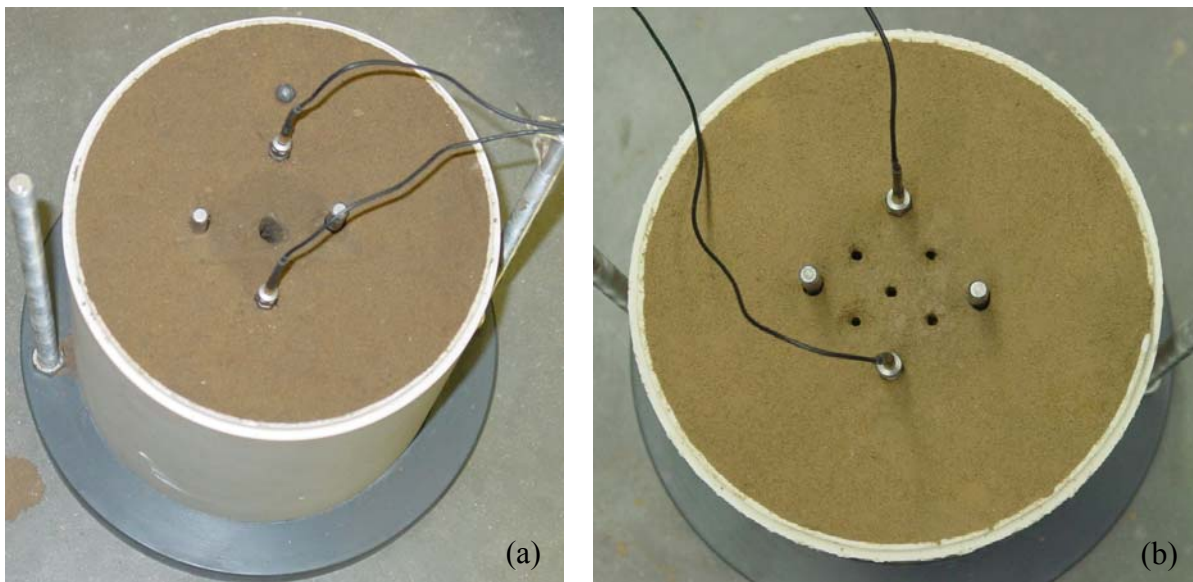


Figure 3.19. (a) Single injection hole configuration, (b) multiple injection hole configuration

By examining the results of the inversion procedure, it can be seen that while no significant improvement was achieved in the water content measurements, the density error values for the two sets of experiments with the multiple hole configuration were noticeably higher than the ones with the single hole configuration. This is believed to be caused by the higher soil disturbance caused by drilling the five injection holes instead of one. Therefore, although more time is required for the water to dissipate into the soil, the single injection hole configuration seems to yield better results than the multiple hole configuration.

3.5.3 Effect of Combined Inversion Errors

As mentioned earlier, several factors could affect the quality of the measured and calculated parameters necessary to calculate the density and water content. Therefore, it is important to examine the effect of such errors on these parameters. Equations 3.13 through 3.15 indicate that the main parameters affecting the values of the density and moisture content are porosity (n) and volumetric water content (θ_v). Therefore, a sensitivity study was performed to assess the effect of measurement errors of these two parameters on the calculated density and water content values.

$$\rho = (1 - n)\rho_s + \theta_v \rho_w \quad (3.13)$$

$$w = \frac{\rho}{\rho_d} - 1 \quad (3.14)$$

$$\rho_d = (1 - n)\rho_s \quad (3.15)$$

To perform the sensitivity analysis, initial values for “ n ” and “ θ_v ” were assumed and used to calculate the initial density and moisture content: ρ_o and w_o . Then, an error of $\pm 10\%$ was added to the porosity and volumetric water content, and the values of density and moisture content were calculated for the different combinations of the parameters ranging from -10% to $+10\%$. The resulting error in the calculated density and moisture content for each combination was calculated using Equation 3.16. The resulting error plots are illustrated in Figure 3.20.

$$\text{Density: } E_{i,j}(\%) = \frac{\rho(n_i, \theta_j) - \rho_o}{\rho_o} \cdot 100\% \quad (3.16a)$$

$$\text{Water Content: } E_{i,j}(\%) = \frac{w(n_i, \theta_j) - w_o}{w_o} \cdot 100\% \quad (3.16b)$$

Figure 3.20 shows that the full spectrum of porosity and volumetric water content error values yield an error of about $\pm 6\%$ in the calculated density values. On the other hand, the same

range of error values yields and error of $\pm 15\%$ in the calculated water content values. This indicates that the error in the calculated water content is more sensitive to errors in the porosity and volumetric water content. This also justifies the higher error values obtained during the inversion of the experimental data.

3.6. Theoretical Framework for an Alternative Methodology

As seen from the previous experimental study and error analysis, although the proposed methodology yielded close estimates of the density and moisture content, similar error values in the inversion algorithm caused uncertainty in the obtained solutions. Therefore, to further improve the current methodology, other physically meaningful parameters should be added to introduce more constraints to the inversion procedure. In this section, the feasibility of adding the shear wave velocity V_s to the obtained measurements is investigated. A new model will be derived assuming that V_s is measured along with V_p and θ_v .

Soils bulk modulus B_{soil} is calculated as the sum of the bulk moduli of the soils skeleton B_{sk} and the suspension B_{sus} . The definitions of these parameters are presented in Section 3.2. B_{sk} can be expressed as a function of the shear modulus and Poisson's ratio (Equation 3.17). On the other hand, substituting the values of B_{fl} from Equation 3.4 into Equation 3.5 yields the expression for B_{sus} in Equation 3.18. B_{soil} can then be calculated by adding B_{sk} and B_{sus} (Equation 3.19).

$$B_{\text{sk}} = \frac{2}{3} G_{\text{soil}} \cdot \left(\frac{1 + \nu}{1 - 2\nu} \right) \quad (3.17)$$

$$B_{\text{sus}} = \left(\frac{n}{\left(\frac{S_r}{B_w} + \frac{1 - S_r}{B_a} \right)^{-1}} + \frac{1 - n}{B_g} \right)^{-1} \quad (3.18)$$

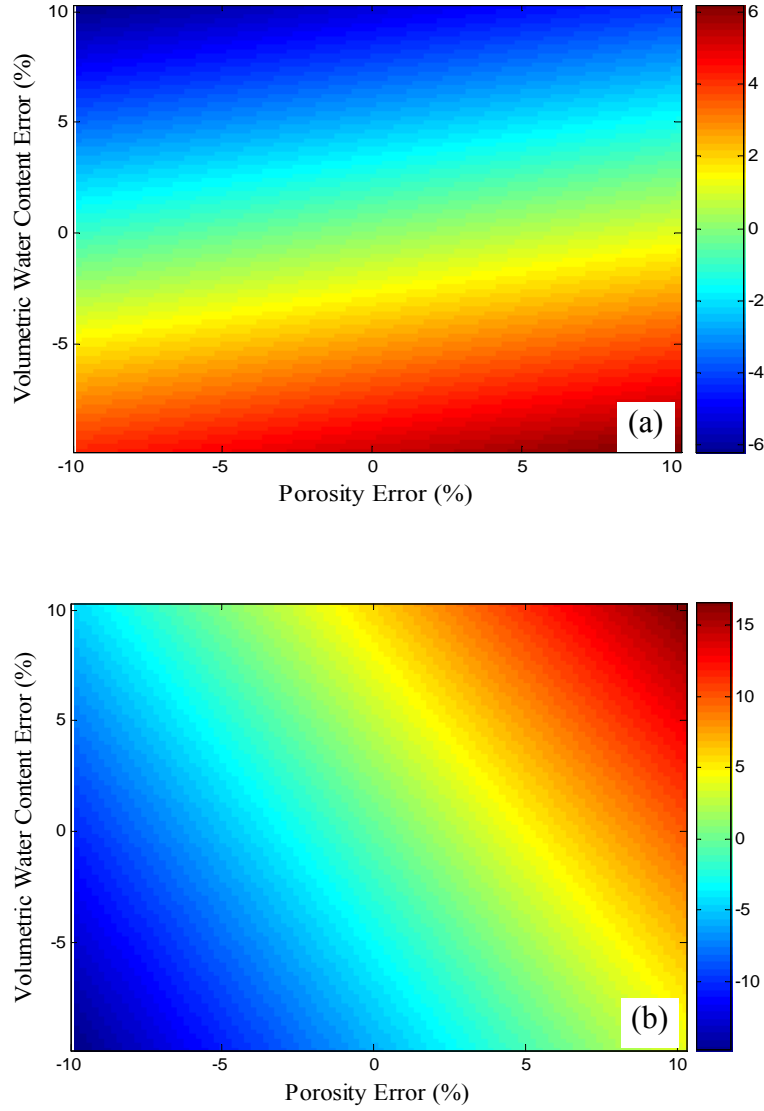


Figure 3.20. Percent change in (a) density and (b) water content with $\pm 10\%$ change in porosity and volumetric water content

$$B_{\text{soil}} = \left(\frac{n}{\left(\frac{S_r}{B_w} + \frac{1-S_r}{B_a} \right)} + \frac{1-n}{B_g} \right)^{-1} + \frac{2}{3} G_{\text{soil}} \cdot \left(\frac{1+\nu}{1-2\nu} \right) \quad (3.19)$$

The shear stiffness G_{soil} can be calculated from the density (ρ) and V_s using Equation 2.17. Substituting the values of G_{soil} and ρ from Equations 2.17 and 3.2 into Equation 3.19 yields Equation 3.20. Moreover, the P-wave velocity (Equation 3.3) is dependent on B_{soil} , G_{soil} , and ρ .

Substituting the values these parameters from Equations 3.20, 2.17 and 3.2 , respectively, into Equation 3.3 yields a new P-wave velocity model (Equation 3.21).

$$B_{\text{soil}} = \left(\frac{n}{\left(\frac{S_r}{B_w} + \frac{1-S_r}{B_a} \right)} + \frac{1-n}{B_g} \right) + \frac{2}{3} ((1-n) \cdot \rho_s + \theta_v \cdot \rho_w) \cdot V_s^2 \cdot \left(\frac{1+v}{1-2v} \right) \quad (3.20)$$

$$V_p = \sqrt{\frac{\left(\frac{n}{\left(\frac{S_r}{B_w} + \frac{1-S_r}{B_a} \right)} + \frac{1-n}{B_g} \right) + \frac{2}{3} ((1-n) \cdot \rho_s + \theta_v \cdot \rho_w) \cdot V_s^2 \cdot \left(\frac{1+v}{1-2v} \right) + \frac{4}{3} [(1-n) \cdot \rho_s + \theta_v \cdot \rho_w] \cdot V_s^2}{(1-n) \cdot \rho_s + \theta_v \cdot \rho_w}} \quad (3.21a)$$

$$V_p = \sqrt{\frac{\left(\frac{n}{\left(\frac{\left(\frac{\theta_v}{n} \right)}{B_w} + \frac{1-\left(\frac{\theta_v}{n} \right)}{B_a} \right)} + \frac{1-n}{B_g} \right) + \frac{2}{3} ((1-n) \cdot \rho_s + \theta_v \cdot \rho_w) \cdot V_s^2 \cdot \left(\frac{1+v}{1-2v} \right) + \frac{4}{3} [(1-n) \cdot \rho_s + \theta_v \cdot \rho_w] \cdot V_s^2}{(1-n) \cdot \rho_s + \theta_v \cdot \rho_w}} \quad (3.21b)$$

It can be seen from Equation 3.21b that since V_s , V_p and θ_v can be experimentally measured, v can be estimated using Equation 2.22, and the bulk stiffnesses of the water, air and the solid grains are known, the only unknown in Equation 3.21b is the porosity (n). Therefore, using this model, with a *single* reading of V_p , V_s and θ_v , the soils porosity can be calculated. The density and moisture content can then be calculated using Equations 3.2 and 3.10.

Using this model requires in addition to the system described in the experimental study section a measurement of V_s . Since the drop ball method is not able to produce S-waves, other

techniques should be used. Potential alternatives include using piezoelectric elements or solenoids (see for example Schneider, 2007). Further experimental investigation is needed to verify the ability of this method to accurately predict the density and moisture content.

3.7. Chapter Summary

This chapter presented a new methodology to non-destructively evaluate soils' in-situ density and moisture content by combining TDR and P-wave measurements. It also presented the results of an experimental program aiming to assess the validity of this methodology along with a numerical sensitivity analysis of the inversion procedure used to calculate these parameters.

It was noticed that although this procedure yielded promising results for the determination of the density and moisture content, the majority of the calculated density values were within 10% of the measured values, and the majority of the water content were within 20% of the measured values. These differences between the measured and calculated values could be caused by several sources of error during acquiring and interpretation of the experimental data as well as numerical errors during the inversion procedures. Further work is required to eliminate or reduce these sources of error, including:

- Introducing more constraints to the inversion procedure to reduce the effect of the similar error values and reduce the area where the solution could be located.
- Utilizing a water injection hole configuration that achieves a fairly homogeneous water distribution with minimal soil disturbance.
- Using a more robust method to generate and acquire P-waves. Signal stacking could result in clearer wave traces reducing the error associated with determining the wave travel time. The use of bender elements is a potential alternative to the drop-ball method.

A theoretical framework for a new methodology for the determination of in situ density and moisture content was also presented in this chapter. It requires a single measurement of V_s ,

V_p and θ_v to calculate the soils porosity, density and moisture content. The accuracy and range of applicability of this method should be determined by an experimental study; however, it presents a promising technique to acquire much faster measurements since it only requires one data point.

CHAPTER 4

A SUCTION-CONTROL APPARATUS FOR THE MEASUREMENT OF P AND S-WAVE VELOCITY IN SOILS

4.1 Introduction

The strength and stiffness characteristics of soils depend on the interaction of the different phase constituents (i.e., particles, air, and water phases) and the acting stresses. In most geotechnical analysis, two-phase systems are assumed where soils are either fully saturated or dry. However in most real engineering applications, soils are neither completely dry nor fully saturated but rather partially saturated. For such conditions, the behavior of soils also depends on the interaction of capillary and chemical forces between particles, air, and water. Since water is attracted to most soil particle surfaces and can develop surface tension, capillary menisci form between particles in partially saturated soils. Capillary forces play a major role in governing the strength, stiffness, and volume change characteristics of unsaturated soils (Matyas and Radhakrishna 1968; Fredlund and Rahardjo 1993, Vinale et al. 1999).

The effective stress σ' is defined as the stress that controls the strength and deformation behavior of granular materials (Terzaghi 1936, Bishop and Blight 1963). In saturated soils, the effective stress is defined by the Terzaghi's effective stress principle as a function of the total stress σ and the pore fluid pressure u :

$$\sigma' = \sigma - u \quad (4.1)$$

However, in the case of three phase systems such as unsaturated soils, Terzaghi's principle is no longer valid. The presence of air in the soil pores induces a new force balance that incorporates capillary forces acting at the soil-air-water interface. Such forces have a direct effect

on the forces acting on particle contacts and greatly influence the macroscopic behavior of surface soils.

The effective stress in unsaturated soils depends on the total stress σ , air pressure u_a and water pressure u_w . Many researchers have proposed different forms of effective stress relationships (e.g., Croney et al. 1958; Aitchinson 1961; Jennings 1961). A definition of equivalent state of stress of unsaturated soils can be expressed as a function of two stress state parameters: net stress and matric suction,

$$\sigma' = \underbrace{(\sigma - u_a)}_{\text{net stress}} + \chi \underbrace{(u_a - u_w)}_{\text{matric suction}} \quad (4.2)$$

where χ is an experimentally determined coefficient that depends on the degree of saturation, porosity, soil type, history of wetting and draining, among other parameters (Bishop and Blight 1963; Aitchison 1960). Equation 4.2 shows how net stresses and matric suction determine the effective stress on unsaturated soil thus controlling its strength and stiffness. However, while Terzaghi's effective stress equation for saturated soils considers only global actions (total stress and pore water pressure); Equation 4.2 mixes global actions with local actions acting on the particle level (matric suction), which causes an inconsistency in describing the effective stresses in unsaturated soils (Cho and Santamarina 2001). A better representation of unsaturated soil behavior is obtained by using net stress and matric suction as two independent state variables to describe the state of stress (Bishop and Blight 1963; Fredlund et al. 1978).

Elastic wave measurements provide a unique tool to examine the small-strain stiffness response of soils. The propagation velocities of compression and shear waves through soil depend on the bulk and shear stiffnesses, respectively (Santamarina et al. 2005; Fratta et al. 2005). Therefore, monitoring the P and S-wave velocity changes in an unsaturated soil medium while controlling the net stress $(\sigma - u_a)$ and matric suction $(u_a - u_w)$ provides an excellent

opportunity to determine the effects of different stress components on the small strain stiffness of soils. It also allows monitoring the stabilization process of the capillary forces after the application of the pressure increments.

Elastic wave propagation measurements have been used in many geotechnical testing apparatuses to examine the small shear strain stiffness of soils (e.g., Dyvik and Madshus 1985, Thomann and Hryciw 1990, Gohl and Finn 1991, Silva et al. 2002). Brignoli et al. (1996) developed an apparatus capable of measuring elastic wave velocities in dry and saturated soil specimens tested under isotropic compression. On the other hand, Hoyos and Macari (2001) developed a stress-suction controlled true triaxial apparatus which was later modified by E. Macari and co-workers (Silva et al. 2002, Porras-Ortiz 2004) to include bender elements to monitor the S-wave velocity variations through unsaturated soil specimens tested under true triaxial conditions. Bender elements were also mounted on the top and bottom end platens of an oedometric drying cell by Cho and Santamarina (2001) to non-destructively examine the effect of capillarity on soil stiffness at low water contents.

This Chapter presents a new apparatus capable of monitoring both P and S-wave velocities in unsaturated soil specimens with controlled net stress and matric suction under isotropic or axisymmetric loading conditions. The simultaneous measurement of P and S-wave velocity in soils permits full characterization of small-strain elastic parameters and provides the ability to verify proposed models for evaluating the effect of saturation on the soil stiffness (i.e., Cho and Santamarina 2001; Fratta et al. 2005; Sawangsuriya 2006). The main advantage of this new test setup is that it permits simultaneous evaluation of the P and S-wave velocity under controlled suction and net stresses conditions while allowing the application of axial and lateral compressions and extensions according to the desired stress paths. A detailed description of the apparatus, pressure control system, and evaluation of experimental results are presented. The

effect of matric suction on the P and S-wave velocities in unsaturated soil specimens under isotropic state of stresses is investigated to better understand the role of matric suction on the inter-particle forces and small strain stiffness of unsaturated soils.

Experimental results obtained from this apparatus will enhance our understanding about elastic wave propagation properties in unsaturated soils. They could assist in the development of new wave velocity models that can be used to improve non-destructive methods for the evaluation of soil properties using elastic waves.

4.2 Elastic Wave Propagation in Soils

As presented in Sections 2.2, S- and P-wave velocities in soils are dependent on the soils shear and bulk moduli as well as the soils density (Equations 2.17 and 2.18). Therefore, by monitoring these velocities, important information about these parameters could be obtained. Moreover, the ratio of the P- to the S-wave velocities can be used to estimate Poisson's ratio ν of soils at small strains (Equation 4.3). Figure 4.1 depicts Poisson's ratio values for a typical range

of $\frac{V_p}{V_s}$.

$$\nu = \frac{\frac{1}{2} \left(\frac{V_p}{V_s} \right)^2 - 1}{\left(\frac{V_p}{V_s} \right)^2 - 1} \quad (4.3)$$

A typical range of P- to S-wave velocity ratio of 1.5 to 1.9 corresponds to Poisson's ratio values between 0.1 to 0.3. The bulk and shear moduli of soils depend on the stiffness of the mineral, air and water phases, porosity, degree of saturation, the interparticle force magnitude, and type of contact between mineral particles (Richart et al. 1970, Mavko et al. 1998). Wave propagation theory and mixture models show that both the P and S-wave velocities are related by the same porous media parameters (Fratta et al. 2005):

$$V_p = \sqrt{\frac{B_{\text{sus}} + B_{\text{ske}} + \frac{4}{3}(G_{\text{sus}} + G_{\text{ske}})}{\rho_{\text{soil}}}} = \sqrt{\frac{B_{\text{sus}} + \frac{4}{3}G_{\text{sus}} + \frac{2-2\nu_{\text{ske}}}{1-2\nu_{\text{ske}}}G_{\text{ske}}}{\rho_{\text{soil}}}} \quad (4.4)$$

$$V_s = \sqrt{\frac{G_{\text{sus}} + G_{\text{ske}}}{\rho_{\text{soil}}}} \quad (4.5)$$

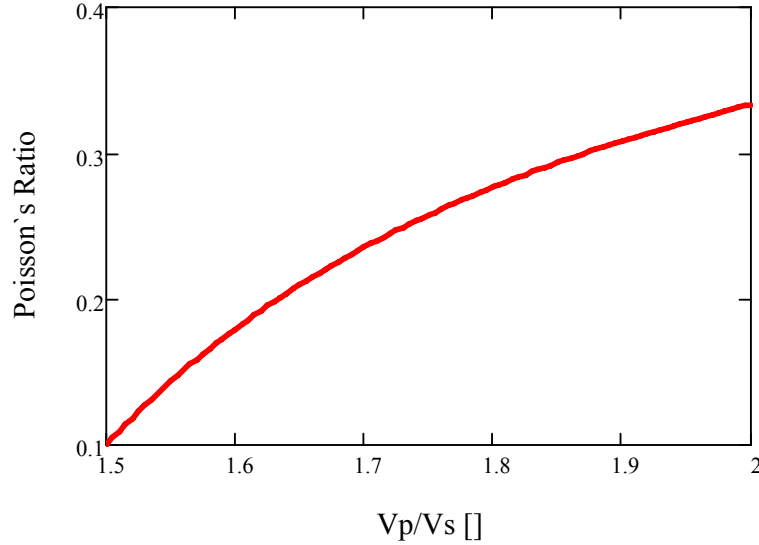


Figure 4.1. Poisson's ratios corresponding to a range of P- to S-wave velocity ratios

where B_{sus} and G_{sus} are the bulk and shear moduli of the suspensions (i.e., the mineral, water, and air phases without considering the interactions between mineral particles) and G_{ske} and ν_{ske} are the shear stiffness and Poisson's ratio of the skeleton. The shear stiffness of the skeleton depends on both the net stresses and matric suction. In unsaturated conditions, the presence of the air phase greatly reduces the bulk modulus of the suspension, therefore $B_{\text{sus}} + 4/3 \cdot G_{\text{sus}} \ll G_{\text{ske}}$ and the equations for P and S-wave velocities can be simplified to:

$$V_p = \sqrt{\frac{\frac{2-2\nu_{\text{ske}}}{1-2\nu_{\text{ske}}}G_{\text{ske}}}{\rho_{\text{soil}}}} \quad (4.6)$$

$$V_s = \sqrt{\frac{G_{\text{ske}}}{\rho_{\text{soil}}}} \quad (4.7)$$

In Equations 4.6 and 4.7, the shear stiffness of the skeleton depends on the effective stress; that is, it depends on both the net stresses and matric suction.

4.3 Apparatus Description

The laboratory apparatus used to measure P and S-wave velocities in unsaturated soil specimens under drained compression and controlled matric suction consists of four main components (Figure 4.2):

- A triaxial cell with an axial load piston modified to enable independent control of air pressure u_a , water pressure u_w , and cell pressure σ . The triaxial cell also houses the piezoelectric transducers required to generate and receive P and S-waves.
- A fluid pressure-control system consisting of three Geotac DigiFlow mechanical pumps with feedback control. The DigiFlow pumps are used to both generate and control the confining pressure in the triaxial cell as well as the air and water pressures applied to the soil specimen's pore space. The water and cell pressure pumps use digitally controlled DC servo-motors to move a piston in and out of a cylinder to pressurize the fluids and track the volume changes. The air pressure pump consists of a similar arrangement where a DC-servo motor moves a piston in and out of a pressure regulator to control the air pressure output. The pumps can run under pressure- or volume-controlled modes via a Geotac custom-made computer software. Figure 4.3 illustrates the DigiFlow flow pumps.
- A computer control system. Custom-made computer software was used to communicate with the pressure control system through a network module. It acquires and records the pressure and volumetric data from the three pressure pumps. It also controls the applied pressures and maintains them at the desired values.
- An elastic wave system consisting of piezoelectric elements (i.e., disks and bender elements) imbedded in the end platens. A 20-V step electrical signal can be sent through an AC signal generator to excite the source piezoelectric element in the bottom end platen and a digital oscilloscope records the data from source. Receiver piezoelectric elements

are mounted in the top end platen. A digital oscilloscope records data from source and receiver piezoelectric elements and transfers it to the computer for further analysis.

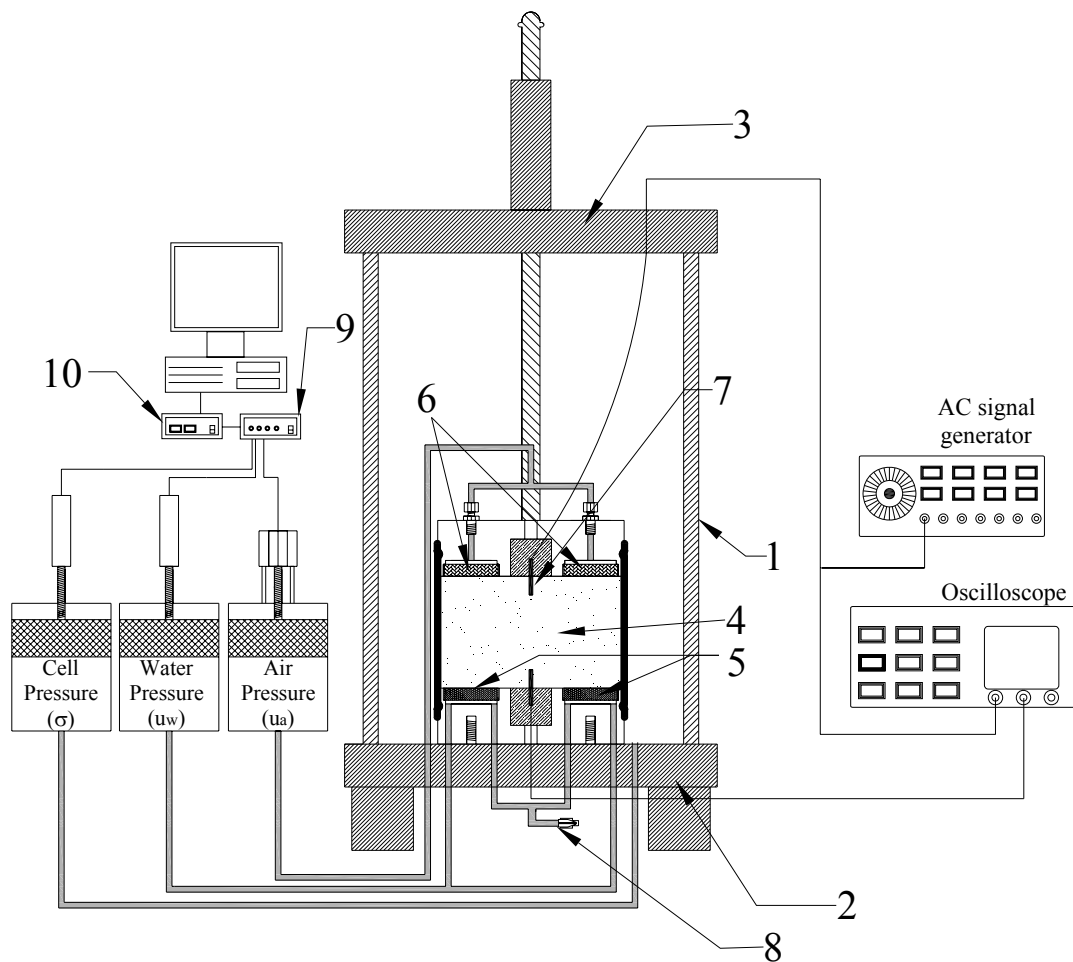
The following sub-sections provide a detailed description of the piezoelectric elements and the modified end platens of the triaxial cell.

4.3.1 Piezoelectric Elements

Piezoelectric elements are composed of crystals that lack central symmetry. These crystals become electrically polarized when they deform. They also deform under the presence of an electric field (Ikeda 1990). Piezocrystal elements are commonly used in geotechnical testing to send and receive elastic waves (e.g., Shirley and Hampton 1978, Brignoli et al. 1996, Lu et al. 2004). The type of wave generated depends on mounting configuration of the element. The most common form is cantilever bender elements mounted in platens on both ends of a triaxial or an oedometer cell to generate and receive S-waves (e.g. Blewett et al. 1999; Fam and Santamarina 1995; Cho and Santamarina 2001; Choi et al. 2004). P-waves can be generated and received using other kinds of configurations, piezoelectric disks (e.g., Fratta and Santamarina 1996, Ismail et al. 2005). When a voltage is applied to two-layer disks, a drum-like motion is generated, creating a P wave. A similar disk located on the other side of the specimen can be used to receive the waves.

The measurement of elastic waves with piezocrystal elements is simple in principle; however, electromagnetic noise may mask the arrival of elastic wave signals. Electromagnetic coupling between the source and receiver bender elements can result in an electromagnetic cross-talk which manifests as an output signal in the receiver bender elements with an early component that is quasi-simultaneous to the input signal (Lee and Santamarina 2005). This problem arises when using piezocrystal elements in conductive media such as unsaturated and saturated fine grained soils. In these soils, the voltage generated by the electromagnetic cross-talk can be

several times larger than the particle motion voltage generated by the elastic waves and therefore may result in masking the true arrival of waves. Bender elements should be properly prepared to reduce or eliminate electromagnetic cross talk. This can be achieved by shielding the element with a conductive coating which allows detecting the arrival of elastic waves by reducing electromagnetic interferences.



- | | |
|---------------------|------------------------|
| 1- Acrylic Jacket | 6- Coarse Porous Stone |
| 2- Cell Base | 7- Bender Element |
| 3- Cell Cap | 8- Flushing Valve |
| 4- Soil Specimen | 9- Data Acquisition |
| 5- HAE Porous Stone | 10- Network Module |

Figure 4.2. Schematic of the modified triaxial apparatus, test control system and wave propagation equipment.

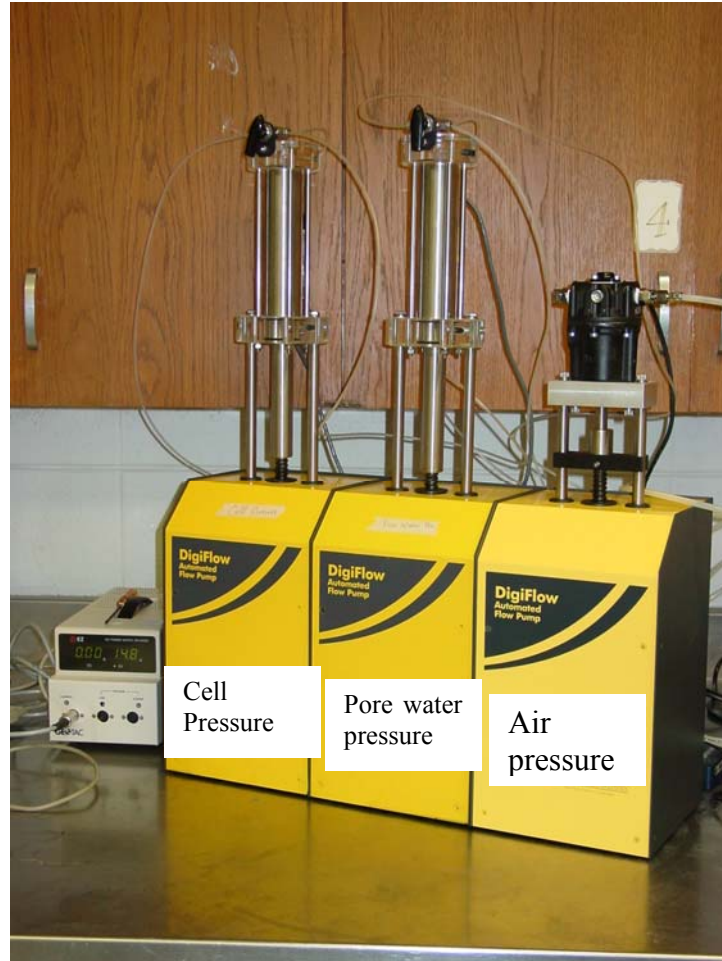


Figure 4.3. Servo-controlled air, water and cell pressure pumps

4.3.2 End Platens

In order to control the air and water pressures independently, the application of the pressures in the air and water phases have to be separated. This is achieved using high air-entry (HAE) porous disks. When saturated, a HAE disk resists flow of air while allowing water flow, as long as the difference between the air and water pressures is less than the air-entry value of the disk (Hoyos and Macari 2001, Lu and Likos 2004). The test cell bottom platen (Figure 4.4) contains two 3-bar (300-kPa) HAE disks with a diameter of 28 mm each. The pore water pressure pump connects to a cell drain line that terminates in a void space beneath the HAE

disks. The pump controls the water pressure in this void space and also tracks the volume of water entering or leaving the specimen during the experiment.

The platen also includes two counter bores to accept small metal cups which house either a bender element or a piezocrystal disk. The bender element that generates (or receives) S-waves is centered in the platen to avoid reflections from the specimen boundaries. Space limitations required that the piezocrystal disks that generate (or receive) P-waves to be placed off center. However, this was considered acceptable since reflections from the sides arrive at later times than the direct P-waves. Figure 4.5 illustrates a picture of the bottom end platen with the HAE porous stones, bender element, and piezoelectric disk.

The bender elements and piezocrystal disks are housed in metal cups and are fixed in place using epoxy resin. The metal cups are imbedded into the designated end platen counter bores and sealed using a silicon-based sealant. A metal tube attached to the cup bottom allows electric wires to pass through to the outside of the test cell where they are connected to the signal generator, oscilloscope, and data acquisition module. Figure 4.6 depicts two bender elements fixed in the metal cups and ready to be installed in the end platen.

Although saturated HAE disks prevent the passage of air through them, due to the long times required for testing unsaturated soils, some air can still diffuse through water to form bubbles under the HAE entry disks. Such bubbles can cause errors in pore water pressure measurements; therefore, a mechanism has to be implemented to periodically flush these bubbles (Fredlund and Rahardjo 1993). Figure 4.4b illustrates the flushing drain lines in the bottom plate and platen. When the flushing valve is closed, the water in the system is pressurized to produce the required pore water pressure, however, when it is opened, water flows through the drain lines and out the flushing valve, removing air bubbles from beneath the HAE disks.

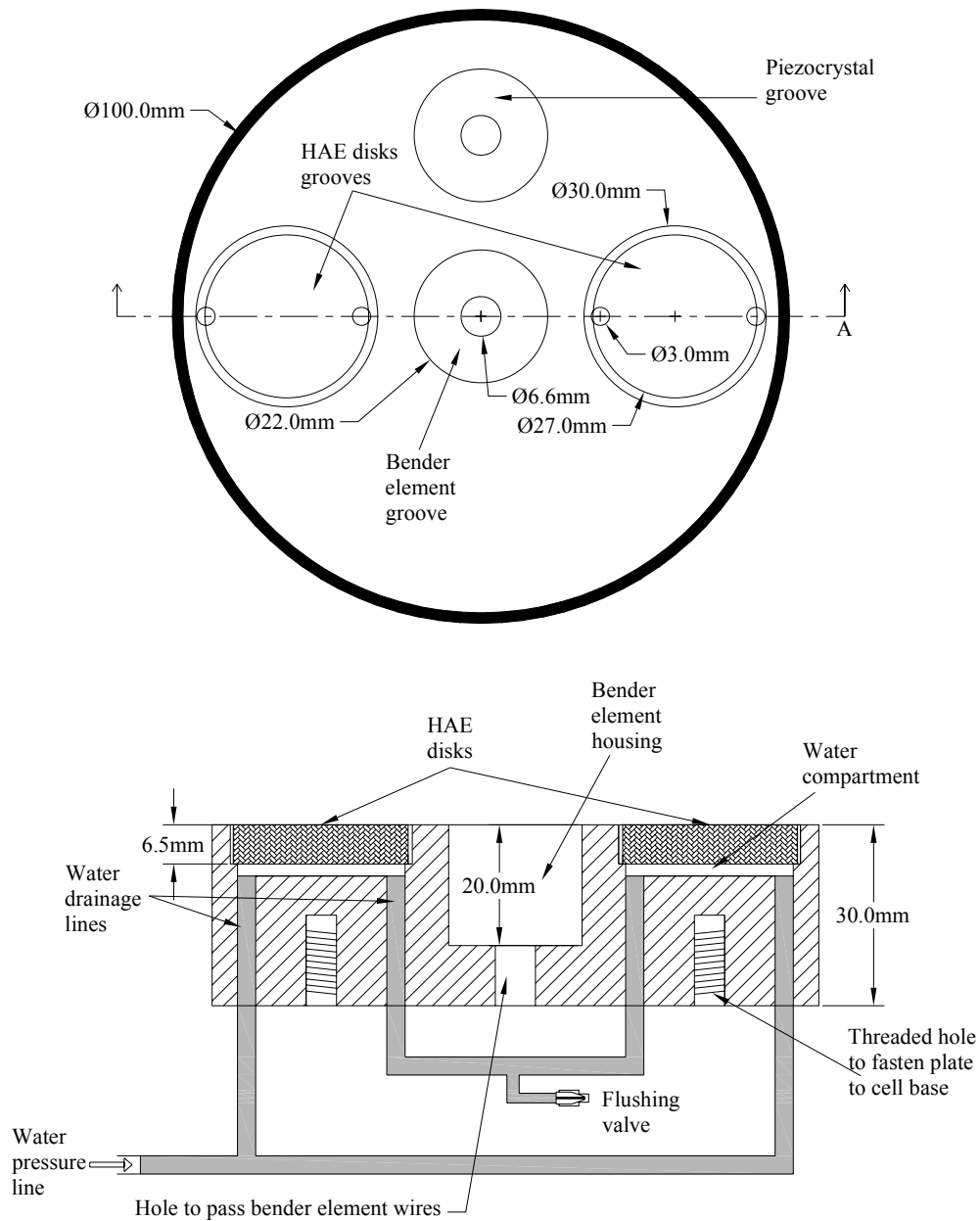


Figure 4.4. Details of the bottom end platen: plan and cut views.

The air pressure pump, used to control the pore air pressure, connects to a cell drain line that terminates at two coarse porous stones imbedded in the test cell top platen. The top platen also contains a bender element and a piezocrystal arranged with a similar geometry to the ones in the bottom end platen.



Figure 4.5. Bottom end platen



Figure 4.6. A pair of bender elements fixed in metal cups and ready to be implemented in the end platens

4.3.3 Distance between Elastic Wave Sources and Receivers

Due to the bending motion of bender elements, S-waves are generated in the direction of the bender element plane; however, two other P-wave side lobes are generated in a direction perpendicular to that plane (Lee and Santamarina 2005). Since the travel velocity for P-waves is always higher than S-waves, the arrival of the P-waves reflected from the specimen boundaries may precede the direct S-wave arrival. The end platen dimensions must be designed so that the travel time for the direct S-wave is shorter than the travel time for the reflected P-wave (Lee and Santamarina 2005):

$$R = \frac{L}{2\sqrt{1-2\nu}} \quad (4.8)$$

where R is the minimum radius of the specimen for a tip-to-tip distance L between bender elements (Figure 4.7). For a typical Poisson's ratio ν of 0.15, the minimum required ratio of R to L is about 0.6. The specimen has a height of 50 mm and the length of the bender elements is about 10 mm, the tip to tip distance is about 30 mm. Therefore, the R to L ratio is about 1.67 which is sufficient to ensure the first arrival of S-waves for Poisson's ratios up to 0.45. Having this large radius gives the flexibility for increasing the specimen height while satisfying an acceptable range of R/L ratio, and therefore permitting the possibility of testing specimens of different heights and under deviatoric stress conditions.

4.3.4 Saturation of the High Air Entry Disks

In order for the high air entry disks to operate properly, allowing the flow of water while preventing the flow of air through them, they must be fully saturated. A procedure suggested by Fredlund (1973) was used to saturate the high air entry disks in the bottom end platen of the triaxial cell. To saturate the porous disks the triaxial cell was assembled and filled with water to a height of about 25 mm above the HAE disks. 600 kPa of air pressure was applied to the cell for

about an hour to induce flow through the disks and drive the air bubbles to the bottom compartment. The bubbles collected beneath the disks were periodically flushed using the flushing valve. Then, all the valves under the disks were closed so that the water over, under, and inside the disks takes the value of the applied (600 kPa) pressure. During this process, the air inside the HAE disks dissolves in water. The pressure should be applied for approximately one hour, after which the valves under the porous stones were opened again for about 10 minutes to allow the water to flow into the bottom compartment. The bubbles are then flushed using the flushing valve where the water was allowed to flow under the disks with a pressure of 2 kPa. This procedure was repeated for at least six times (usually 10 times) to ensure the proper saturation of the HAE disks. After this procedure, the HAE disks were submerged in water until an experiment was ready to start.

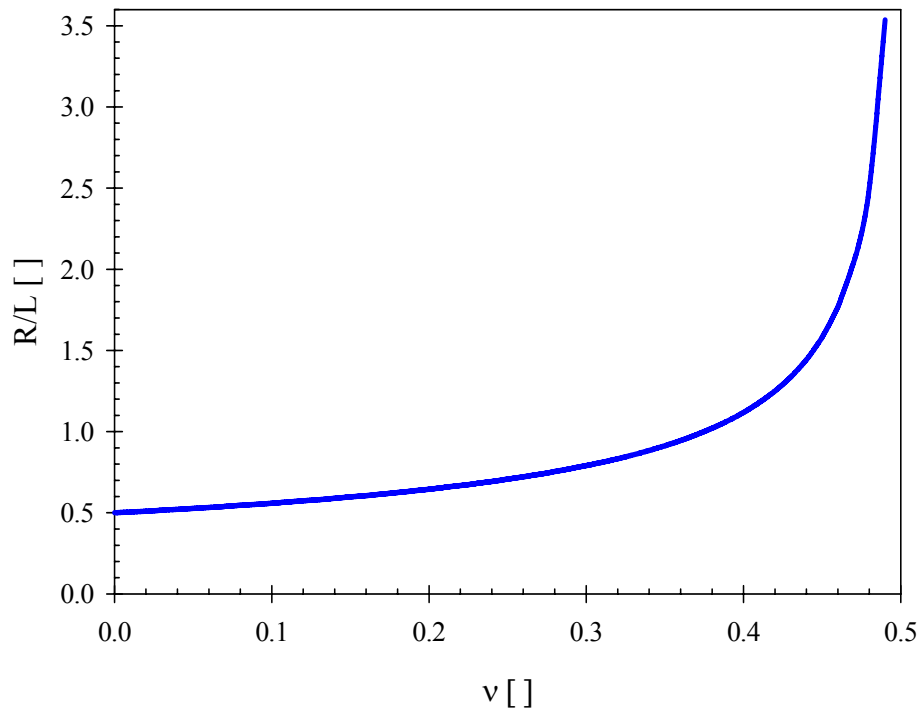


Figure 4.7. Minimum radius to “tip to tip” distance ratio required to ensure the first arrival of the S-wave.

4.4 Specimen Preparation and Test Procedure

Three unsaturated soil specimens were tested under drained isotropic compression conditions. During each test, the matric suction was changed while the net stress was kept constant. The axis translation technique was used to elevate the reference pore water pressure value to avoid cavitation resulting from negative pore water pressure values (Hilf 1956). During the experiment, pore water pressure was held constant while air pressure and cell pressure were increased by the same amount to increase the matric suction and keep the net stress constant. Wave velocity measurements were performed at different time intervals during each matric suction increment.

Remolded soil specimens were prepared by first mixing the soil and water to achieve the desired initial water content. The soil was placed inside the mold in multiple layers (typically three layers) while applying uniform tamping to each layer to achieve the desired density. The top end platen was then placed on the top of the specimen and was sealed using the rubber membrane. The split mold was then removed and vacuum was applied to the specimen pore space to prevent disturbance. The test cell was then assembled and the air, water, and cell pressure lines were connected to their pumps and the cell was filled with the confining fluid. Silicone oil, which is non-conductive, was used as the confining fluid to prevent short circuits between the electrical wires emanating from the top platen.

After filling the cell with silicone oil, the cell, water and air pressures were applied to the specimen to achieve an initial equilibrium condition. Once initial equilibrium was established, pressures were adjusted to achieve equilibrium at a series of matric suction values. For each change or step, a hydrostatic pressure rate of 6 kPa/hr was used based on recommendations by Porras-Ortiz (2004) to ensure a minimum disturbance of the capillary menisci. Water drained from the specimen pore space during each step, and was monitored to determine when

equilibrium was reached. The equilibrium was reached when the amount of water leaving the pore space was less than 0.04% per day as recommended by Aversa and Nicotera (2002). Wave velocity measurements were acquired to monitor the change in soil stiffness as the capillary forces were developing in the soil specimen.

After equilibrium was attained on a step, a new suction increment was applied while the net stress was held constant. This procedure was repeated for different suction increments until the maximum desired suction value was reached.

4.5 Experimental Results

A series of experiments was conducted to verify the operation of the apparatus. Three soil specimens were tested under drained isotropic compression. P and S wave traces were captured as the matric suction was increased. The specimens were prepared by mixing different proportions of a river silty soil and F-75 Ottawa sand. A description of the soil properties, specimen composition, initial density and moisture content are listed in Table 4.1. Grain size distribution for the various soils is presented in Figure 4.8. The net stress was held constant at 20 kPa throughout the experiments while the matric suction was incrementally increased. The net stress was kept at a value that corresponds to a soil depth of approximately 1.5 m. This low net stress value was selected to emphasize the effect of change in matric suction on the wave propagation velocity.

Table 4.1. Summary of Soil Properties

Soil	River Silt (%)	F-75 Ottawa Sand (%)	Plastic Limit (%)	Liquid Limit (%)	Plasticity Index (%)	D ₅₀ (mm)	Specific Gravity	Initial Void Ratio	Initial Dry Density (kg/m ³)	Initial Water Content (%)
1	50	50	14.8	17.5	2.7	0.30	2.67	0.64	1,628	19.07
2	67	33	17.1	21.2	4.1	0.21	2.67	0.54	1,735	13.57
3	100	0	22.1	28.7	6.6	0.21	2.66	0.75	1,516	26.10

4.5.1 Wave Velocity Results

Typical recorded traces in the time and frequency domain plots are shown in Figures 4.9 and 4.10. The power spectra density shows frequency contents that vary between 1.5 to 7.0 kHz that for typical wave velocities of 200 to 500 m/s yields wavelengths that range from 0.07 to 0.13 m. These results are compared to the separation between sources and receivers (0.03 m) and show that the wave propagation measurements are in the near-field regime (i.e., wavelength $\lambda > L/4$; Sanchez Salinero et al. 1986; Sawangsuriya et al. 2006). Propagation in the near field regime complicates the evaluation of the S-wave arrival because it is masked by the arrival of the P-wave even when the generated particle motion is perpendicular to the direction of wave propagation.

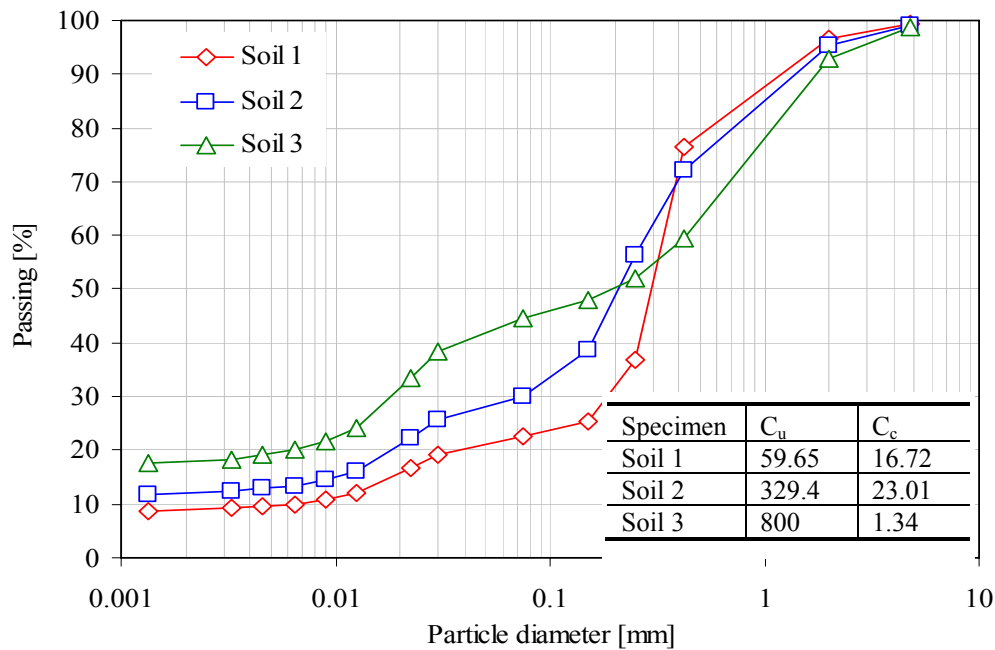


Figure 4.8. Grain size distribution of tested soils

The P and S-wave velocities were obtained by measuring the travel times between the source and receiver piezoelectric elements. The first arrival time of the signal to the disk piezoelectric element (P-waves) can be manually obtained from the first break in the wave traces

(Figure 4.9a). However, in the case of S-waves, due to the presence of near field effects and to some difficulties in the selection of an accurate travel distance, other methods were used to facilitate the determination of the wave velocities (for example, Lee and Santamarina 2005). Multiple reflections detected by the same element provide a simple and robust method to accurately determine travel times. When multiple reflections are captured by a single bender element, the time shift between the first arrival and the reflection can be determined by a simple cross correlation between the two parts of the trace. Then using two times the distance between the two end platens, the velocity can then be calculated. Figure 4.9b shows examples of S-wave traces where the first arrival and the reflection are illustrated.

Figure 4.10 shows that the frequency content increases with higher matric suction values. These trends are an indication of the effect of suction on the stiffness of the soil by generating higher frequency contents. Furthermore, as the matric suction increases the amplitude of the power spectra density decreases while the range of frequency becomes wider (Figure 4.10b). These phenomena are caused by two different effects: as matric suction increases, (a) the increase in soil stiffness reduces the deformation triggered by the source bender elements and therefore reducing the signal obtained at the receiver bender elements, and (b) the bender element resonant frequency increases with increasing soil stiffness and the bender element can better reproduce the wide range of frequencies present in the step input signal.

After each increment of suction, wave velocity measurements were acquired at different time intervals throughout the stabilization process. Figure 4.11 illustrates a typical change in P and S-wave velocities versus time after a matric suction increment is applied. A rapid increase in the wave velocity is observed right after the suction increment is applied, and then the rate of change in velocity and water volume decreases with time until little or no change is observed when equilibrium is reached.

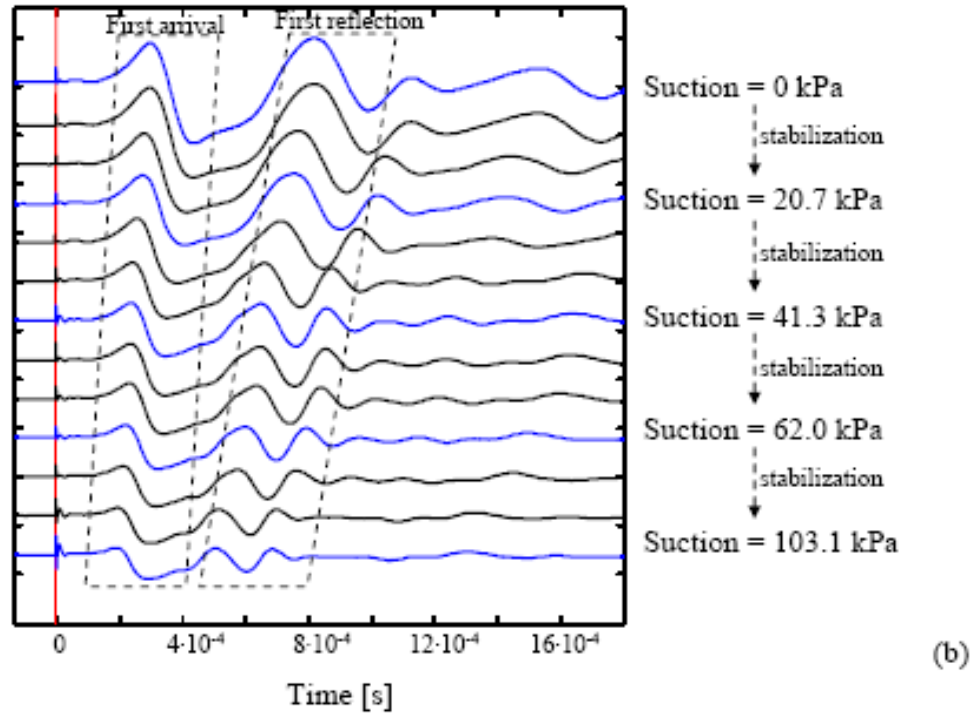
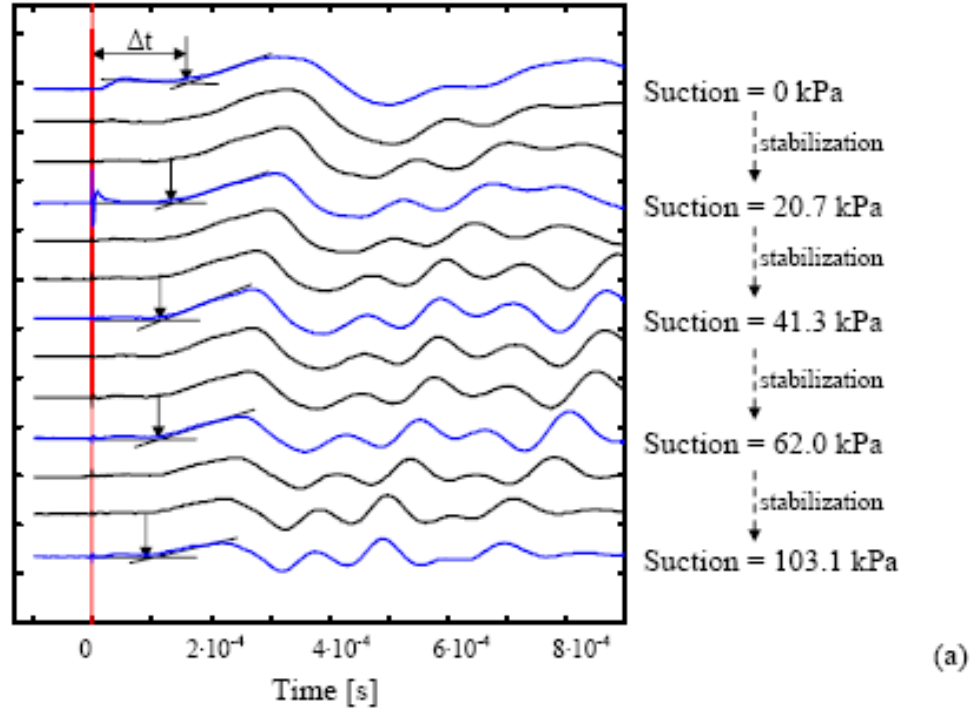


Figure 4.9. Wave signatures obtained for specimen 3 at increasing matric suction values. (a) P-wave traces. (b) S-wave traces. The thick lines correspond to the wave traces at stabilized matric suction values while the thin lines indicate wave traces during the stabilization process.

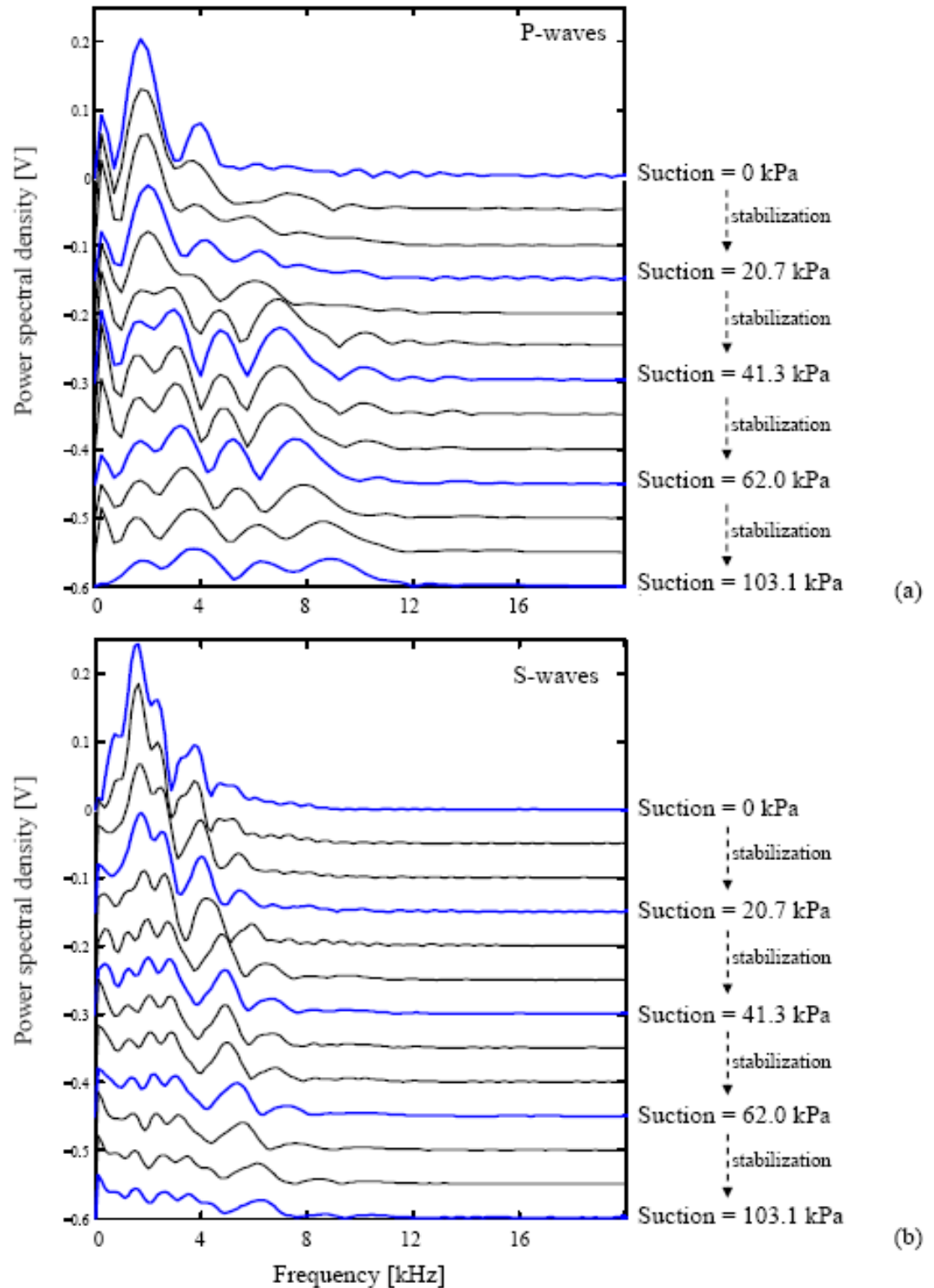


Figure 4.10. Wave power spectra density calculated for specimen 3 at increasing matric suction values. (a) P-wave traces. (b) S-wave traces. The thick lines correspond to the wave power spectra at stabilized matric suction values while the thin lines indicate wave power spectra during the stabilization process.

To completely illustrate the effect of matric suction on the elastic wave velocity, the final P and S-wave velocities obtained at the end of each loading step are shown in Figure 4.12 and 4.13. It can be seen that increasing the matric suction causes an increase in both P and S-wave velocities. This is a result of the increase in bulk and shear moduli caused by the increase in matric suction that increases the interparticle forces. Furthermore, higher P and S-wave velocities were obtained with specimens with greater fine content. This observation seems to support resonant column results obtained by Qian et al. (1991). Qian and coworkers tested angular sand specimens with varying degree of saturation and fine content. They documented an increase in the normalized small-strain shear modulus G_{ske} with increasing fine content in sandy specimens.

4.5.2 Evaluation of Results

To evaluate the wave velocity results obtained during the three experiments, the results were compared to wave velocity values calculated using a theoretical wave velocity model. For this purpose, a model that considers the effect of matric suction on the shear wave velocity was used to predict V_s values to be compared to the experimental data. Santamarina et al. (2005) proposed the following semi-empirical relation:

$$V_s \approx V_s^{S_r=1.0} \cdot \left[1 + S_r \frac{u_a - u_w}{\sigma'} \right]^\beta \quad (4.9)$$

where $V_s^{S_r=1.0}$ is the shear wave velocity at saturation, σ' is the mean effective stress that corresponds to the saturated soil, S_r is the degree of saturation, and β is an experimentally determined exponent that depends on the soil type (i.e., type of contact) and the stress history. The term in parenthesis shows the interplay between matric suction and degree of saturation. When soils dry, matric suction increases while saturation decreases capping the maximum interparticle forces created by the water menisci (Vinale et al. 1999).

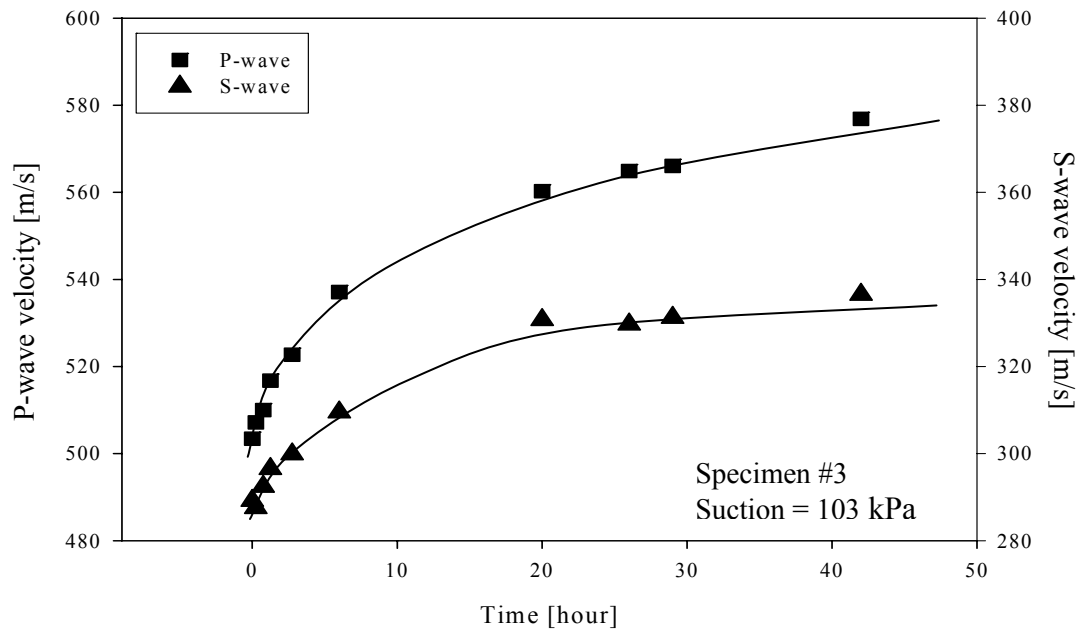


Figure 4.11. Typical wave velocity response after the application of a matrix suction increment (the lines just indicate trends).

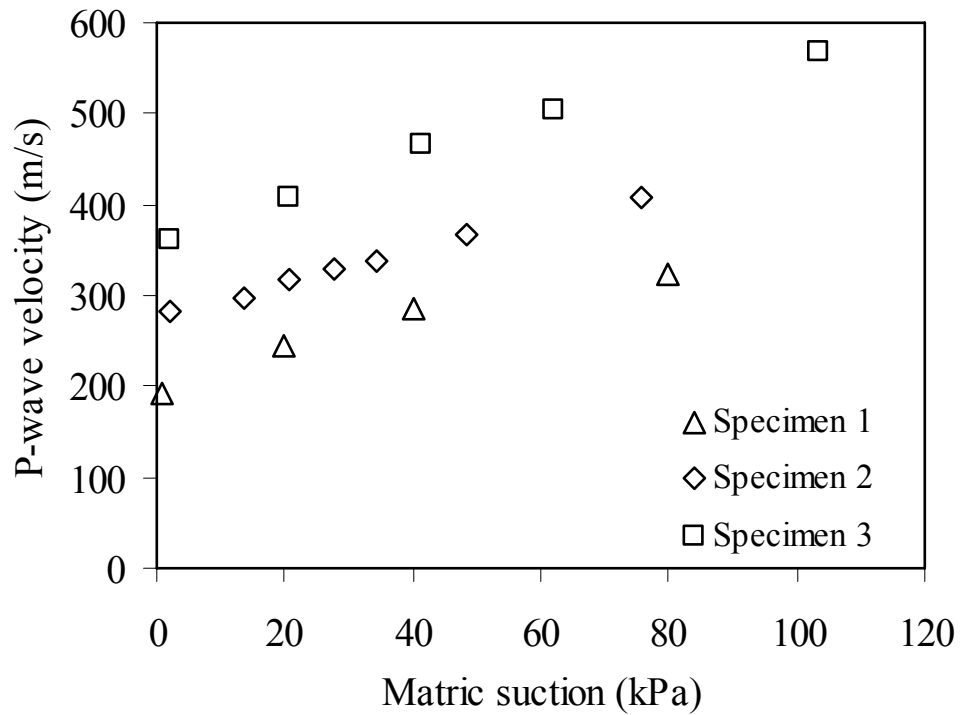


Figure 4.12. Measured P-wave velocity at increasing matrix suction values.

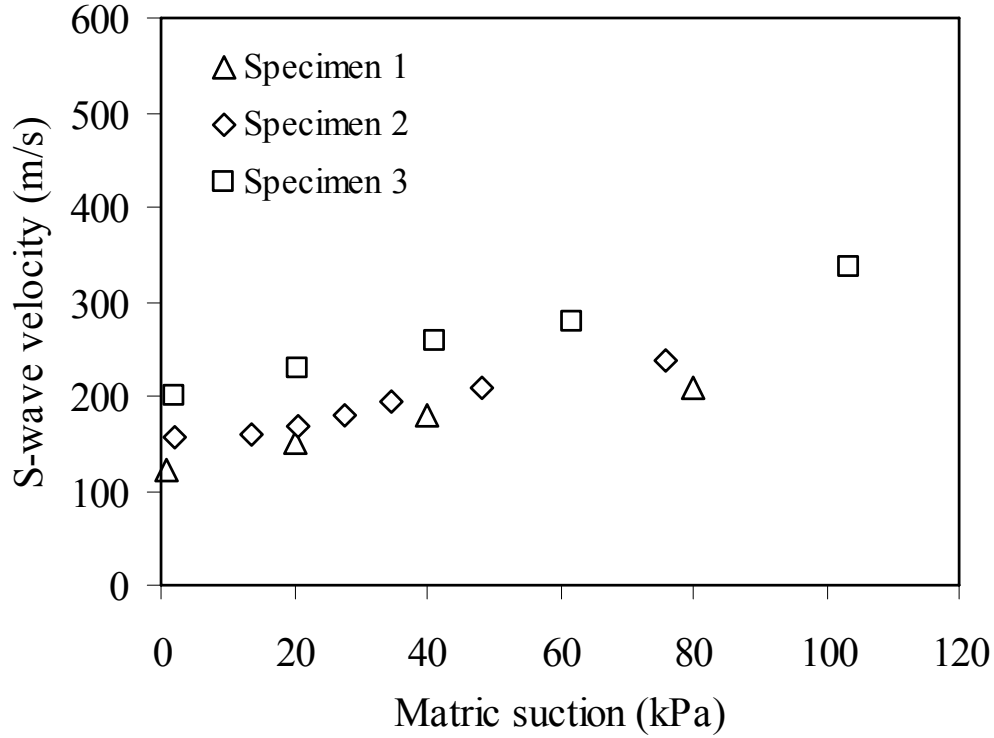


Figure 4.13. Measured S-wave velocities at increasing matric suction values

In order to use the wave velocity model described by Equation 4.9 the knowledge of S-wave velocity at saturation is required; therefore, another model was used to estimate that velocity. Santamarina et al. 2001 presented a velocity-stress relationship for isotropically loaded soils based on the fact that the state of stiffness of granular materials (and therefore V_s) is determined by the state of the effective stress (Equation 4.10).

$$V_s^{S_r=1.0} = \alpha \left(\frac{\sigma - u_a}{\sigma_{ref}} \right)^\beta \quad (4.10)$$

where $\sigma_{ref} = 1$ kPa is the reference stress α and β are experimentally determined parameters (α is the velocity of the medium subjected to 1 kPa confinement). Santamarina et al. (2001) also presented a plot of α and β values for various granular materials (Figure 4.14), and suggested an empirical relationship between the two parameters:

$$\beta \approx 0.36 - \frac{\alpha}{700} \quad (4.11)$$

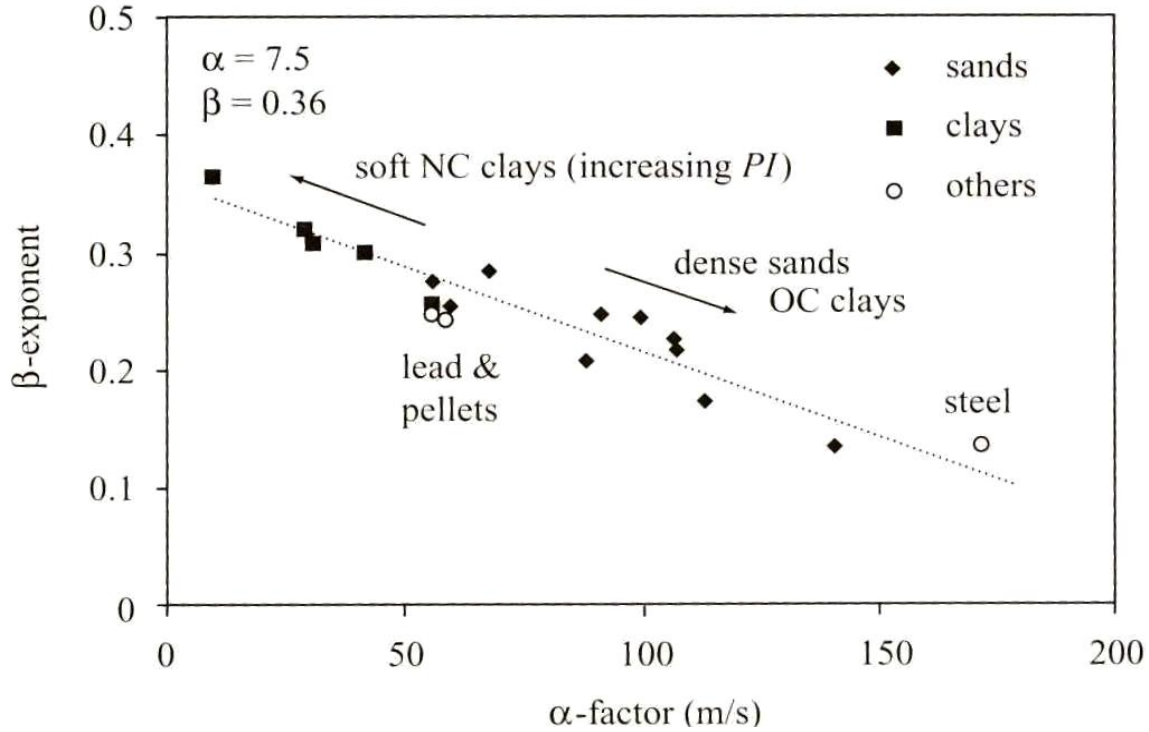


Figure 4.14. Typical values for α and β parameters for different types of granular materials (Santamarina et al. 2001).

Incorporating the value of the S-wave velocity at saturation from Equation 4.10 into the general S-wave velocity model in Equation 4.9 yields a new equation for the S-wave velocity in unsaturated soils considering the effect of matric suction:

$$V_s \approx \alpha \left(\frac{\sigma - u_a}{\sigma_{\text{ref}}} \right)^{\beta} \cdot \left[1 + S_r \cdot \frac{u_a - u_w}{\sigma - u_a} \right]^{\beta_{\text{unsat}}} \quad (4.12)$$

where $\sigma_{\text{ref}} = 1$ kPa is the reference stress. It should be noted that Equation 4.12 now contains two different β exponents. Therefore, the second exponent was denoted β_{unsat} . The exponents β and β_{unsat} must be different not only in the magnitude (the matric suction term is normalized with respect to the net stresses) but in the physical processes they represent. That is, the change of S-wave velocity may be controlled by different mechanisms along each of the two

stress state parameters: net stresses and matric suction. While first exponent β represents the nature of the contact between particles (Fratta and Santamarina 2003), the second exponent β_{unsat} includes not only the type of contact but also the effect of the menisci on the contact. The effect of menisci may include particle migration, osmotic suction, salt deposition, and cementation. These effects are expected to be most important at low degree of saturation where salt precipitation cements the contacts and contributes to the skeleton shear stiffness. Further research is necessary to confirm this hypothesis.

Equation 4.12 was then used to evaluate the changes in S-wave velocity induced by increasing matric suction in the three tested soils. Moreover, knowing that the ratio of the P-wave to the S-wave velocity is mainly controlled by Poisson's ratio, Equation 4.13 was used to calculate the P-wave velocity. The calculated P-wave velocities were also compared to the experimentally determined P-wave velocity values.

$$V_p = \sqrt{\frac{2-2\nu}{1-2\nu}} \cdot V_s \quad (4.13)$$

The results and fitting parameters for V_s and V_p are plotted and summarized in Figure 4.15. It can be seen that the unsaturated exponent β_{unsat} was in all cases much larger than the measured exponent β obtained for dry or saturated soils (Santamarina et al. 2001). These results indicate a dominant effect of matric suction on shear stiffness at low net stresses. The value of the unsaturated exponent β_{unsat} is expected to decrease with increasing net stresses.

To confirm these results, resonant column data presented in Vinale et al. (1999) were fitted with Equation 4.12. The small-strain shear modulus data were obtained for compacted Metrano silty sand specimens prepared at optimum and wet of optimum water contents. The shear modulus data were converted to shear wave velocity by assuming constant mass density $\rho=1900 \text{ kg/m}^3$. The data and the fitted model results are presented in Figure 4.16. The data were

only fitted in the low matric suction range to allow the comparison with the results documented in this study. Consistently, with the model parameters obtained with this new test setup, the data from Vinale et al. (1999) yields values of β_{unsat} greater than published values of β . Furthermore, the new Equation 4.12 helps represent the changes of S-wave velocity with both matric suction and net stresses, the two controlling stresses in the problem of wave propagation in unsaturated soils. It is also important to indicate that as suction increases, the trend shown by Equation 4.12 would fail to match the experimental results presented in Figure 4.16a. This inconsistency can be justified by the reduction in degree of saturation that caps the gain in stiffness in the soil. This effect is not captured in Equation 4.12 as the degree of saturation at high matric suction values is not available.

4.6 Chapter Summary

A new apparatus for the measurement of P and S-wave velocities in unsaturated soil specimens under controlled net stress and matric suction was presented in this chapter. Matric suction-control experiments on three different soils were performed to evaluate the system performance. Elastic wave velocities in unsaturated soil specimens were monitored while matric suction was increased. The stabilization of capillary pressures was observed through monitoring the rate of change of wave velocities after suction increments were applied. It was realized that both P and S-wave velocities increase as the matric suction increases following semi-empirical models presented in the literature. These models also help to better understand the interaction between the different elastic components of the wave velocity equations in unsaturated soils.

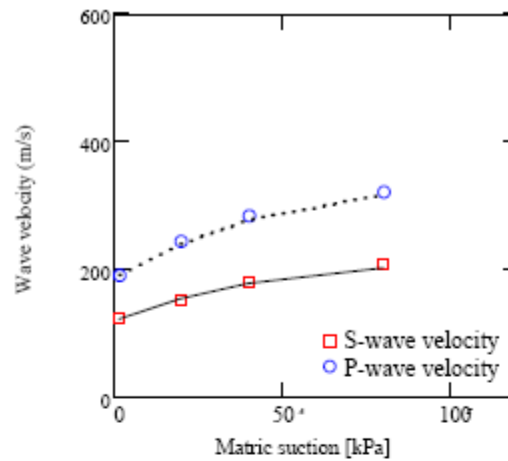
Specimen 1:

$$\alpha = 50 \text{ m/s}$$

$$\beta = 0.29$$

$$\beta_{\text{unsat}} = 0.45$$

$$\nu = 0.15$$

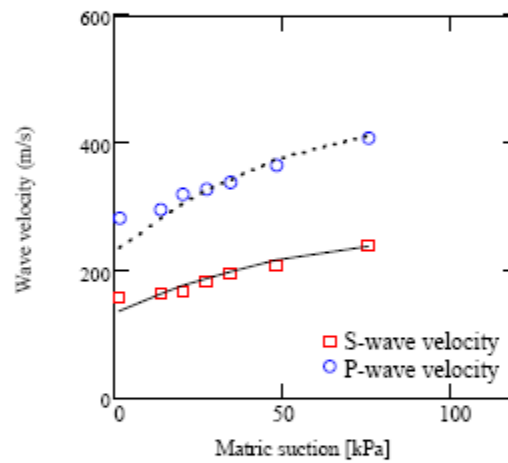
**Specimen 2:**

$$\alpha = 56 \text{ m/s}$$

$$\beta = 0.28$$

$$\beta_{\text{unsat}} = 0.58$$

$$\nu = 0.25$$

**Specimen 3:**

$$\alpha = 100 \text{ m/s}$$

$$\beta = 0.22$$

$$\beta_{\text{unsat}} = 0.32$$

$$\nu = 0.25$$

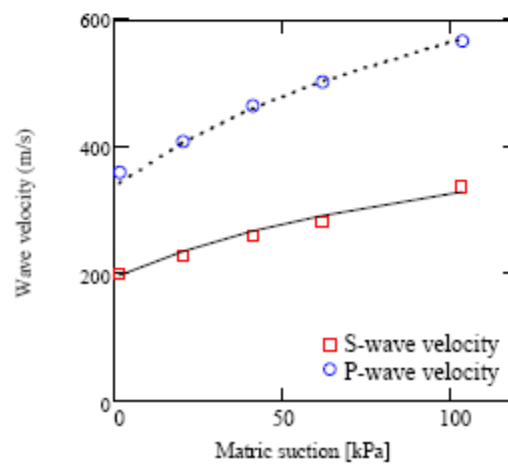
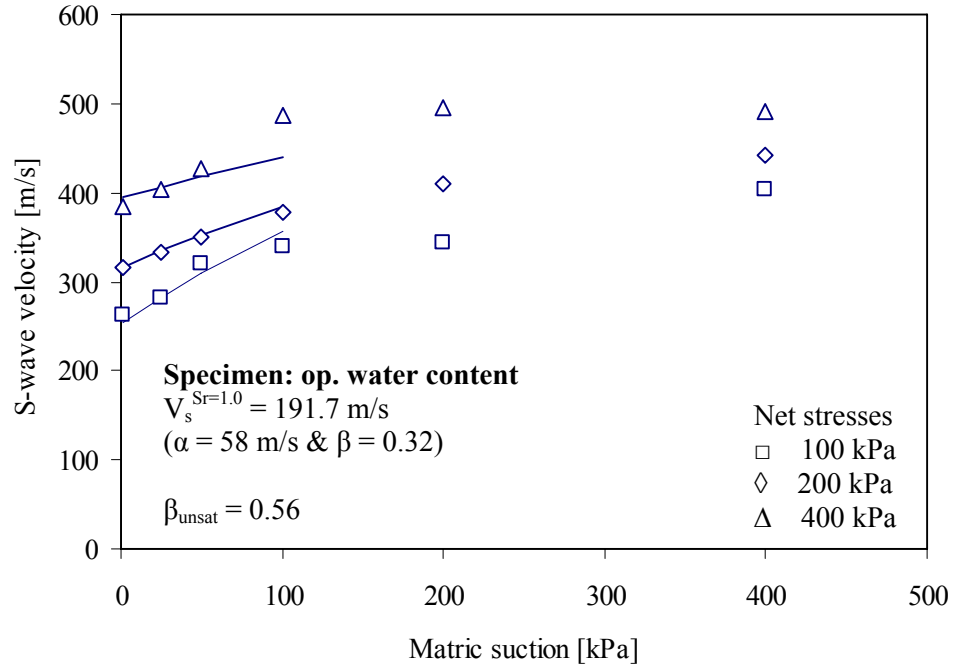
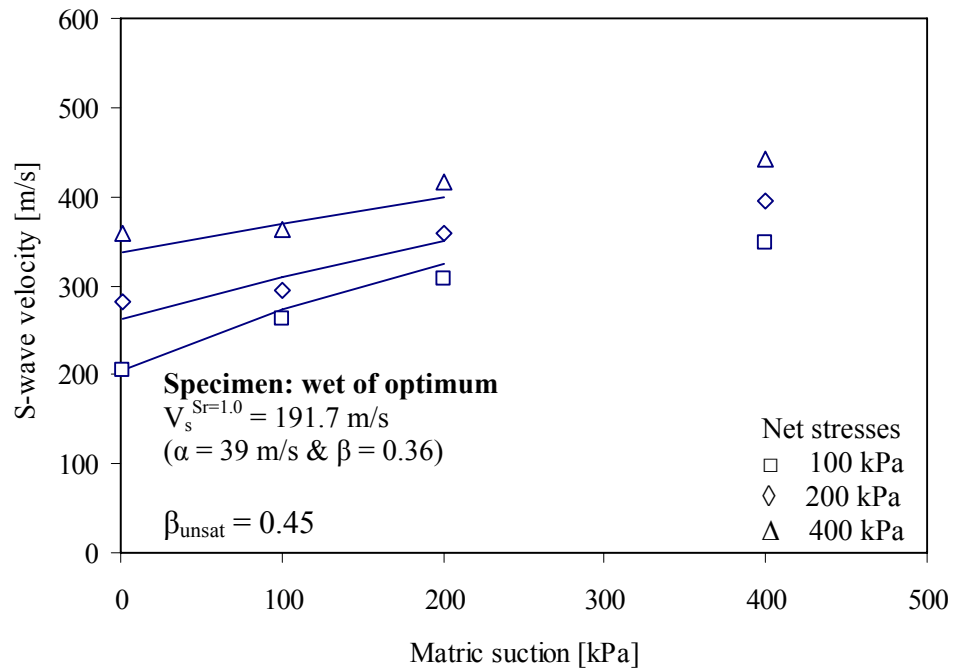


Figure 4.15. Modeling of wave velocities and Poisson's ratios for soils under increasing matric suction levels.



(a)



(b)

Figure 4.16. Calculated S-wave velocity from shear stiffness using resonant column measurements in compacted silty sand specimens versus matric suction and net stresses data after Vinale et al. 1999): (a) Compacted at optimal water content and (b) compacted wet of optimum. The data was fitted with Equation 4.12. Model parameters are shown in the figures.

CHAPTER 5

THE EFFECT OF FINE PARTICLE MIGRATION ON THE SMALL STRAIN STIFFNESS OF UNSATURATED SOILS

5.1 Introduction

It was demonstrated in the previous two Chapters that the small strain shear stiffness of near surface soils is to a great extent controlled by the soils' degree of saturation. As the degree of saturation decreases, matric suction increases, causing an increase in the interparticle forces, consequently, the overall soil shear stiffness increases. In most experimental studies on unsaturated soils, the pore fluid was either tap water or deionized deaired water. However, natural water could in many instances have different properties than the water used in these experiments. Usually, fine materials such as clay or silt particles are transported by water, and the presence of such fines in a granular medium can alter the pore fluid properties. Therefore, it is important to study the effects of the presence of fines on the overall properties of unsaturated soils such as small strain shear, constraint, and bulk moduli.

This Chapter presents an experimental study on the effects of the presence of fines (silt and clay) in the pore fluid on the soils' small strain shear, bulk and constraint moduli by using elastic P and S waves. The results were also confirmed by studying the effect of the fines on the properties of the pore fluid using synchrotron x-ray micro-computed tomography (SMCT). This Chapter includes a description of the apparatus used to carry out the experiments, description of the specimens and the results of the experimental study. It also presents a description of the x-ray CT system, the acquired images, image analysis results and a discussion of the results.

5.1.1 Drying in Unsaturated Soils

During drying, unsaturated soils go through different stages at different degrees of saturation. Starting from a fully saturated condition, drainage or evaporation causes the outside menisci to be pulled inward increasing the suction pressure (Figure 5.1a). Although a very small change in the water content takes place at this stage, it has an important global effect on the soil mass which is basically still saturated away from the boundaries (Santamarina et al. 2001). As drainage or evaporation continues, the air pressure causes the air phase to invade the water pore structure. The pressure at that point is called the air-entry value or the bubbling pressure. It is a function of the sizes of the pore. Smaller particles are usually associated with smaller pore space which makes it more difficult for air to break into them; therefore, fine soils usually have higher air entry values. Most soils have air entry values at a degree of saturation between 0.9 and 1.

After air entry, the water phase enters what is known as the “funicular regime” (Leverson and Lohnes, 1995). At this stage, although the soil is unsaturated, water still exists in a continuous phase. In a funicular regime, many pores are still filled with water; however, some are filled with air (Mitarai and Nori, 2006). Suction in this stage increases in a quasi-linear trend while drying proceeds. Water pressure is homogenized throughout the mass through pressure diffusion through the continuous water body (Santamarina et al. 2001). A schematic illustration of the funicular regime is depicted in Figure 5.1b.

The next stage commences when the water phase becomes disconnected which is known as the “pendular regime” where water forms bridges at the contact points between grains (Figure 5.1c - Mitarai and Nori, 2006). Suction at this stage acts at the contact level and exhibits a significant increase because of the small radii of the menisci (see Equation 2.14). Suction changes within a meniscus are felt by the other menisci through the corresponding change in vapor pressure in a slow homogenization process (Santamarina et al. 2001).

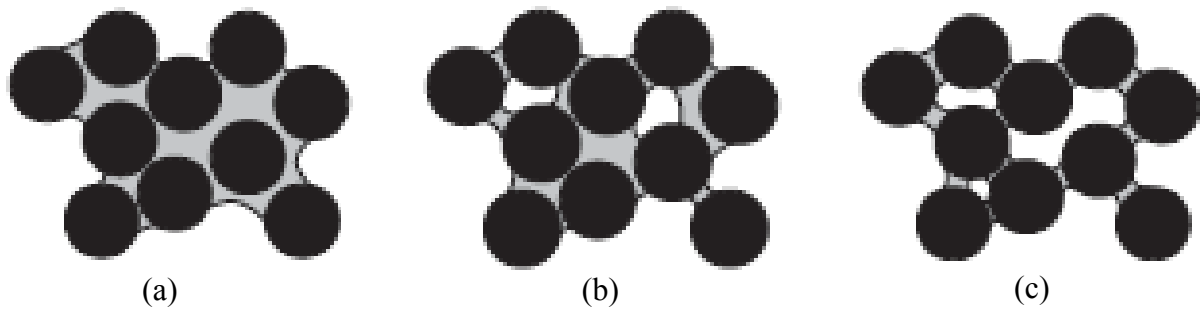


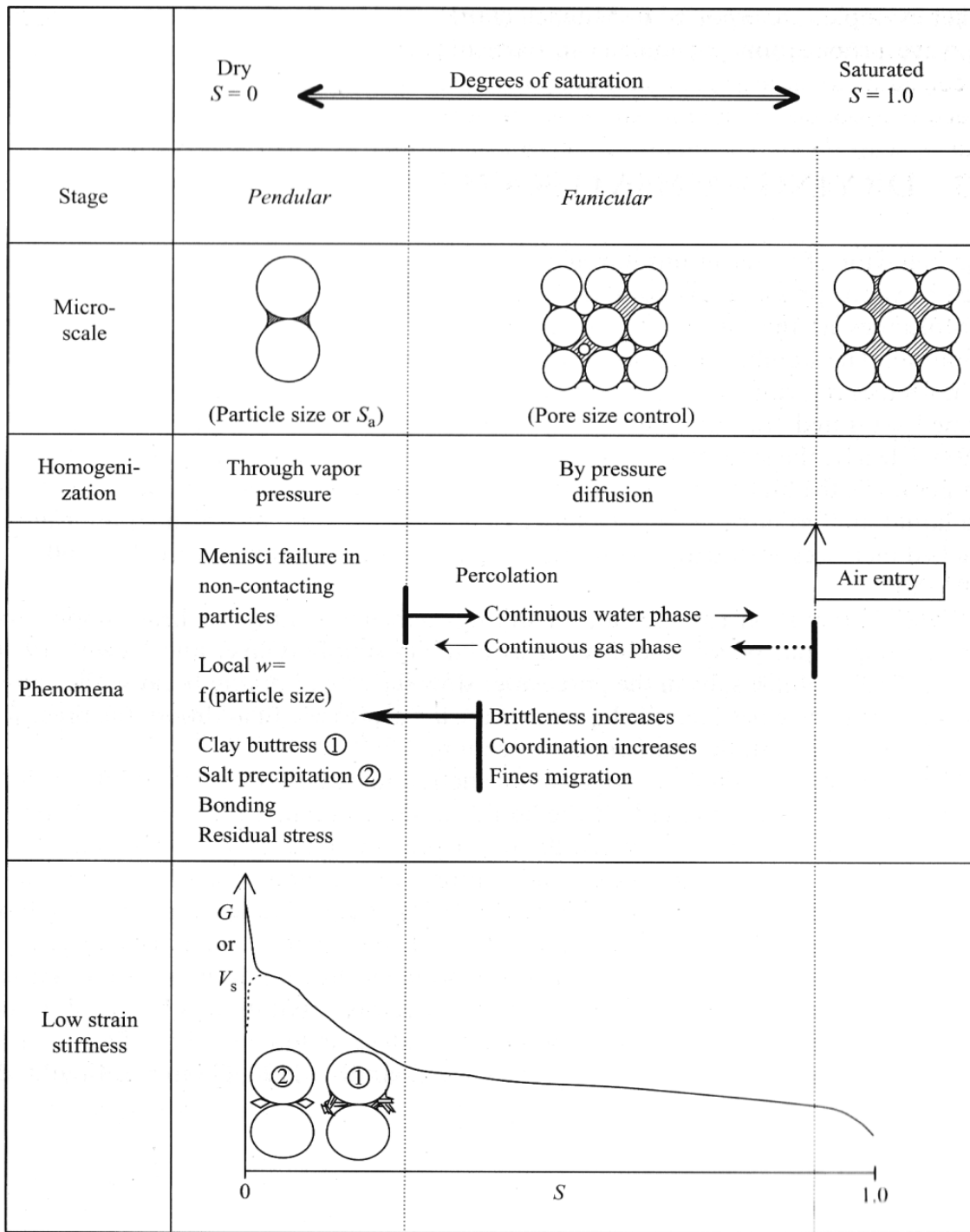
Figure 5.1 Schematic plots of the different saturation stages, (a) fully saturated, (b) funicular regime, (c) pendular regime (after Mitarai and Nori, 2006).

As the soil dries during the pendular regime, fine particles migrate to contact areas and form buttresses between larger particles, causing an increase in the stiffness of the granular skeleton formed by the coarser grains. The ionic concentration also increases with further drying causing salt precipitation at the particle contacts and contributing to the stiffness of the skeleton (Santamarina et al. 2001). Figure 5.2 depicts the different drying stages of unsaturated soils and the different phenomena associated with each stage.

5.2 Experimental Work

5.2.1 Apparatus Description

In order to investigate the effect of capillarity and the presence of fine particles in the pore fluid on the small strain stiffness and elastic waves velocity through soils several drying experiments were conducted using a specially designed drying cell. The drying cell consists of an acrylic cylinder and two end platens that accommodated the elastic waves' sources and receivers. Eight cotton threadlike wicks pass through the wall of the cell to accelerate the drying process. Figure 5.3 shows a schematic of the drying cell showing the locations of the cotton threads and the bender elements.



Notation: G is shear modulus, V_s is shear wave velocity

Figure 5.2. Stages of unsaturated conditions and related phenomena (Cho and Santamarina, 2001)

Each end platen houses two types of bender elements: a vertical (cantilever) bender element located at the center of the end platen to generate and receive S-waves, and a horizontal

bender element located at an eccentric position to send and receive P-waves. It is important to have the S-wave source and receiver at the center of the end platen to ensure that the direct S-wave arrival before the P-wave side lobes reflected from the cell walls. The end platens also have grooves to run the electrical wires. After the bender elements were installed in place, all the grooves were sealed with epoxy resin to prevent water leakage. Figure 5.4 depicts a schematic of the end platen. A dead weight of 500 grams was placed on top of the top platen to provide a surcharge of 1.7 kPa. Figure 5.5 illustrates pictures of the ends platens and the drying cell.

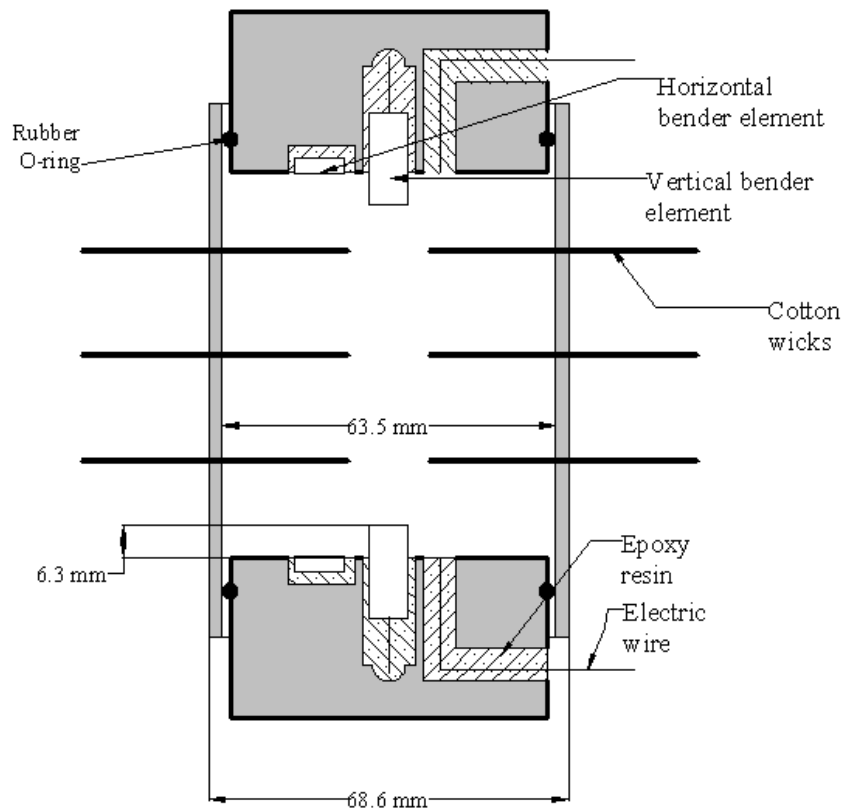


Figure 5.3. Schematic plot of the drying cell

The elastic wave velocity system used to generate and receive the elastic waves throughout the drying experiments consisted of a digital oscilloscope (Agilent 54624A, 100

MHz) to record the generated and received elastic wave traces, an AC signal generator (Krohn Hite 1200A) to generate the signals to excite the source bender element, and a filter/amplifier (Krohn Hite Model 1944) to reduce noise and amplify the signals from the receiver bender element. A 20V step signal was used to excite the source bender element.

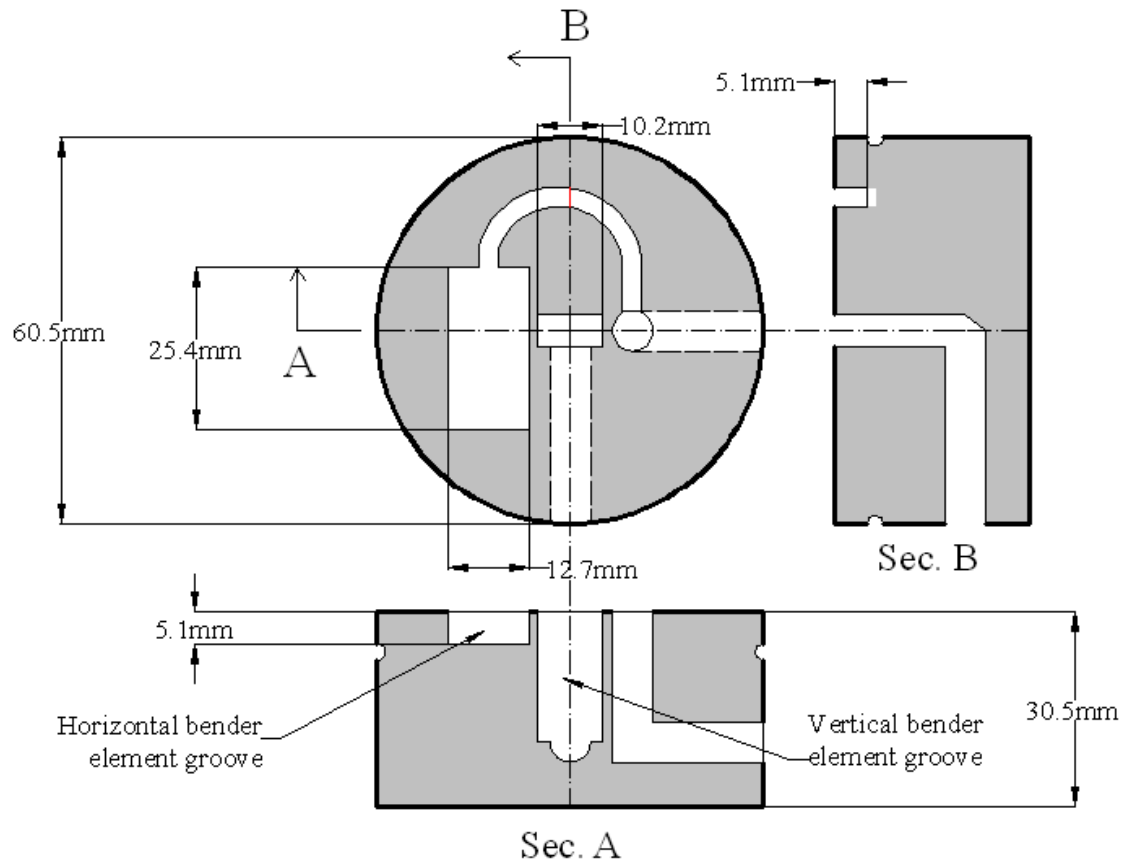


Figure 5.4. Schematic of the drying cell end platen

5.2.2 Specimens Description

Drying experiments were performed on glass beads specimens. Figure 5.6 shows the particle size distribution curve of the glass beads. Three experiments were conducted with three types of pore fluids: distilled water, distilled water with silt (12% by weight) and distilled water with Kaolinitic clay (12% by weight). The properties of the three specimens used in the investigation are summarized in Table 5.1.

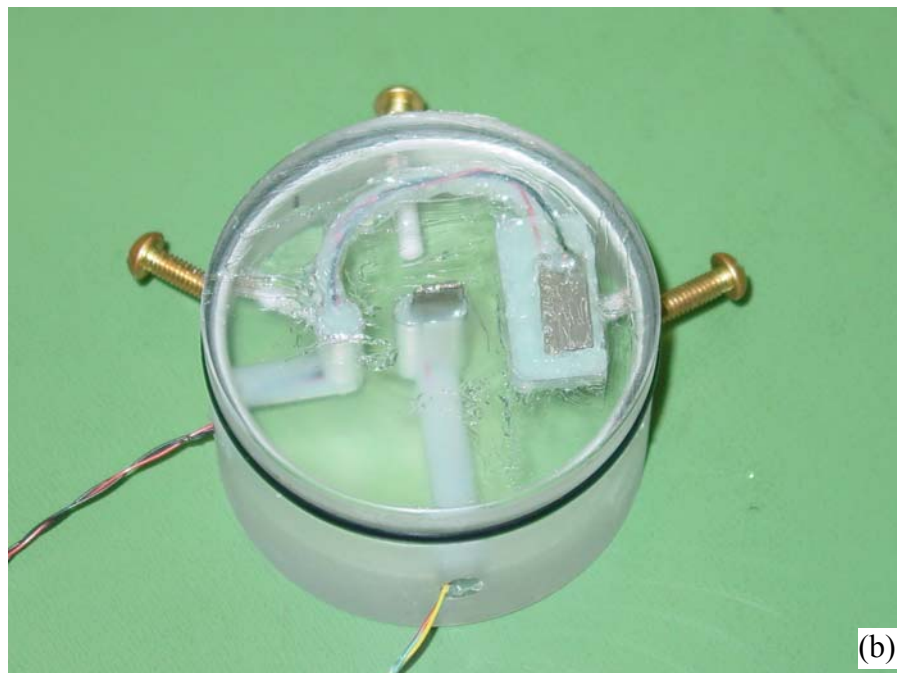
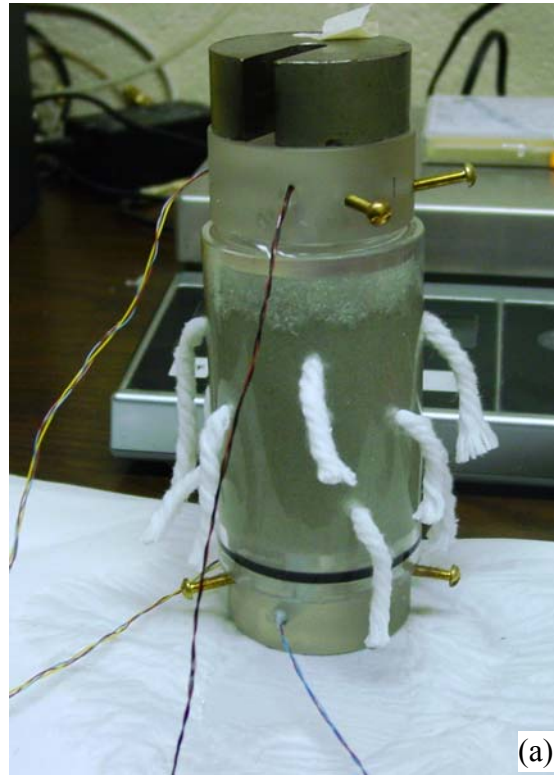


Figure 5.5. (a) Assembled drying cell with glass beads specimen (b) end platen with vertical and horizontal bender elements.

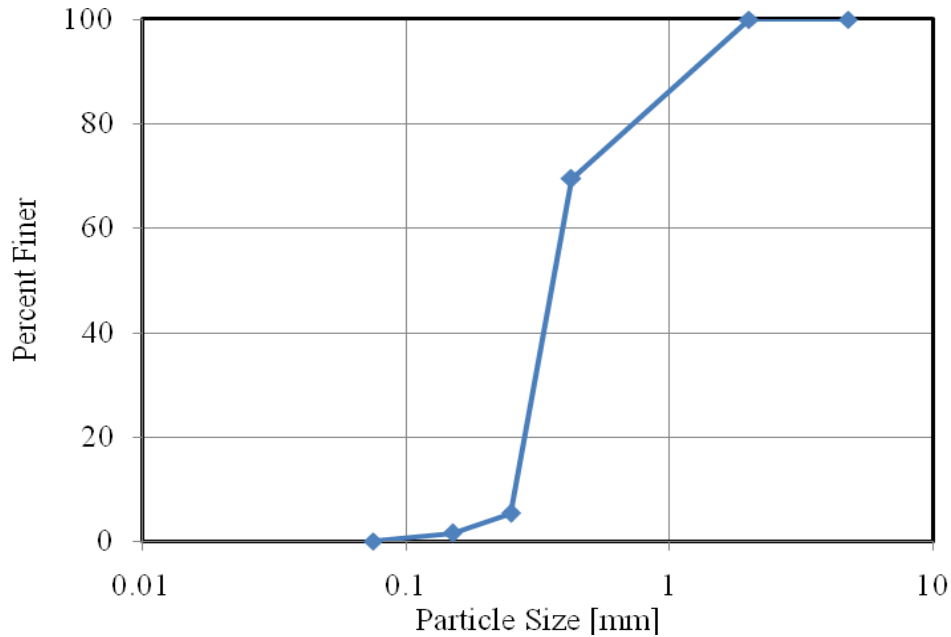


Figure 5.6. Particle size distribution of glass beads

Table 5.1. Summary of Specimens' Properties

Experiment Number	Pore Fluid	Porosity	Dry Density (kg/m ³)	D ₅₀ (mm)	C _u	C _c
1	Distilled Water	0.337	1757	0.36	1.54	0.87
2	Distilled Water + silt	0.340	1750			
3	Distilled Water + clay	0.336	1759			

In order to avoid salt precipitation or any other surface impurities, the glass beads were thoroughly cleaned using an acid washing procedure. The glass beads were first soaked in nitric acid for approximately five hours followed by thorough rinsing with deionized distilled water. The beads were then allowed to dry on a large open surface to avoid clumping of the particles.

Saturated specimens were prepared by pouring the glass beads in the cell full with water in three layers, and tapping each layer with a plastic rod to compact the specimen and eliminate air pockets. After taking the initial P- and S- wave velocity readings, the drying experiments were performed while the weight of the specimen was monitored in order to calculate the specimens' degree of saturation at any point throughout the test. Two light bulbs were placed on

the sides of the specimen to maintain a temperature of about 35° C to accelerate the evaporation process. The temperature was continuously monitored to maintain a constant temperature throughout the experiment to avoid changes in the surface tension of water and changes in confinement resulting from the expansion and shrinkage of the cell walls (Cho and Santamarina 2001). The experiments were carried out along the full range of saturation (from $S = 1$ to $S = 0$). P- and S-wave velocity readings were recorded at different stages until no further change in the weight of the specimen was recorded.

5.3 Experimental Results and Discussion

5.3.1 Wave Velocity Results

After each experiment, the arrival times of the S- and P wave were evaluated from the recorded wave traces. A MATHCAD algorithm was developed to evaluate the time shift between the different traces and calculate the wave velocities. Knowing the distance between the bender elements and the time between the sent and received signals, the wave velocities were calculated. Figure 5.7 depicts typical wave traces obtained throughout the drying process. The wave velocities for the three experiments were then compared to determine the effect of the presence of fine particles. Figure 5.8 depicts the obtained wave velocities for the three specimens. Moreover, in order to compare the percentage of increase and the different wave velocity patterns between the three experiments, the velocity values for each experiment were normalized with respect to the minimum (initial) velocity value of that experiment. The normalized velocity values are presented in Figure 5.9. Poisson's ratio (ν) was calculated from the recorded P- and S-wave velocities using Equation 5.1 to verify that the velocity values yielded Poisson ratio values within the typical range for soils. The calculated ν values for the three experiments are shown in Figure 5.10. It can be seen that the calculated Poisson's ratio values were between 0.15 and 0.2, where typical values for soils are between 0.1 and 0.3.

$$v = \frac{0.5 \left(\frac{V_p}{V_s} \right)^2 - 1}{\left(\frac{V_p}{V_s} \right)^2 - 1} \quad (5.1)$$

It can be seen from Figures 5.8 and 5.9 that the presence of the fine particles in the pore fluid has a profound effect on the values of the P- and S- wave velocities during drying. The P- and S- wave velocities in Experiment 1 (distilled water) exhibited very small change during the initial drying stages, and then their values started increasing at degrees of saturation under 20%. This is an expected result for this type of granular material; since it is a poorly graded (monosized particles) material ($C_c = 1.54$) it is expected that all pores will have similar sizes and the capillary forces take effect at low degree of saturation (Leong and Rahardjo, 1997, Yang et al., 2004). On the other hand, the P- and S- wave velocities in experiments 2 and 3 increase steadily from the early drying stages until the end of the experiments, this can be caused by two different processes: (1) the capillary forces acting within the silt and clay particles where it is known that silt clay have the ability to develop higher capillary forces at higher degrees of saturation (see for example Figure 2.7, Miller et al. 2002, Cho and Santamarina 2001). (2) during drying, fines migrate with the pore water to the contact area between glass beads, and with further drying they adhere together forming small bridges that increase the inter-particle contact stiffness and the overall stiffness of the medium. This process was further investigated using synchrotron x-ray microtomography and will be discussed in Section 5.4. It can also be seen from Figure 5.9 that the presence of clay in the pore fluid caused a larger increase in the wave velocity values than silt. This could be caused by the higher capillary forces that can develop in clays, and the stronger aggregation between clay particles.

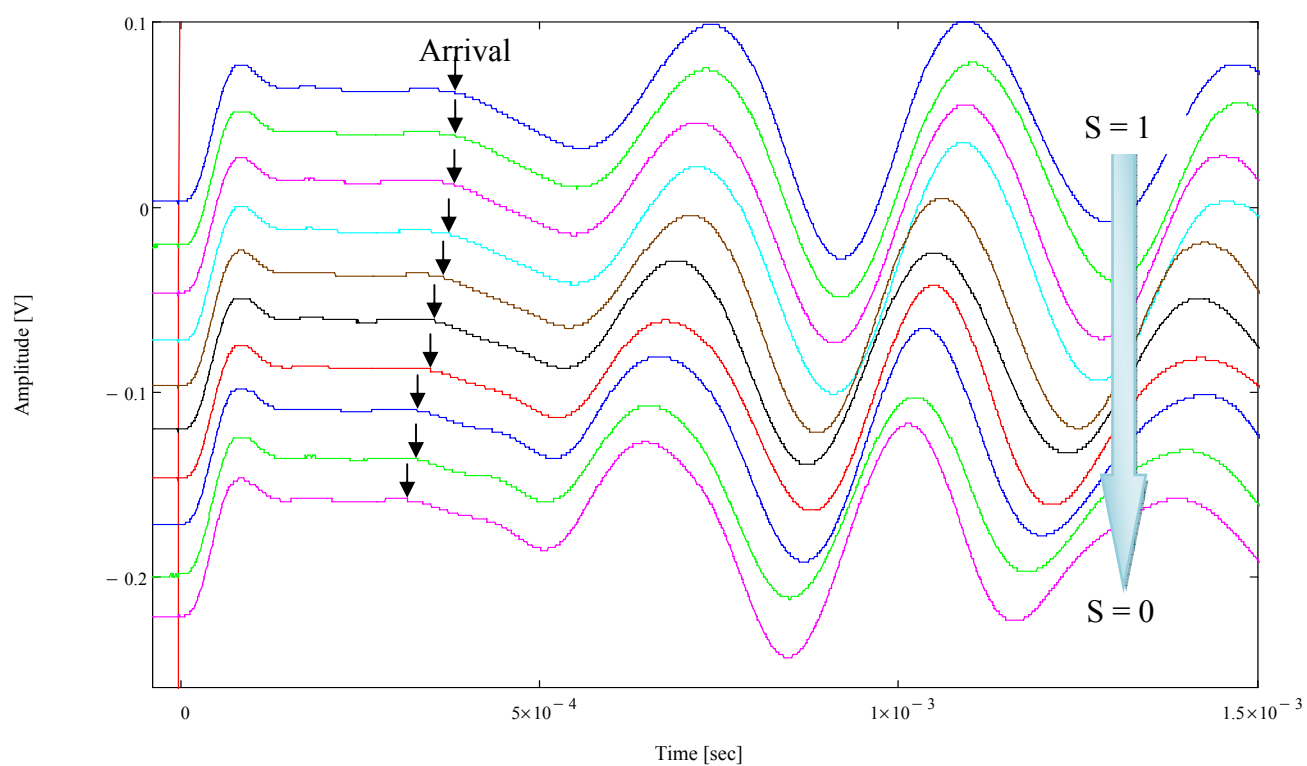
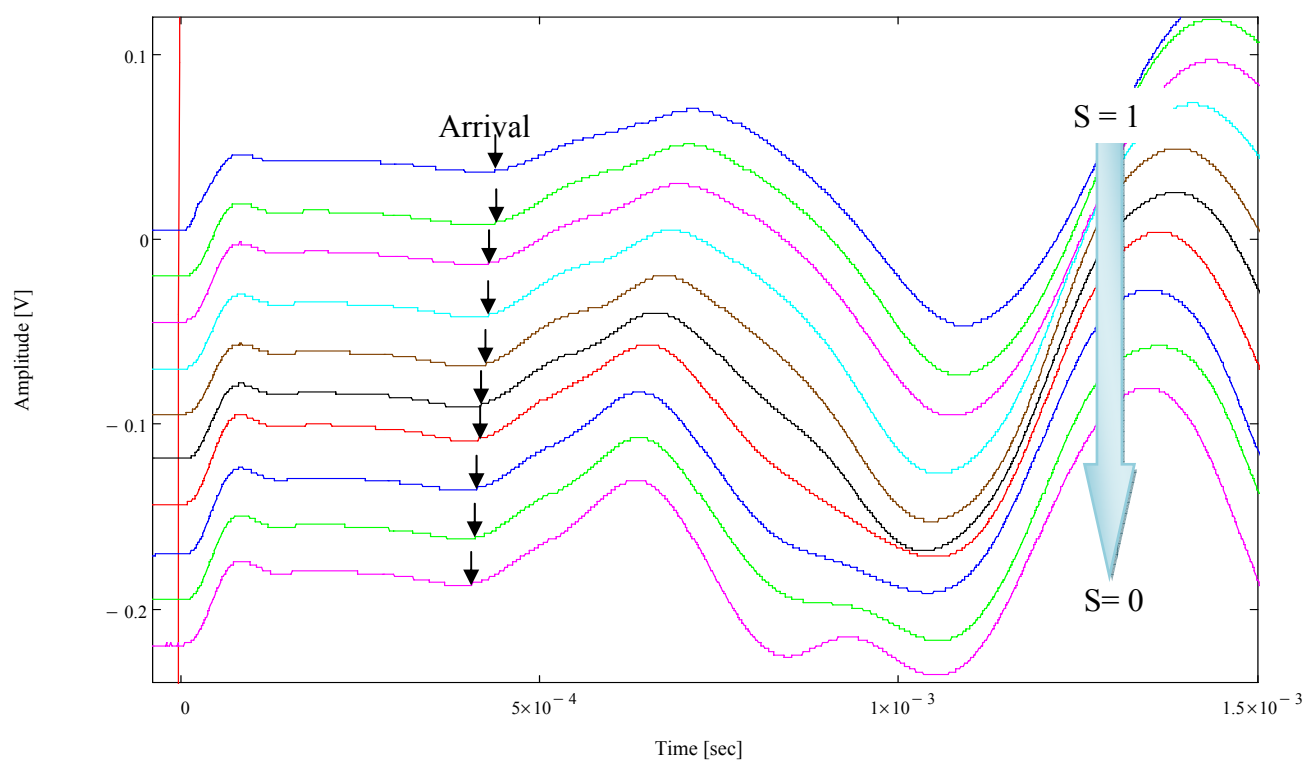


Figure 5.7. Typical wave traces obtained throughout the drying process

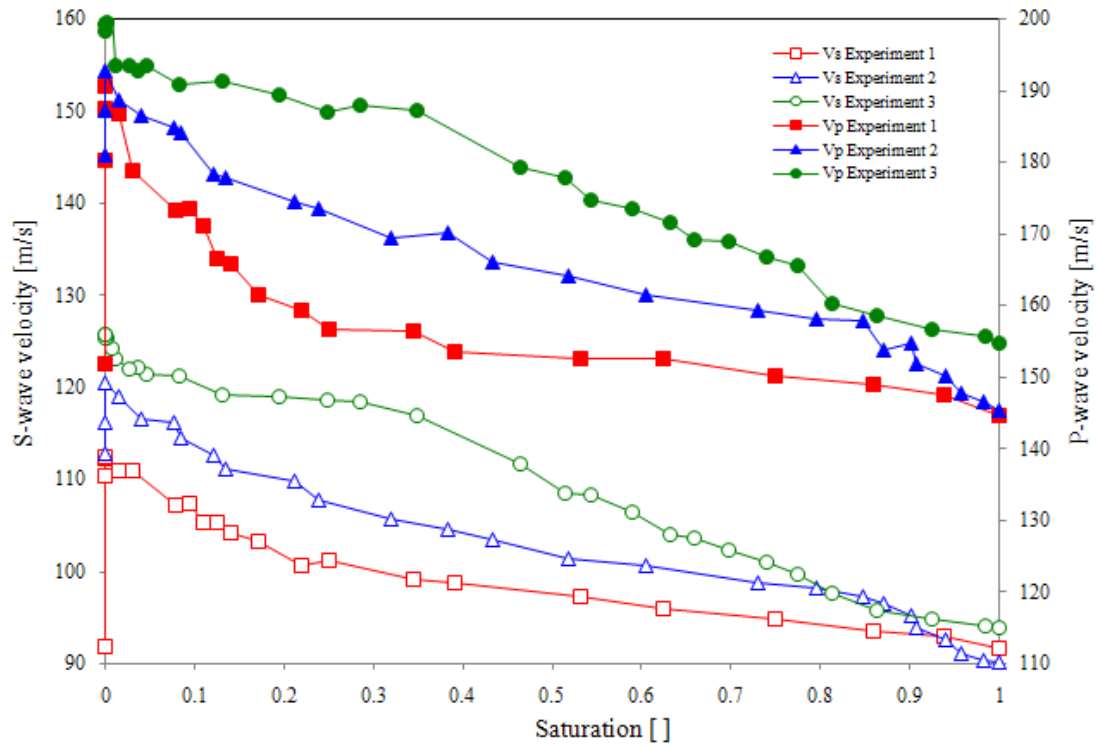


Figure 5.8. S- and P- wave velocities for the three drying experiments

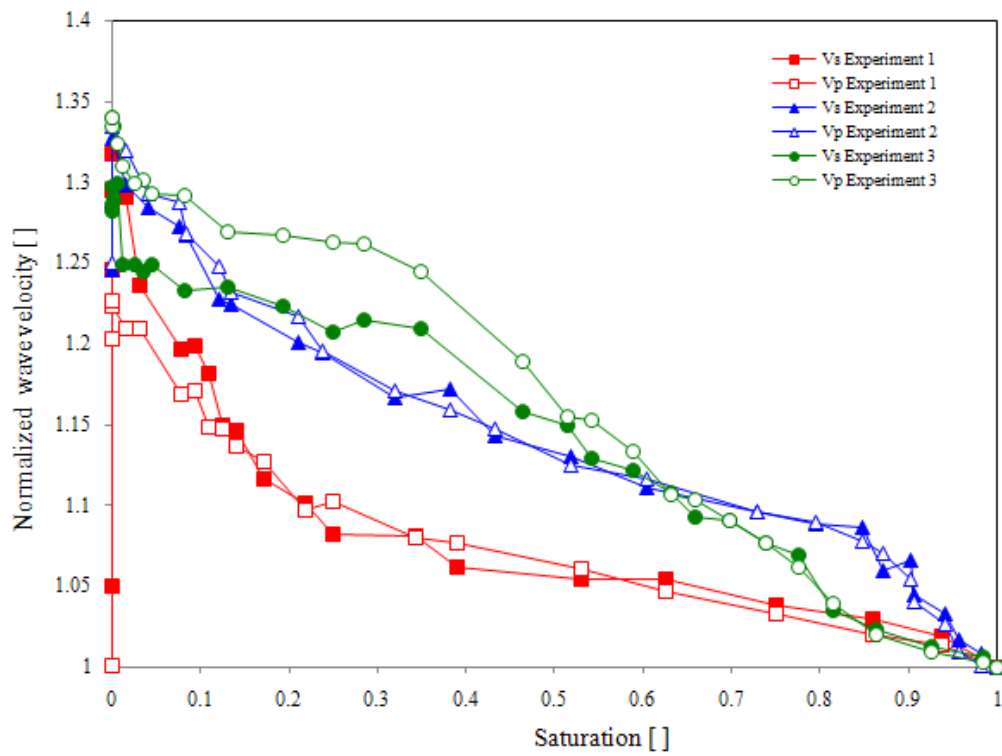


Figure 5.9. Normalized S- and P- wave velocities for the three drying experiments

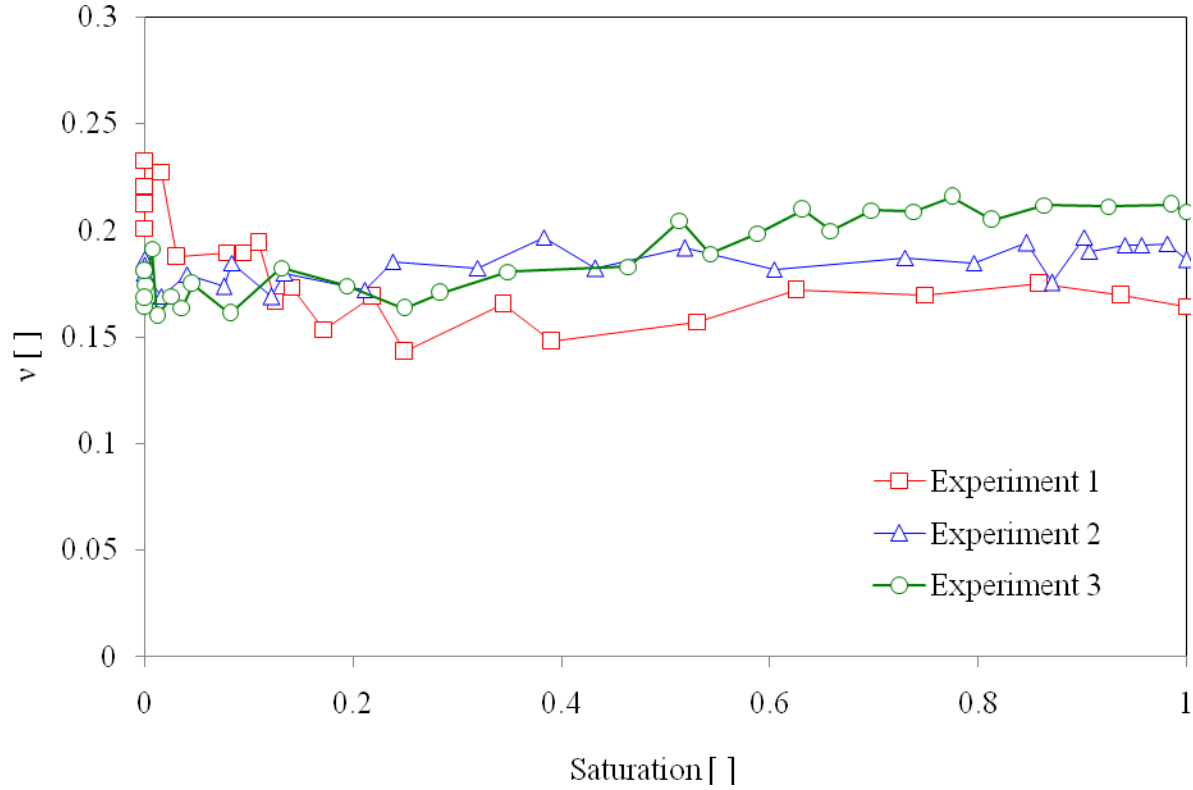


Figure 5.10. Calculated Poisson Ratio values for the three drying experiments

5.3.2 Effect on Soil Stiffness

To evaluate the effect of the presence of the fine particles in the pore fluid on the small strain shear stiffness (G) and the constraint modulus (M) of the soil, G and M were evaluated from the wave velocity data and compared for the whole drying range. Equation 5.2 was used to calculate G from the shear wave velocity and the density. Similarly, M was calculated from the measured P-wave velocity (Equation 5.3). Moreover, in order to determine the relative variations in the stiffness values and to highlight the effect of the presence of the fine particles in the pore fluid, the calculated G and M values were normalized to their initial (minimum) values. Figures 5.11 and 5.12 illustrate the calculated and normalized values of G and M , respectively, throughout the drying process. Furthermore, knowing the soils shear and constraint moduli, the soils' bulk modulus (B) was calculated using Equation 5.4. The normalized B values are depicted in Figure 5.13.

$$G = \rho \cdot V_s^2 \quad (5.2)$$

$$M = \rho \cdot V_p^2 \quad (5.3)$$

$$B = M - \frac{4}{3} \cdot G \quad (5.4)$$

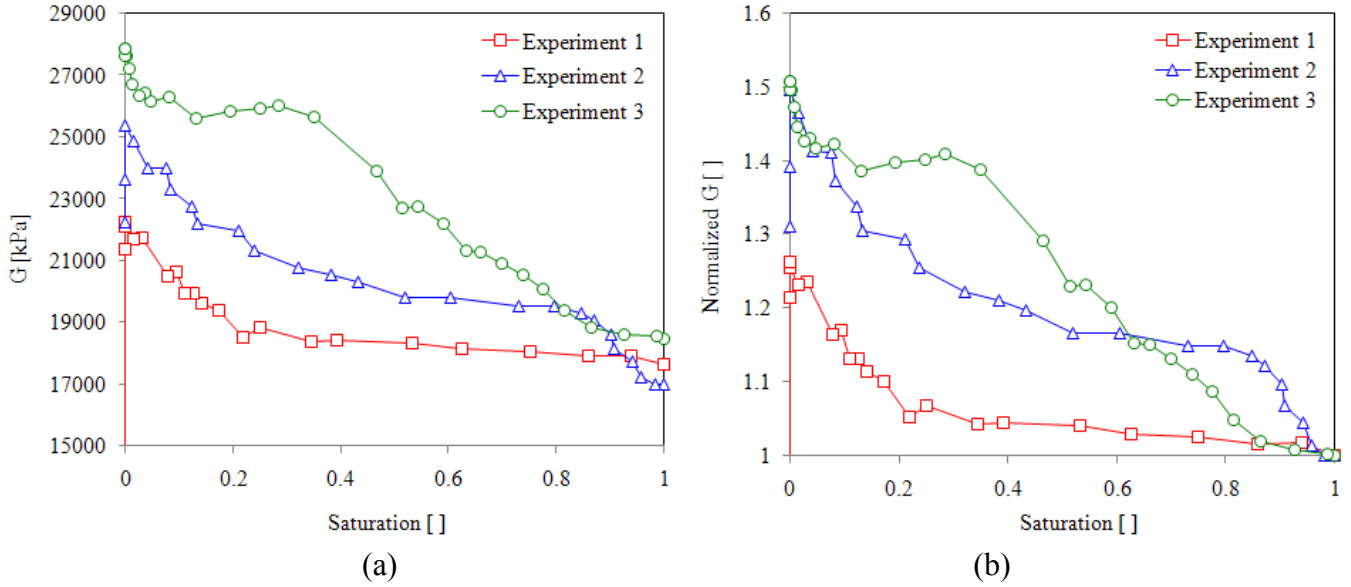


Figure 5.11. Small strain shear stiffness (G) values during the drying process. (a) Calculated values, (b) Normalized values.

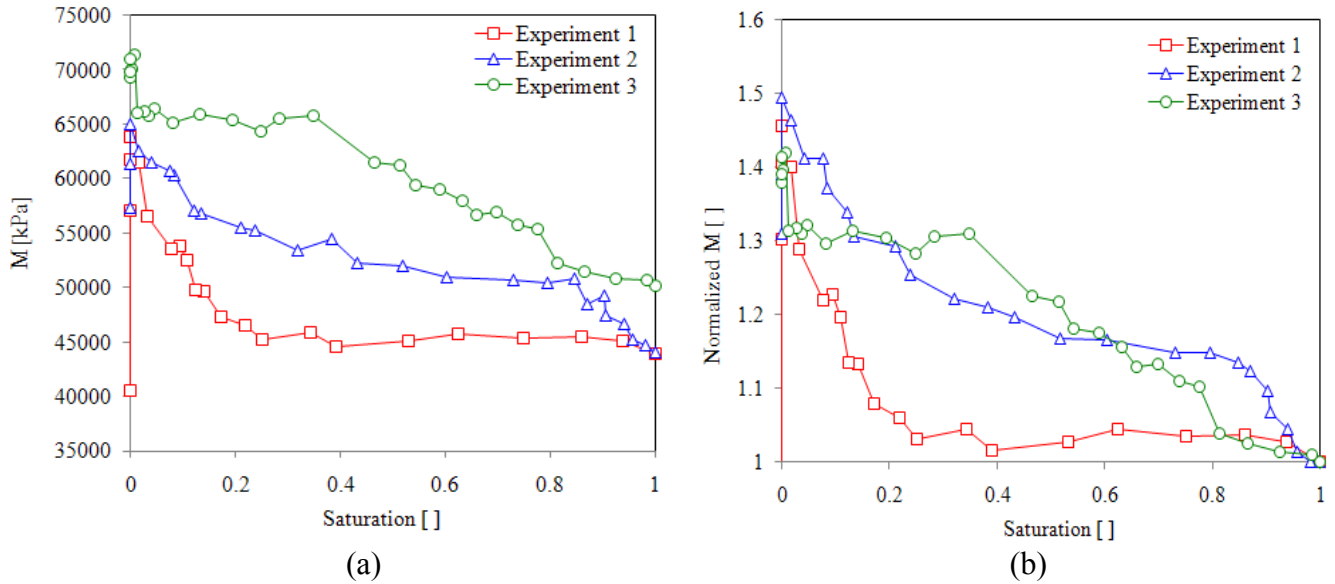


Figure 5.12. Soils' constraint modulus (M) values during the drying process. (a) Calculated values, (b) Normalized values.

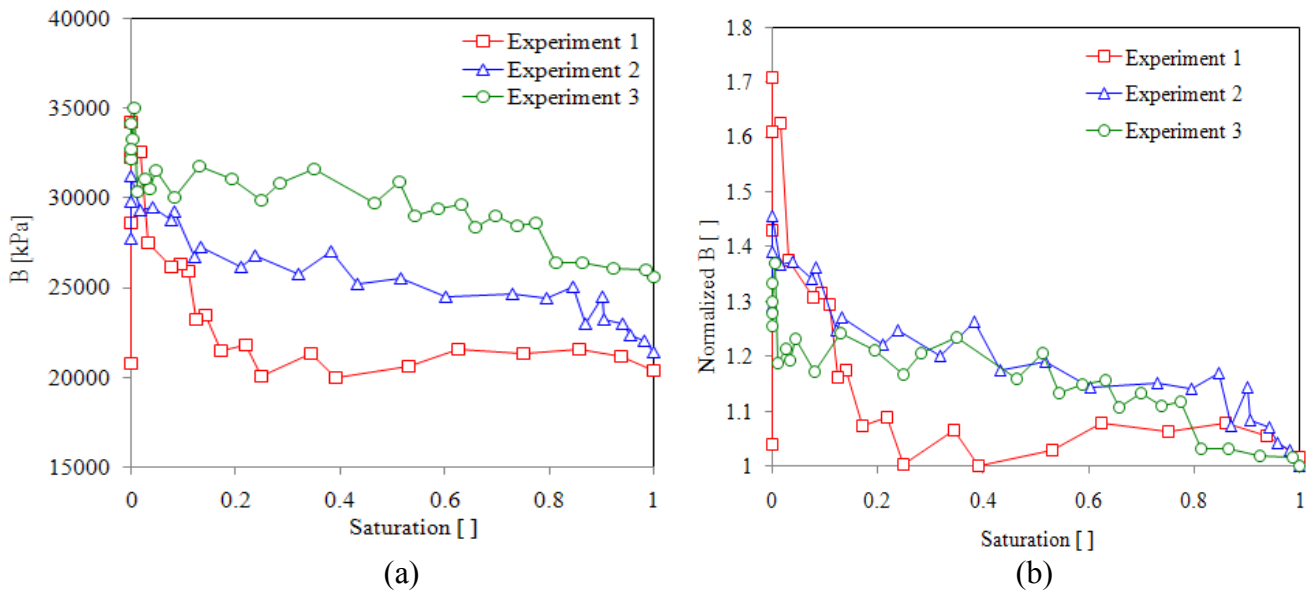


Figure 5.13. Soils' bulk modulus (B) values during the drying process. (a) Calculated values, (b) Normalized values.

It can be seen from Figures 5.11 through 5.13 that the presence of the fines in the pore fluid has a profound influence on the values of G, M and B. During drying, the value of G increased by about 25% when using distilled water. However, adding silt or clay to the pore fluid caused an increase of about 50% at the end of the experiments. Moreover, the three experiments showed different trends during the drying process. This could be caused by the difference in the particle sizes and time required for the fine particles to settle on the surface of the glass beads. It can be seen for example the delayed effect of the clay in comparison to the silt in Figure 5.12b, this could be due to the clays' ability to stay in suspension for a longer time than silt. Similarly, higher M values were obtained for the clay specimen, followed by silt and the pure water sample; however, the three specimens exhibited an increase in the value of M of about 40% to 50% at the end of the drying process. A similar trend was also noticed for the bulk modulus B where the clay specimen had the highest B values followed by the silt specimen.

It can be seen that the presence of fine particles not only increases the values of the values of the shear, constraint and bulk moduli of soils, but also highly affects the shape of the

drying curve. The properties of the pore fluid could have a significant effect on the stiffness of unsaturated soils at different degrees of saturation. Therefore, it is important to study the changes in the properties of the pore fluid throughout the drying process. The following section presents a study of the pore fluid properties using synchrotron X-ray computed tomography and image analysis techniques.

5.4 Synchrotron X-ray Tomography

5.4.1 Background

In recent years, X-ray microtomography (μ CT) has emerged as a very useful tool to non-destructively visualize and study the internal structure of geomaterials. It is an imaging technique in which an object is placed between an x-ray source and a detector, and the object is rotated while the x-ray passes through it collecting information about its internal structure. The information is reconstructed to create a two-dimensional cross section “slice” of the object that can be used to view the internal structure of the object (Cromwell, 1984). Slices can also be stacked together to create a three dimensional volume view of the scanned object.

μ CT was first used in diagnostic medical fields to acquire 3D images of the internal body organs and tissues. First it was used to scan the human head, and then it was used all over the body. In the field of materials research, it was widely used in petroleum engineering for core analysis, μ CT images were used for characterization of reservoir rocks, for mineral formation and rock property evaluation (e.g. Honarpour et al. 1985). It was also used for monitoring two and three-phase experiments in porous media (e.g. Bansal et al. 1991) and to verify and analyze the multiphase fluid flow in porous media (e.g. Cromwell et al. 1984). In the field of geotechnical engineering, μ CT was used to observe localized deformations and shear bands in specimens undergoing triaxial compression, void ratio evolution inside shear bands, void ratio variation in soil specimens, and characterization of failure in soils. (Desrues et al. 1996, Otani et

al. 2000, Alshibli et al. 2000a,b). On the other hand, Matsuhima et al. (2006) visualized and quantified the motion of each grain inside a triaxially compressed specimen using μ CT images.

A significant enhancement to the traditional x-ray tomography systems was the use of synchrotron radiation. Synchrotron radiation sources generate higher intensity beams (10^6 times the conventional x-ray beam). Moreover, such sources are able to generate monochromatic x-ray beams reducing or eliminating the beam hardening effect that could otherwise be misinterpreted as density variations (Alramahi and Alshibli, 2006). The synchrotron sources are also capable of generating highly coherent beams spanning the full range of x-ray energies from infrared to hard X-rays up to 400 keV (Boller et al. 2006). Images obtained by synchrotron radiation sources are characterized by very high spatial resolution, low signal to noise ratio, and enhanced quantitative exploitation (Boller et al. 2006). These features make synchrotron radiation a very useful tool to study geomaterials. Image analysis operations such as phase segmentation, tracking density variations and analyzing spatial patterns can be performed more accurately using high resolution synchrotron images.

The enhanced capabilities of the synchrotron micro-computed tomography (SMCT) were utilized by many researchers to characterize the internal structure and monitor different processes in geomaterials. Culligan et al. (2004) used SMCT to perform interfacial area measurements between wetting and non wetting fluids during primary and secondary drainage and imbibition cycles. On the other hand, Wildenschild et al. (2005) used SMCT images to quantitatively study the multi-phase flow through sands and to examine the drying rate dependency of the residual water saturation. SMCT images were also utilized by Ham and Willson (2005) to characterize the pore geometry and topology and entrapped residual air phase. They also examined the correlation between the pore structure and the non wetting blobs.

5.4.2 The Interaction of X-ray with Matter and CT Number

When the x-ray passes through matter, the photons are either scattered with some loss of energy, or completely absorbed in a photoelectric interaction, as a result, some photons will be lost from the beam while passing through matter. Only part of the photons will remain in the beam and reach the detector, which means that a beam of photons is not degraded in energy as a result of passing through matter, just attenuated in intensity. The probability of photon attenuation can be expressed per thickness of the attenuator, with the Beer equation:

$$I(x) = I_0 e^{-\mu x} \quad (5.5)$$

where I is the x-ray intensity at certain distance x , I_0 is the x-ray intensity when it is emitted from the source, and μ is the linear attenuation coefficient which can be defined as the probability that an x-ray or gamma-ray photon will interact with the material it is traversing per unit path length traveled. It is usually reported in cm^{-1} . The linear attenuation coefficient depends on the photon energy, the chemical composition and physical density of the material. The output of a CT scan represents the attenuation values of the different parts in the object expressed as CT numbers. Usually, higher CT numbers represent higher attenuation values, and higher density. Therefore, by tracking the changes in CT numbers in different scans, density variations could be traced throughout the scanning process.

5.4.3 Synchrotron X-ray Facility

Scanning for this study was performed at the GeoSoilEnviroCARS 13-BM-D synchrotron x-ray beam at Argonne National Laboratory Advanced Photon Source (ANL-APS) facility (Figure 5.14). This beam line uses an Si (220) channel-cut monochromator covering an energy range of 5 to 45 keV. The transmitted x-rays are imaged with a single crystal YAG scintillator, a microscope objective and a 1242 X 1152 pixel fast CCD detector (Rivers et al. 1999). This system is capable of collecting the data for the entire vertical field of view (~5mm) all at once

without the need to vertically translate the specimen on the x-ray stage. This significantly reduced the time required to collect the data at different heights of the scanned object. Scanning was performed at 35 keV energy, with a spatial resolution of 20.35 μm . The resulting vertical field of view was about 6 mm. Scanning time was about 15 minutes per scan.

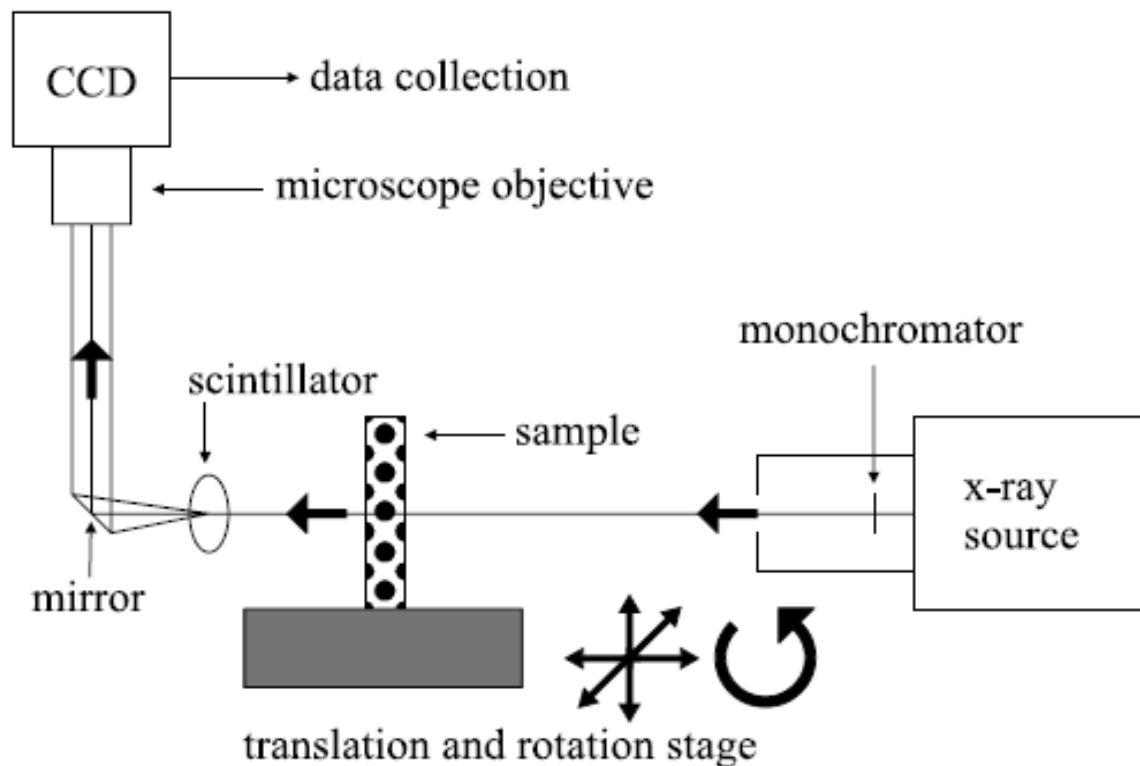


Figure 5.14. Schematic of the synchrotron x-ray system 13-BM-D at ANL (Culligan et al. 2004)

5.4.4 Description of Tomographic Specimens

To compare the effects of the presence of fine clay and silt particles on the properties of the pore fluid, CT scans were performed on glass beads specimens containing the two types of pore fluids: water with silt and water with clay. The glass beads and pore fluids were similar to the ones used in the drying cell experiments. Saturated specimens were prepared in small glass vials with an internal diameter of 10 mm. Glass beads were poured into the vials while continuously tapping the walls to eliminate any air pockets. Figure 5.15 depicts a typical glass

vial used to prepare the specimens. Three identical specimens were prepared for each type of pore fluid, and then they were dried to different degrees of saturation. Scanning was performed for each type of pore fluids at three different stages of the drying process. Figure 5.16 presents example CT slices representing horizontal and vertical cross sections of a scanned specimen. It can be seen that the liquid phase in this slide is in the funicular regime where air pockets are present, however, most of the liquid phase forms large interconnected bodies. It can also be seen that in some dry areas, the liquid exists in small separated bodies at the interparticle contacts that are denoted as pendular rings.

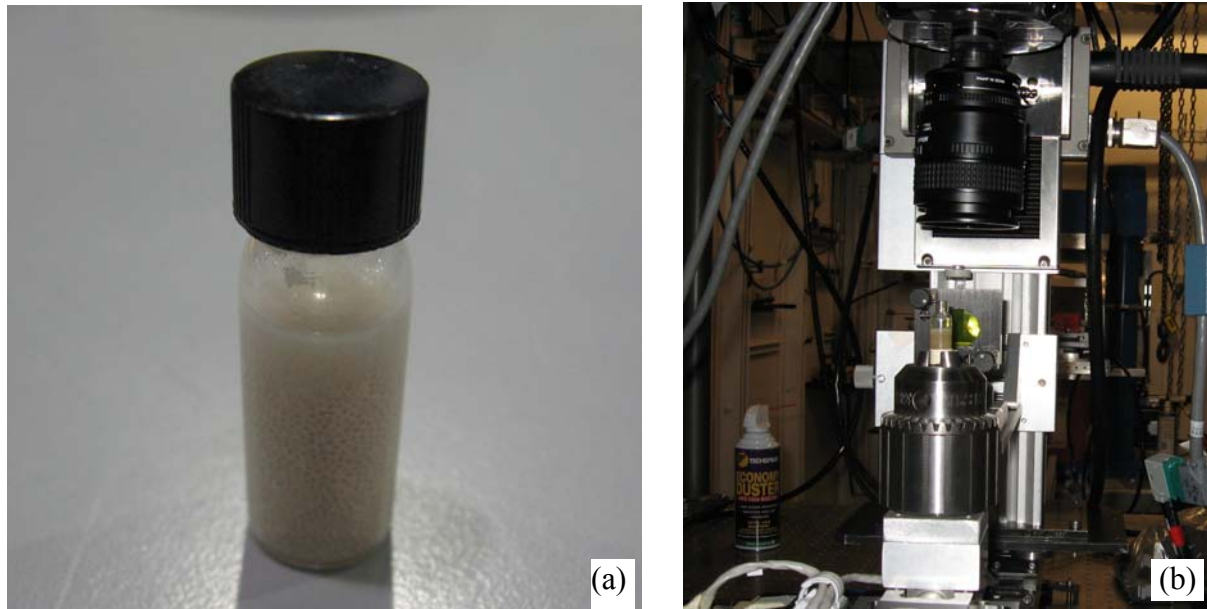


Figure 5.15. (a) Typical glass vial containing glass beads. (b) glass vial mounted on the CT stage

5.4.5 Image Analysis

In order to study the properties of the pore fluid in each specimen, image segmentation was performed to separate the three phases (solids, air, and water) in each image. IDL algorithms were developed to separate the different components. The algorithms use a combination of threshold value and phase continuity criteria to assign every pixel to one of the three phases. It

starts with comparing the pixel CT values to predetermined threshold CT values for each of the three phases. However, because there are no clearly defined threshold values for every phase in the medium, a different algorithm examines the connectivity of each group of pixels in order to generate spatially connected groups of pixels assigned to one of three possible phases. This technique along with image filtering yielded binary CT images where each pixel was labeled with the phase it represents. Figure 5.17 illustrates an example CT slice where the water pixels were extracted from the rest of the image to be studied separately.

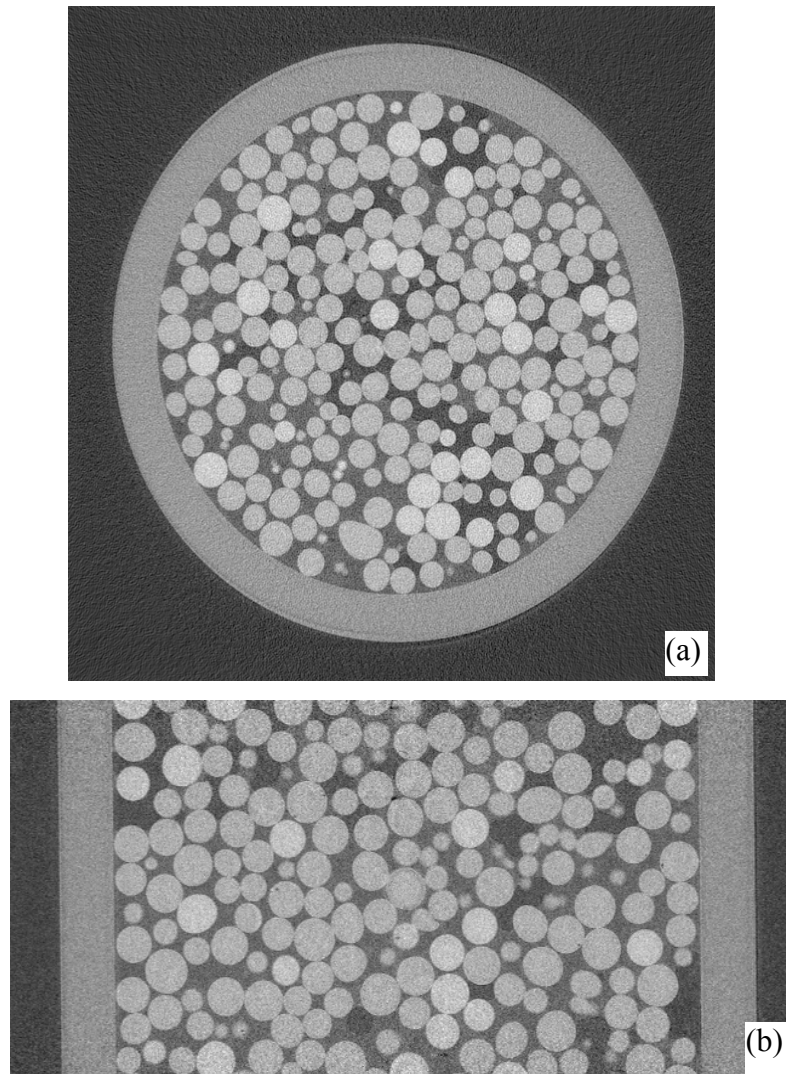


Figure 5.16. Typical CT images (a) horizontal cross section, (b) vertical cross section

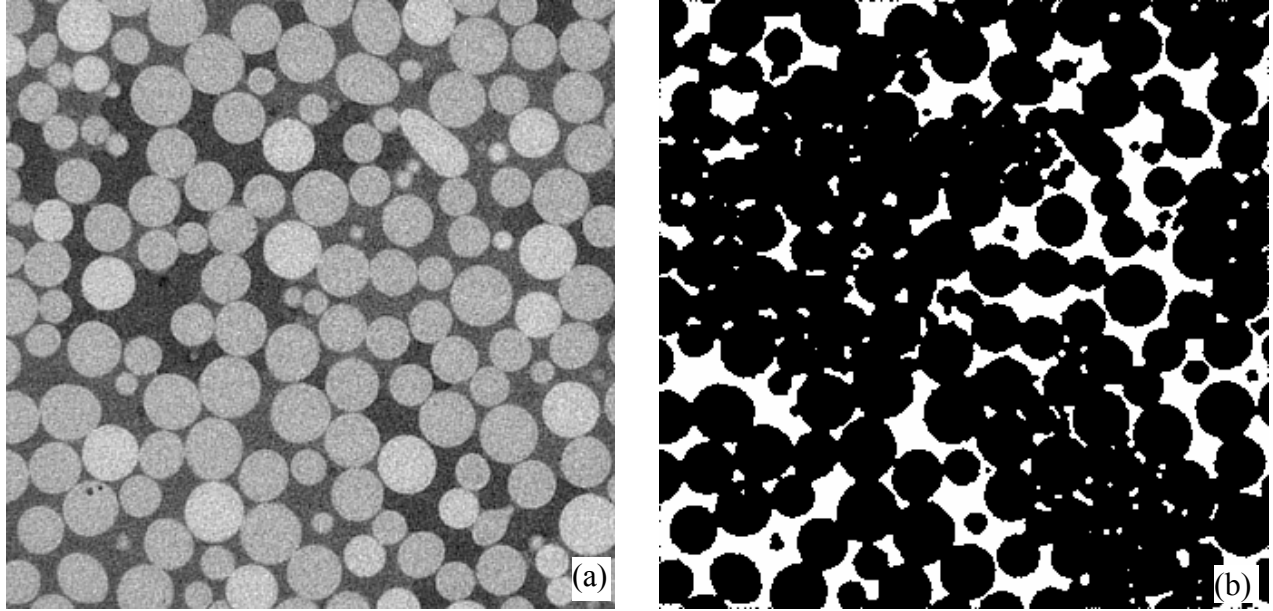


Figure 5.17. (a) Original grayscale image (b) binary image representing the liquid phase

Before performing any measurements on the obtained volumes, it is important to determine whether their sizes were large enough to be representative elementary volumes (REV) that represents the properties of the whole specimen. Therefore, the minimum size of an REV should be determined to serve as a guideline for the size of the volumes to be used in the analysis. An IDL algorithm was developed to perform this task, where it considers cubic volumes of increasing sizes and calculated their porosity. The minimum REV size is determined as the size of the cube where the calculated porosity becomes constant. This procedure was performed for all the scanned volumes and the results are depicted in Figure 5.18. It can be seen that having very small cube sizes yielded erroneous porosity values since the volume could lie entirely inside a void or a solid space. The values of the porosity for all the volumes stabilized for cube sizes of about 125 pixels (yielding a volume of about 2 million pixels). Therefore, this was taken to be the minimum volume size for any further analysis.

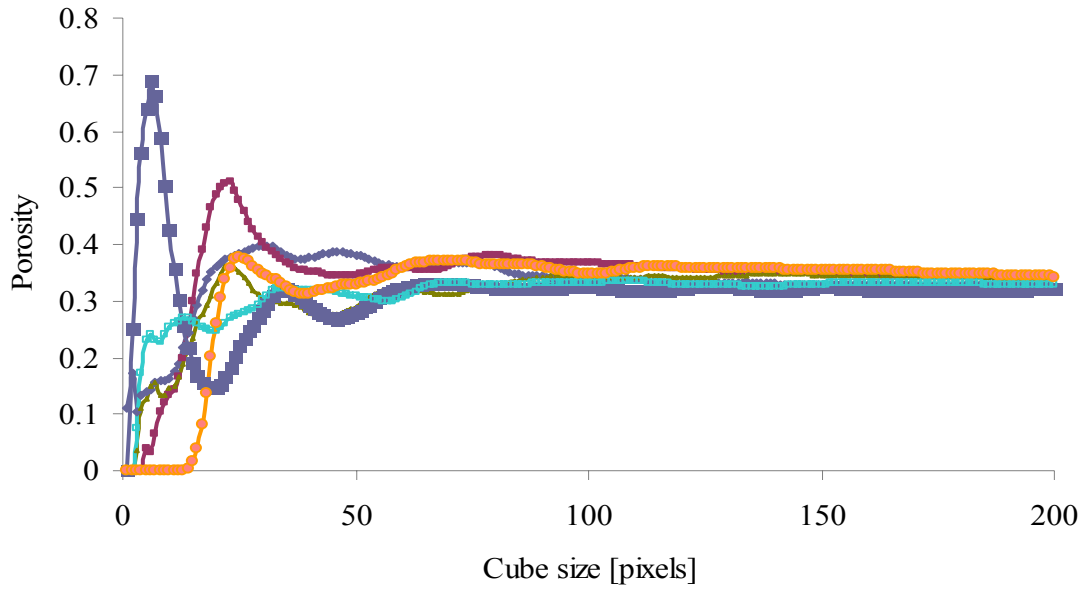
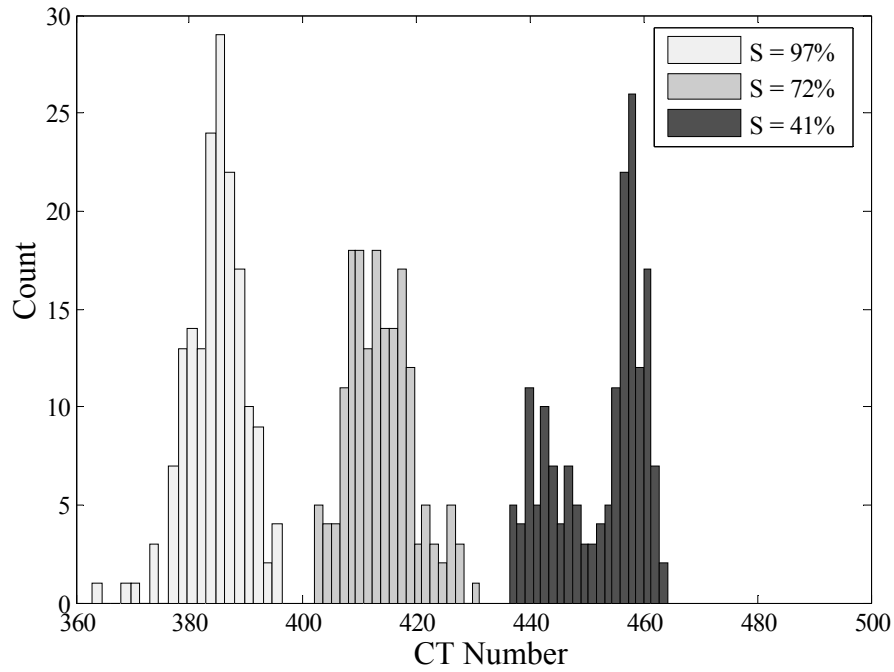
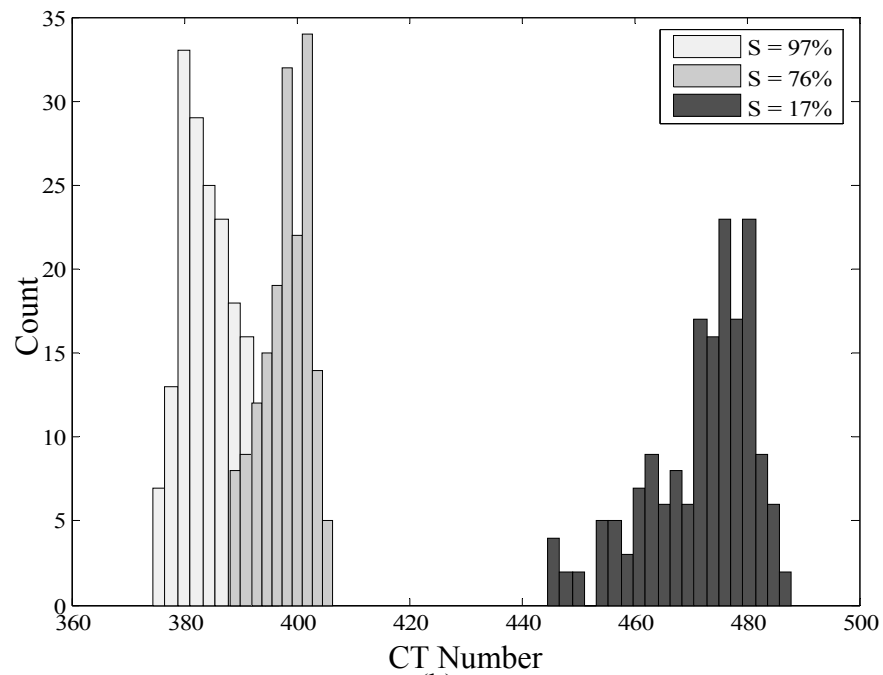


Figure 5.18. Determination of the minimum REV size

After the pore fluid pixels were separated, their CT values were analyzed to track the changes in the pore fluid density caused by the increase in the concentration of the fine particles during the drying process. An Interactive Data Language “IDL” algorithm was developed to separate the water pixels in each slice of the volume, and calculate the average CT value of the water pixels. The algorithm also counts the water, air and solid pixels and calculates the degree of saturation of each slice. After this process was performed, histograms of the CT values of all the slices of each volume were prepared to determine the changes in the pore fluid density at different degrees of saturation. Figure 5.19 depicts the CT value histograms for the two types of pore fluids. Moreover, the average CT value for each of the volumes was calculated and plotted against the degree of saturation (Figure 5.20). A linear equation was used to fit the data for the silt and clay experiments. It was noticed that there was a strong linear correlation between the average CT number of the pore fluid and the degree of the saturation.



(a)



(b)

Figure 5.19. Histograms of CT values at different stages of drying (a) water and silt, (b) water and clay

To estimate the density values and the change in the concentration of fine particles during drying, a density curve was developed where water and soil mixtures were prepared at different

concentrations. For each mixture, the density was determined using weight and volume measurements. Figure 5.21 illustrates the obtained fluid density curve. The initial density of the pore fluid used in the experiments (12% by weight) was about 1.06 g/cm³.

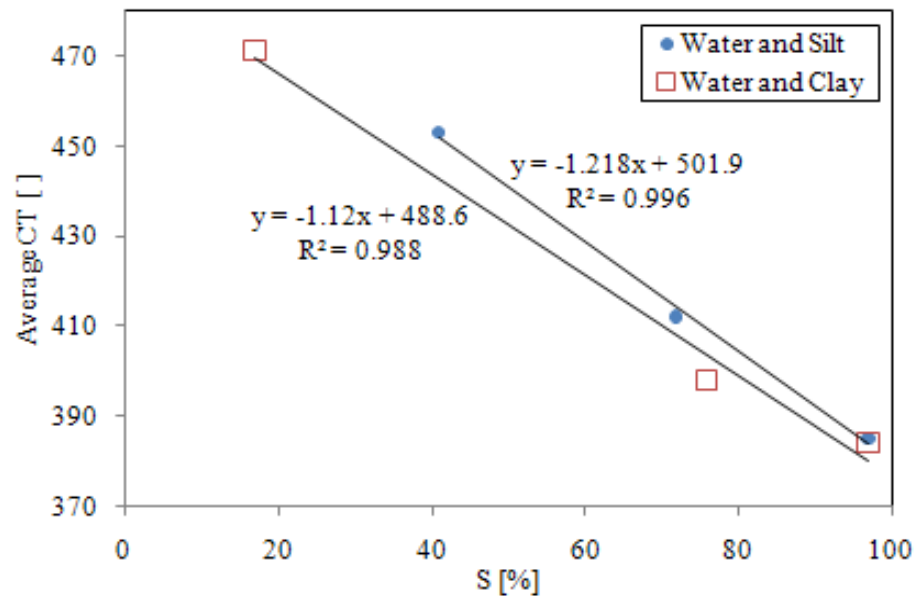


Figure 5.20. Average overall CT numbers during drying

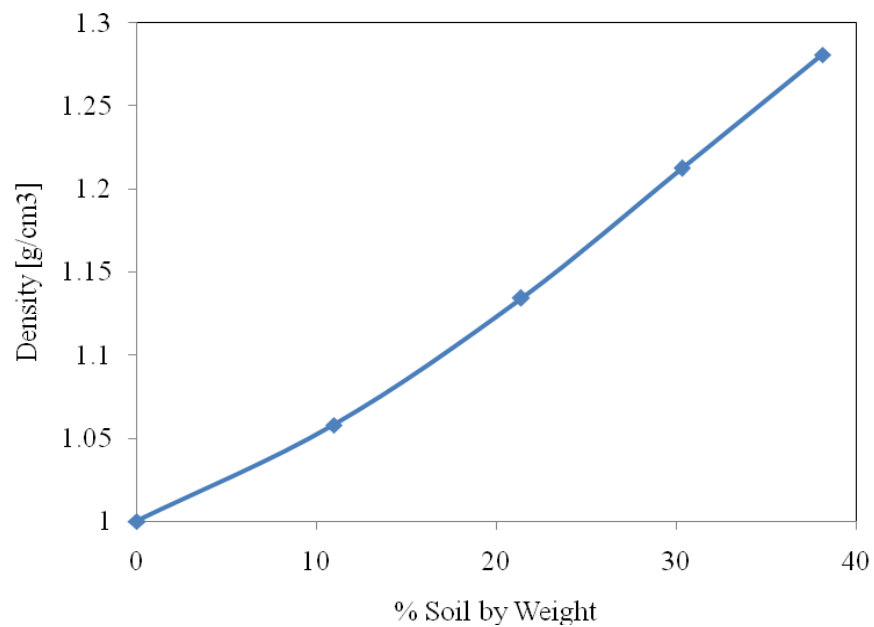


Figure 5.21. Fluid density at different soil concentrations

The changes in the pore fluid density were calculated from the changes in CT numbers assuming a linear relationship between the density and CT number (Otani, 2006; Alramahi et al., 2005). Therefore, by knowing the initial pore fluid density and the percentage of change in the CT numbers, the pore fluid density could be calculated during the drying process. Figure 5.22 illustrates the pore fluid density values for the two experiments at different stages of the drying process.

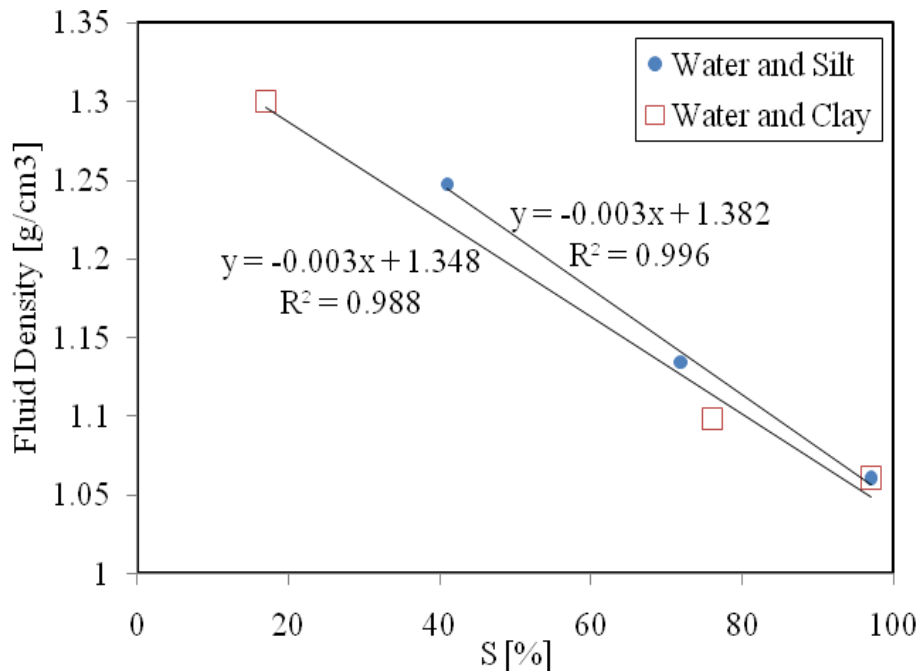


Figure 5.22. Change in fluid density during drying process

It can be seen from Figure 5.22 that the density of the pore fluid increased during the drying process due to the increase of the fine particles' concentration, indicating the migration of the fine particles to the contact area as drying proceeds. At $S = 41\%$ the density of the pore fluid for the water and silt experiment increased to a 1.25 g/cm^3 corresponding to a silt concentration of 34%. On the other hand, in the clay and water experiment, the pore fluid density increased to 1.3 g/cm^3 corresponding to a clay concentration of about 40% at $S = 17\%$. Such increase in the

concentration of the fines was caused by the migration of the fines to the regions where water is located during the evaporation process. Therefore, at low degrees of saturation, high concentrations of the fines are present at the inter-particle contact region forming small bridges between the particles and causing the increase in the small strain stiffness in the medium. Figure 5.23 illustrates an example CT image at a low degree of saturation ($S = 17\%$). It can be seen that the pore fluid, containing high concentrations of fine particle, is located mainly at the inter-particle contact areas.

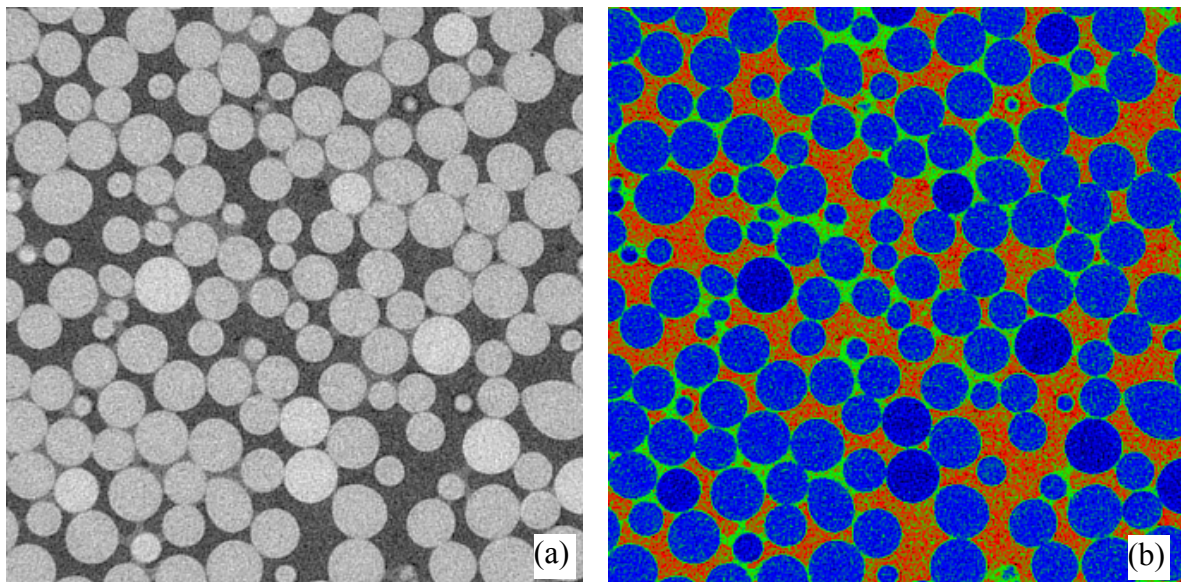


Figure 5.24. Example CT image at low degree of saturation (a) gray scale, (b) enhanced contrast; blue: solid, green: water, red: air

5.4.6 Effect of Fluid Location on the Concentration of Fines

It was demonstrated in the previous Section that as the drying process proceeds, the overall density of the pore fluid increases as a result of the increase of the concentration of fine particles. However, since the fluid density values were averaged across the whole cross sectional area, local variations in the fluid density values were not fully identified. Therefore, to better understand the behavior of the fines in the pore fluid during drying, it is important to take a

closer look at the local density variations within the cross sections. It is also of great importance to determine whether the pore fluid located in close proximity to the interparticle area has different characteristics than the fluid away from it since the presence of the fine particles at the interparticle contact area has a profound effect on the overall soil stiffness.

To study the effect of the location and size of the pore fluid bodies on their density and the concentration of fines, several cross sections along the height of the scanned specimens were analyzed. In each cross section, pore fluid bodies were identified and classified based on their location and proximity to the interparticle contact area to one of two possible classes: pore body fluid or interparticle fluid. The pore body fluids are the large bodies of fluid that fill the large pore volumes away from the interparticle contact area. On the other hand, the interparticle fluids are the fluid bodies that are located in close proximity to the interparticle contact areas or pore throats. This classification aims at determining whether there is a significant difference in the fines concentration in the areas closer to the interparticle contacts. IDL-iImage utility was used to identify the different pore fluid bodies and classify them according to their proximity to the interparticle contacts. Regions of interest (ROI) were generated for each fluid body and their CT values were extracted to determine their mean value. Figure 5.25 depicts an example CT slice where different pore fluid bodies were identified and classified. Red regions denote pore body fluids while yellow regions denote interparticle contact fluids.

Pore fluid bodies were identified in different cross sections along the height of the scans in order to obtain statistically representative samples. The average CT value of the different fluid bodies from the water and silt as well as the water and clay samples were calculated and compared. The results are depicted in Figure 5.26. It can be seen that in both types of specimens the pore fluid located at the interparticle contact area had significantly higher CT values than the pore fluid located away from the interparticle contact area, indicating higher fluid density values

resulting from the increase in the concentration of fines around the interparticle contacts. This is the result of the precipitation of the fine particles at the surface of the glass beads during drying, creating buttresses of silt or clay that increase the interparticle stiffness and the P- and S- wave velocities as was demonstrated in Section 5.3.

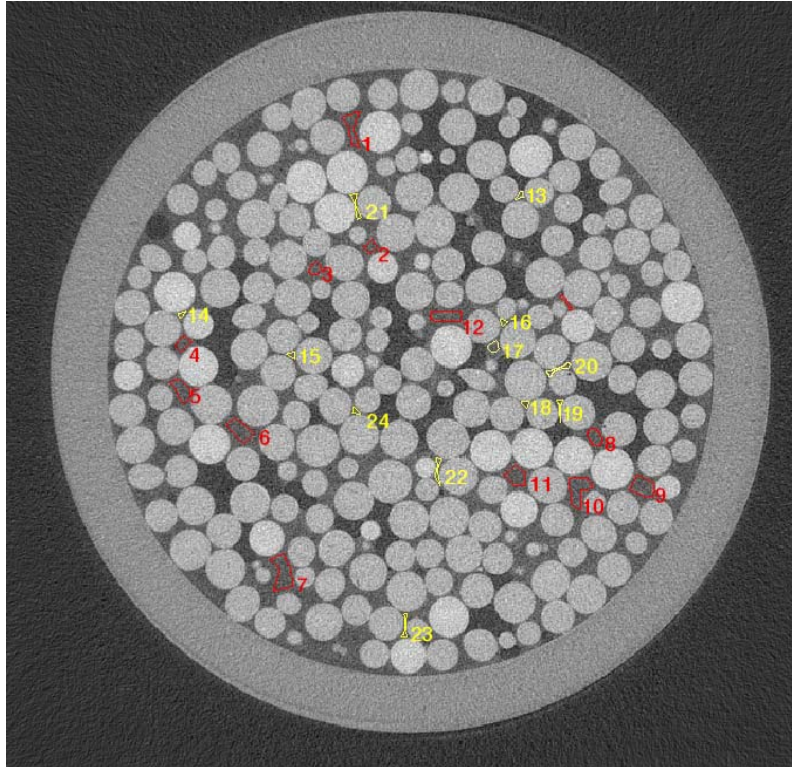


Figure 5.25 Example CT slice with different pore fluid bodies. Red indicates pore body fluid and yellow indicates interparticle fluid

In order to compare the changes in the pore fluid properties between the silt and clay samples, histograms of the CT numbers of the pore fluid bodies located at the pore body and the interparticle contact area were compared. The four histograms are depicted in Figure 5.27. It can be seen that the pore fluids in the clay specimen had slightly higher values than the silt specimen. This is caused by the ability of fine clay particles to stay in suspension for a longer time, causing the pore fluids to maintain higher density values throughout the drying process.

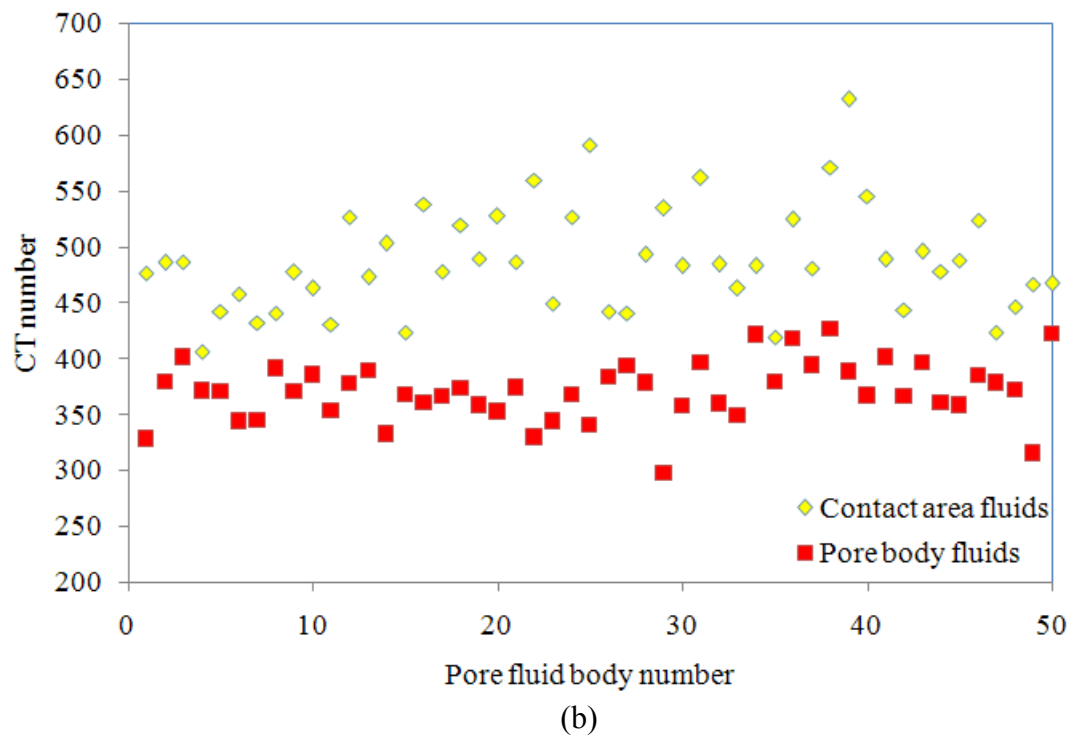
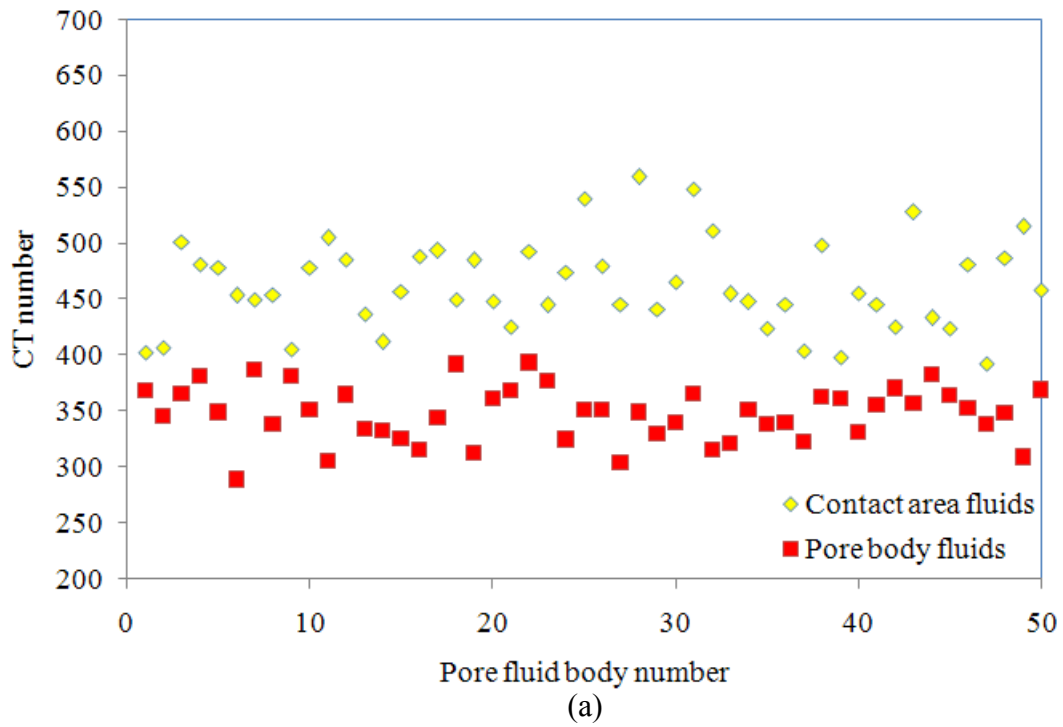


Figure 5.26. Average CT number for pore fluid bodies located in the pore body and interparticle contacts. (a) silt specimen (b) clay specimen.

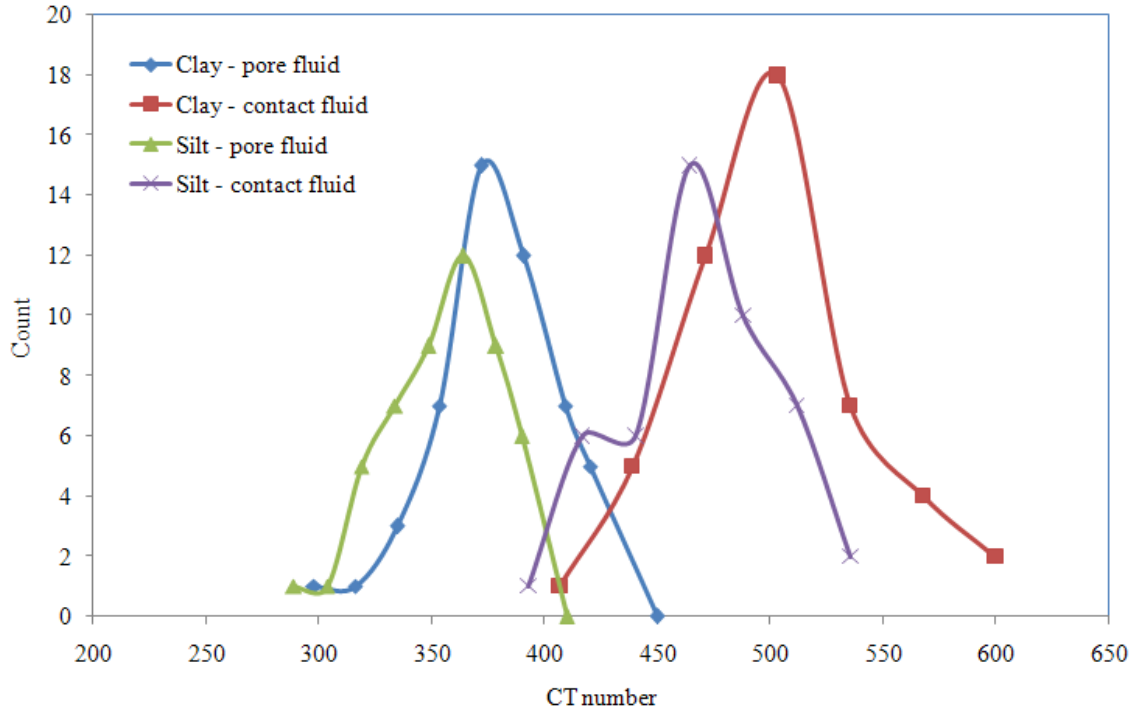


Figure 5.27 Histograms of the pore fluid and contact fluid CT numbers for clay and silt specimens

5.4.7 Measurement of Air-Water Interfacial Area

As discussed in Section 5.2, during drying, the pore fluid in soils goes through different stages, starting from fully saturated, passing through funicular and pendular regimes reaching a fully dry state. Moreover, it can be seen from Equation 2.14 and Figure 2.4 that the capillary forces in unsaturated soils develop across the air water interface. Therefore, it is important to study the changes in the interfacial area between the air and water phases throughout drying to better understand the nature of the physical processes that take place during drying and their impact on the overall capillary forces and small strain stiffness of soils. In this Section, the air-water interfacial area is studied for the different specimens throughout the drying process. The measurements are then used to monitor the transition between the different drying phases.

Amira® Visualization software package was used to perform surface area measurements on all the scanned volumes. In order to perform surface area measurements, the three phases

were separated in each volume, and stored in separate arrays. Then the volumes were resampled in order to obtain smoother surfaces and reduce the effects of pixelization. An Isosurface was then created around each phase using a marching cube algorithm, and measurements of the total surface area of each area were obtained; the three surface area measurements were: the surface areas of the solids (S_s), water (S_w), and air (S_a). The air-water interfacial area was then calculated from the obtained measurements. The obtained surface area for each phase represents the interfacial area of that phase with the other two phases:

$$S_s = S_{s/a} + S_{s/w} \quad (5.6)$$

$$S_a = S_{s/a} + S_{w/a} \quad (5.7)$$

$$S_w = S_{w/a} + S_{s/w} \quad (5.8)$$

Where $S_{s/a}$, $S_{w/a}$ and $S_{s/w}$ are the interfacial areas between solids and air, water and air, and solids and water. Therefore, after measuring the surface area of each phase, Equation 5.9 was used to calculate $S_{w/a}$ for each of the six scanned volumes, similarly, the solid-water interfacial area was calculated using Equation 5.10. The calculated interfacial surface areas are depicted in Figures 5.29 and 5.30.

$$S_{w/a} = \frac{S_w + S_a - S_s}{2} = \frac{S_{w/a} + S_{s/w} + S_{s/a} + S_{w/a} - S_{s/a} - S_{s/w}}{2} \quad (5.9)$$

$$S_{s/w} = \frac{S_w + S_s - S_a}{2} = \frac{S_{w/a} + S_{s/w} + S_{s/a} + S_{s/w} - S_{s/a} - S_{w/a}}{2} \quad (5.10)$$

It can be seen from Figure 5.29 that at the initial stages of drying, the interfacial area is very small because of the very small amount of air in the pores. However, as drying proceeds the amount of air increases in the specimen as the amount of water decreases. At this stage, new fronts of air-water interfaces are created as the amount of air increases; however, most of the water in the specimen still exists as large interconnected bodies, which is a characteristic of the

funicular regime. During this stage, two physical processes that affect the capillary forces take place as drying proceeds: when new air-water interfaces are created capillary forces develop across them resulting in an overall increase in capillary forces in the specimen. On the other hand, the reduction in the amount of water causes some interfaces to dry out, thus losing the capillary forces across them. Therefore, as drying proceeds at this initial stage, these two opposing effects cause the changes in the overall capillary force (and small strain stiffness) to be very small. On the other hand, at later drying stages, a transition is made to the pendular regime, where most of the water bodies are separated and form menisci at the interparticle areas. At this stage, when the interfacial area between air and water reach the maximum values, no more fronts between air and water are created and the interfacial area starts decreasing because of the shrinkage of the water menisci. This reduction of the size of the water menisci plays a very important role in the increase of the capillary forces at this stage as can be demonstrated by Equation 2.14. Therefore, at the later drying stages, a sharp increase in the capillary forces takes place causing an increase in interparticle forces and the overall small strain stiffness. Similar results were obtained by a similar study performed by Culligan et al. (2004). They reported a steady increase in the air-water interfacial area where it reaches a maximum at degrees of saturation between 20 and 35 percent, then starts decreasing. On the other hand, it can be seen from Figure 5.30 that the solid-water interfacial area decreases linearly with drying. This is caused by decrease in the amount of water while the amount of solids stays constant.

5.5 Chapter Summary

It was demonstrated in this chapter that the presence of fine particles in the pore fluid can greatly affect the small strain stiffness of unsaturated soils. By monitoring the P- and S-wave velocities in glass beads specimens during drying, it was shown that adding a small amount of clay or silt to the pore fluid significantly increases the recorded velocity values. This increase

could be caused by the capillary forces within the silt or clay particles in addition to the increase in the interparticle contact stiffness caused by the settlement of the fine particles at the contact area between the glass beads, forming small bridges between the particles and adding to the stiffness of the contacts, and the overall stiffness of the particles.

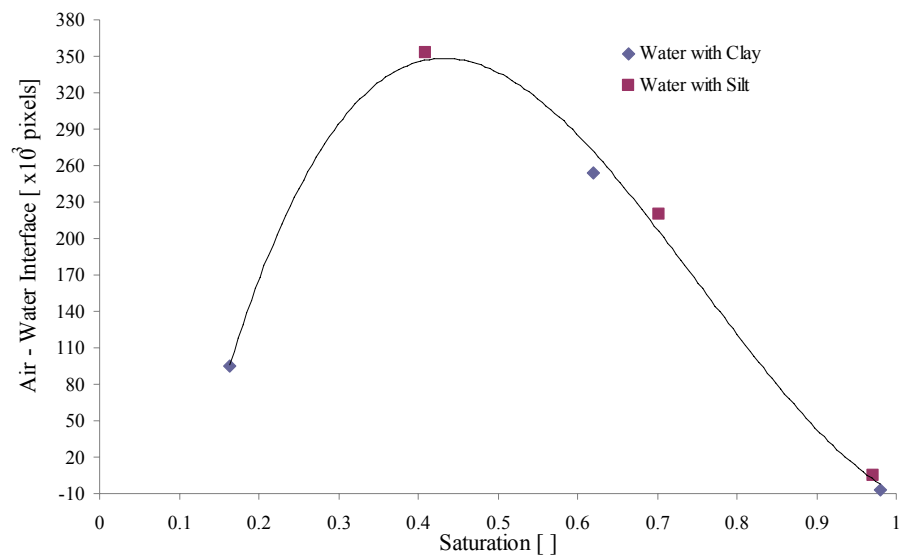


Figure 5.29 Calculated air-water interfacial areas at different drying stages

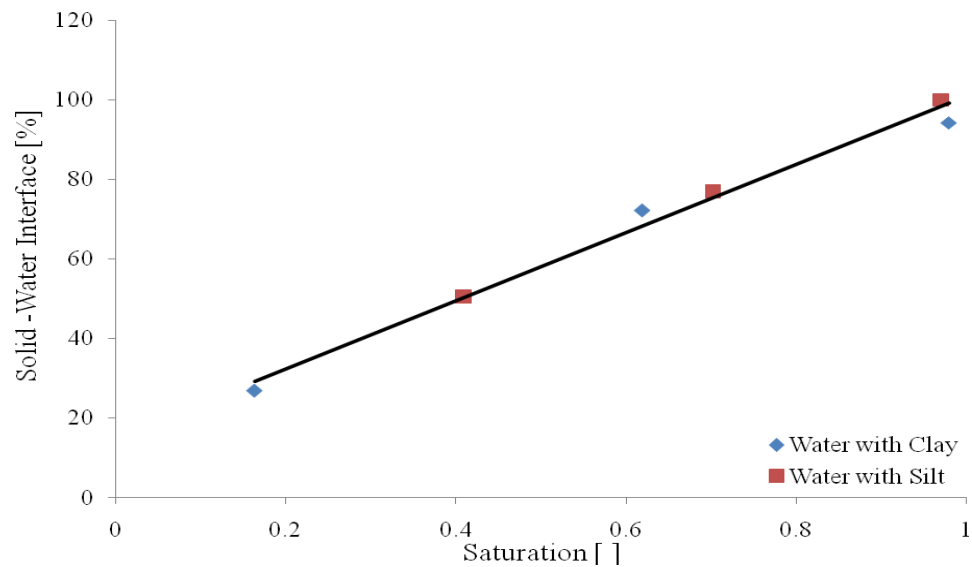


Figure 5.30 Calculated solid-water interfacial areas at different drying stages

It was also shown that synchrotron X-ray tomography is a useful tool for tracking the changes in pore fluid density at different stages of the drying process. The overall density of pore fluid increased during drying due to the increase in the concentration of the fine silt and clay particles, resulting in an increase of the overall stiffness of the glass beads specimen. On the other hand, after studying the local variations in the pore fluid density it was found that the fluids located at the interparticle contact area had a greater concentration of the fine particles than the fluids located in pore bodies. It was also found that the air/water interfacial area steadily increases with drying reaching a maximum value at a degree of saturation of approximately 40%. When the water bodies are separated in the pendular regime, the interfacial area starts dropping because of the shrinkage in the menisci causing a noticeable increase in the capillary forces and small strain stiffness.

CHAPTER 6

CONCLUSIONS AND RECOMMENDATIONS

6.1 Conclusions

This dissertation presented different studies on utilizing non-destructive techniques such as elastic and electromagnetic wave propagation as well as synchrotron x-ray computed tomography to characterize different aspects of the behavior of unsaturated soils. A new method for the determination of field density and moisture content using elastic and electromagnetic waves was presented. Numerical analysis and experimental studies showed that this method was capable of yielding good estimates for the soils' density and moisture content. However, various experimental and numerical error sources resulted in some errors in the estimated density and moisture content values. The density values were in most cases estimated with $\pm 10\%$ of the measured values based on weight-volume measurements whereas the moisture content values were estimated within $\pm 20\%$ of the measured values. To reduce the effects of the error sources and to further improve the suggested methodology, a theoretical framework of a new methodology where the S-wave velocity is added to the acquired measurements was presented. This methodology requires a single V_p , V_s , and θ_v reading instead of the multiple readings required for the original methodology. The suggested methodology and inversion algorithm will be used to develop a prototype of a new instrument for the determination of the in situ density and moisture content of soils. The instrument will also have the ability to monitor the small strain stiffness on in situ soils (see Schneider, 2007) making it a very useful tool for quality control/quality assurance of compaction projects.

On the other hand, a new apparatus was developed to monitor the changes in P- and S-wave velocities in unsaturated soils under controlled net stress and matric suction. This apparatus

was used to determine the effect of matric suction on P- and S- wave velocities in soils, and to monitor the stabilization of capillary forces process after applying different suction increments. It was realized that increasing matric suction, caused an increase in both P- and S-wave velocities due to the increase in the interparticle forces and therefore the overall shear stiffness. It was also noticed that higher wave velocity values were obtained by increasing the percentage of fines in the specimen. A theoretical model that captures the effect of both the net stress and matric suction on the S-wave velocity was presented and was used to determine the relative effects of the physical processes associated with saturated as well as unsaturated soils. The exponent that represents the processes associated with unsaturated soils (β_{unsat}) always have greater values than the exponent that represents the nature of the interparticle contacts (β) highlighting the effect of capillary forces on the small strain stiffness and elastic wave velocity in unsaturated soils.

The effect of the presence of fine particles in the pore fluid on the small strain stiffness of unsaturated soils was also investigated by performing a series of drying experiments on glass beads specimens containing different types of pore fluids while monitoring the P- and S- wave velocities throughout the drying process. The presence of fine particles was found to have a profound effect on the P- and S- wave velocities during drying indicating higher values of the small strain shear, constraint, and bulk moduli. It was also found that adding a small amount of silt or clay to the pore fluid changes the shape of the wave velocity curve during drying. These effects are due to the increase in inter-particle contact stiffness caused by the precipitation of the fine particles at the contact area, as well as the added capillary forces caused by the presence of the fine particles. The results of the x-ray tomography Image analysis showed an increase in the pore fluid density as drying proceeds indicating the migration of the fine particles to the inter-particle contact area at lower degrees of saturation.

The studies presented in this dissertation aimed to highlight the significance and the great value that elastic and electromagnetic wave measurements present for the characterization and monitoring of the effects of different factors on the overall properties of unsaturated soils. It was demonstrated that through monitoring the change in wave velocities, the effects of different factors such as the volumetric water content, matric suction, and pore fluid properties could be quantified and compared. Moreover, other important properties such as the density, moisture content and small strain stiffness could be calculated from these measurements.

6.2 Recommendations for Future Work

The proposed methodology for the determination of density and moisture content could be improved by eliminating or reducing the effect of the potential sources of error. The wave velocity measurement system could be improved by using a repetitive wave source that allows for signal stacking to eliminate the random noise in wave traces and reduce the errors in arrival time measurements. Therefore, the use of bender elements could be a good replacement for the drop ball method. However, the challenge here is to achieve appropriate coupling between the bender elements and the in situ soils especially that it is difficult to imbed the bender elements in stiff soils. Moreover, using bender elements allows for the generation of both P- and S- waves, allowing for the implementation of the alternative methodology presented in Section 3.6. It is also recommended that the alternative methodology should be verified by an experimental program similar to the one performed for the original methodology. Although this methodology theoretically presents a very promising alternative, it still needs to be experimentally verified and tested for different types of soils to assure its proper operation.

The new triaxial apparatus could be used to perform additional experiments to determine the effect of net stress on S- and P- wave velocities to be compared to the effect of matric suction. A two dimensional surface could be developed combining the effects of net stress and

matric suction on different soil types. The experimental program could also be extended to testing soils with different stress paths. The cell could be used to perform triaxial testing on unsaturated soil specimens while monitoring the changes in elastic wave velocities. These tests could also be performed at varying net stress and matric suction values to determine their relative effects on the overall shear strength of the unsaturated soil specimens.

The effect of fine particles on the small strain stiffness could be further investigated by testing different soil as well as different pore fluid types. Different particles shapes result in different types of interparticle contacts, therefore, the developed drying cell could be used to determine the effects of different factors such as particle shape, fine particle concentration and drying rate on small strain stiffness on unsaturated soils. Moreover, synchrotron computed tomography could be used to further investigate the properties and behavior of pore fluids during drying as well as wetting processes.

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